

Study of Ice Accretion and Icing Effects on Aerodynamic Characteristics of DU96 Wind Turbine Blade Profile

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Abstract

In order to optimize the large wind turbines operation in ice prone cold regions, it is important to better understand the ice accretion physics and its effects on aerodynamic performance and power production losses. This paper describes a case study of ice accretion on DU96-W-180 airfoil, which has been used for large wind turbine blades such as NREL 5MW. Analysis has been carried out for glaze and rime ice conditions using icing tunnel experimental data and multiphase Computational Fluid Dynamics (CFD) based numerical approach. Results show a difference in profile surface roughness and heat fluxes during rime and glaze ice accretion process, which leads to a significant change in rate and shape of ice accretion. More complex ice shapes are observed in case of glaze ice conditions that affects the aerodynamic performance differently from rime ice conditions. Numerical results are compared with the experimental data, where a good agreement is found. Results show higher aerodynamic performance degradation for glaze ice conditions particularly at higher angles of attack.

Keywords: Ice accretion; Wind turbine; Aerodynamic; Icing tunnel; Surface roughness; CFD.

1. Introduction

In recent years, activities regarding wind energy projects in ice prone cold regions have increased due to availability of good wind resources, but atmospheric icing on wind turbine blades is considered as a potential hazard in proper utilization of these good wind resources. Atmospheric icing affects the wind turbine aerodynamic performance, which leads to a decrease in wind energy production (Muhammad S. Virk, 2011). Worldwide, installed wind energy capacity in ice prone regions is expected to reach 123 GW in year 2020 (Paul Dvorak, 2013). Wind energy production losses due to icing have been reported to lead up to a 17% decrease in Annual Energy Production (AEP) and 20-50% in aerodynamic performance (Ozcan Yirtici, 2016). This highlights the importance of finding the innovative solutions for wind turbine operations in icing conditions. There is a growing need to

30 improve knowledge about aerodynamic design and performance of large wind turbine rotor blades for optimal
31 operations in icing conditions. The icing conditions within cold climates are insufficiently included in the design
32 limits presently covered by the national and international standards for wind turbine design. The International
33 Energy Agency (IEA) Task 19: '*Wind energy in cold climates*' has also urged the development of new methods to
34 enable better understanding of the effects of ice accretion on wind turbine performance and energy production
35 (Timo Laakso, 2009).

36

37 Atmospheric ice accretion on wind turbine blades is caused by the impingement of super-cooled water droplets.
38 Atmospheric ice is mainly classified as dry rime and wet glaze ice. Rime ice is soft and less dense and happens at
39 very low temperature, when 100% impinging droplets freeze, whereas glaze ice is harder and denser. Glaze ice
40 happens close to freezing temperature, when freezing fraction of impinging droplet is not 100% and some droplets
41 run along blade surface as very thin water film. Accreted ice has a range of shapes resulting from different
42 temperatures and heat balance situations that causes different levels of aerodynamic performance losses. Duncan
43 et al. (T. Duncan, 2008) made differentiation between rime and glaze ice shapes and found that glaze ice caused
44 larger loss than rime ice. Virk et al. (Muhammad S. Virk, 2010) found that the ice mainly accretes along leading
45 edge of the blade and reduces torque. This effects the wind turbine aerodynamic performance and results in Annual
46 Energy Production (AEP) losses. The parameters causing wind turbine blade aerodynamic penalty can be divided
47 into three categories: *small-scale surface roughness*, *large-scale surface roughness* and *ice geometry*. The shape
48 of the accreted ice along the wind turbine blade depends upon many variables such as point of operation, the
49 geometry of the wind turbine blade, relative wind velocity, atmospheric temperature, droplet diameter and the
50 liquid water content (Drage Peter, 2009). Better understanding of ice accretion effects on the blade aerodynamics
51 can help to optimize its design and reduce the AEP losses.

52

53 Ice accretion physics and its effects on wind turbine's aerodynamic performance can be analysed using both
54 experimental and numerical approaches. In recent years, advance CFD based numerical techniques have begun to
55 play a significant role both in simulating and determining the performance of wind turbine blades under icing
56 conditions (Gitsuzo d. Tagawa, 2018; Jia Yi Jin, 2018; Jian Liang, 2018; Pavlo Sokolov, 2018; W. A. Timmer,
57 2010; Woobeom Han, 2018). Lab based icing tunnel experimenters provide more accurate picture of ice accretion,
58 however, it have limited insight of the air flow and droplet behaviour, whereas CFD simulations can provide the

59 economical insight details of air flow physics and droplet behaviours, which is difficult to obtain from icing tunnel
60 experiments.

61

62 DU96-W-180 airfoil has been an attractive choice for large wind turbine blades and has been used for NREL 5MW
63 wind turbine. This is an asymmetric profile with 18% thickness and has aerodynamic features, such as high-lift-
64 to-drag ratio, insensitivity to contaminations and low noise (Hao Guo, 2017; R. P. J. O. M. van Rooij, 2003; W.
65 A. Timmer, 2003). Many Researchers have carried out analysis to investigate DU series wind turbine blade profiles
66 using experimental and numerical techniques (Muhammad S. Virk, 2010; W. A. Timmer, 2001). Linyue Gao et
67 al. have done the experimental study of ice mitigation for DU96-W-180 airfoil.(Linyue Gao, 2018; Linyue Gao,
68 2017), but not much work has been carried out by the researchers to study the performance of DU96 airfoil under
69 icing conditions. Matthew C. Homola et al. (Matthew C. Homola, 2012) have done CFD based numerical study to
70 simulate the ice on DU96 airfoil as part of NREL 5MW wind turbine study, but no published work is available
71 regarding icing tunnel experimentation of DU96 airfoil.

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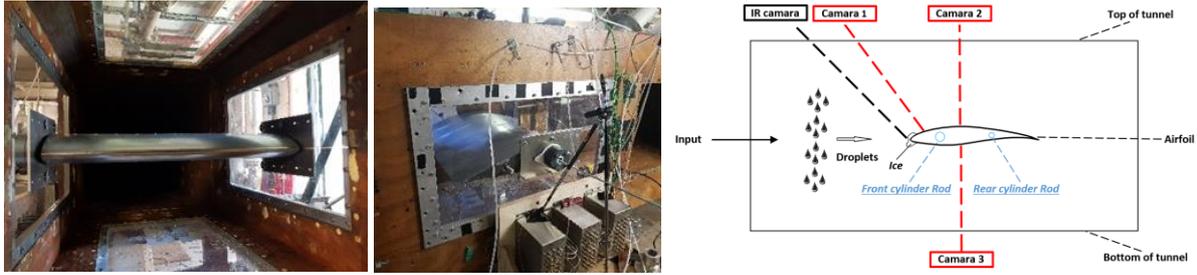
73 This paper is focused on better understating of ice accretion physics and its resultant effects on aerodynamic
74 performance of DU96-W-180 airfoil. Both icing tunnel experiments and CFD based numerical simulations have
75 been carried out to better investigate the ice accretion physics and resultant rate and shapes of accreted ice for both
76 rime and glaze ice conditions. Icing tunnel experiments are carried out at Cranfield University UK, whereas CFD
77 based numerical simulations are carried out using ANSYS-FENSAP ICE and FLUENT.

78

79 **2. Experimental Analysis**

80 **2.1 Experimental Setup**

81 The experimental study is carried out at icing wind tunnel laboratory of Cranfield University (CU), UK. DU96
82 wind turbine blade profile model with a span of 758 mm and chord length of 500 mm was used. The surface of
83 blade profile model was made of Galvanized steel (VGAL.V.D×SID+Z275) with an average surface roughness of
84 0.9 microns. Icing wind tunnel facility of CU is able to create natural icing conditions. This icing tunnel has test
85 section (761×761 mm) and can operate for Medium Volume Diameter (MVD) ranging from 15-80 microns, Liquid
86 Water Content (LWC) from 0.05-3 g/m³ and air temperature from -30 to +30 °C. Figure 1 shows the icing wind
87 tunnel setup and test cross section.



88
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Figure 1. CU icing wind tunnel setup.

90 To closely monitor the ice accretion on the blade profile, three different High Definition (HD) cameras (one for
91 side, one for top and one for top view) were used for video recording and images. Accreted ice shapes were
92 extracted and sketched after each experiment. These experiments were carried out at Reynolds number = 3×10^6
93 for both dry rime and wet glaze ice conditions. Table 1 shows the operating conditions used for the experiments.

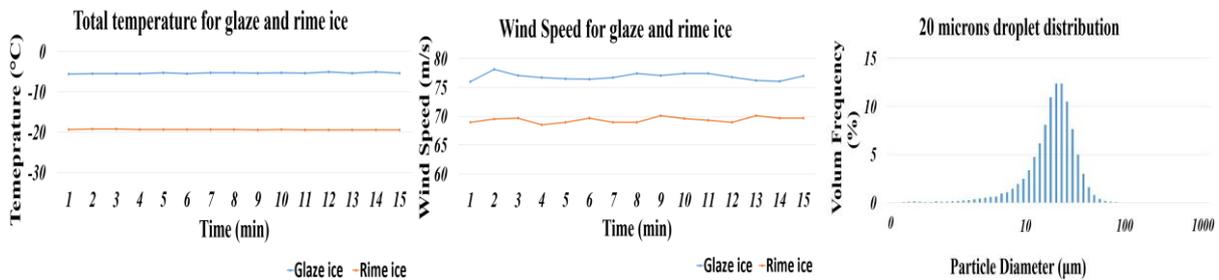
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Table 1. Icing tunnel operating conditions.

Test	Ice Type	Velocity (m/s)	Temperature (°C)	LWC (g/m ³)	MVD (microns)	AOA (degree)	Icing time (mins)
1	Glaze	77	-5	0.35	20	0	15
2	Rime	70	-20				

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To closely monitor the icing tunnel operations, various operating parameters of icing tunnel are closely monitored
to make sure the smooth operations. MVD of 20 microns is used with the droplet distribution spectrum consists of
60 bins. Figure 2 shows the droplet distribution spectrum, variation of icing tunnel air speed and total air
temperature at tunnel test section for both rime and glaze ice conditions.

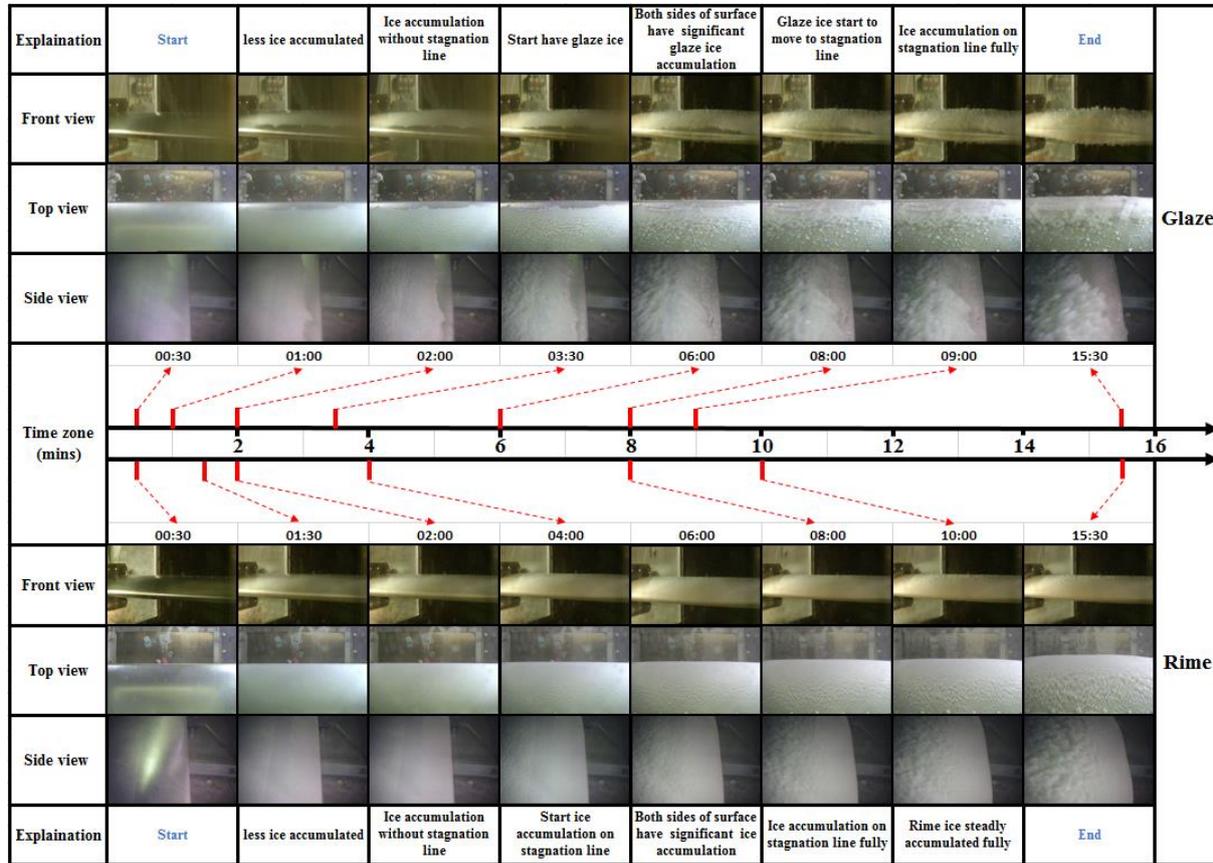


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Figure 2. Variation of wind speed, temperature and droplet distribution spectrum in CU icing tunnel.

105 **2.2 Experimental Results**

106 Figure 3 shows the experimental ice growth with time where results show a significant difference in ice growth
 107 for both rime and glaze ice conditions. Views from three different HD cameras were used to monitor the ice growth
 108 for this study.



109
 110 Figure 3. Overview of experimental ice growth at various time steps for 15 minutes.

111 Figure 3 shows the ice accretion process at different time intervals for both rime and glaze ice conditions. Results
 112 show that ice mainly accreted along leading edge of the blade profile. Analysis shows that for rime conditions,
 113 accreted ice was dense and shape along stagnation line was smooth. Further down from leading edge, tight grain
 114 hard rime ice with direction of feather growing parallel to the flow is observed. For rime conditions, ice accretion
 115 extended till 20-25% of the profile chord length. For glaze conditions, clear ice is observed along stagnation line
 116 with horn shape. Large individual feathery spikes, pointing perpendicular to the profile surface is observed. After
 117 each experiment, the ice shapes were extracted from center section of the blade profile. Experimental ice shapes
 118 with cut-outs and ice thicknesses are shown in Figure 4. Streamline ice shapes are observed for the rime ice
 119 conditions, whereas complex horn type shapes are found in case of glaze ice conditions.

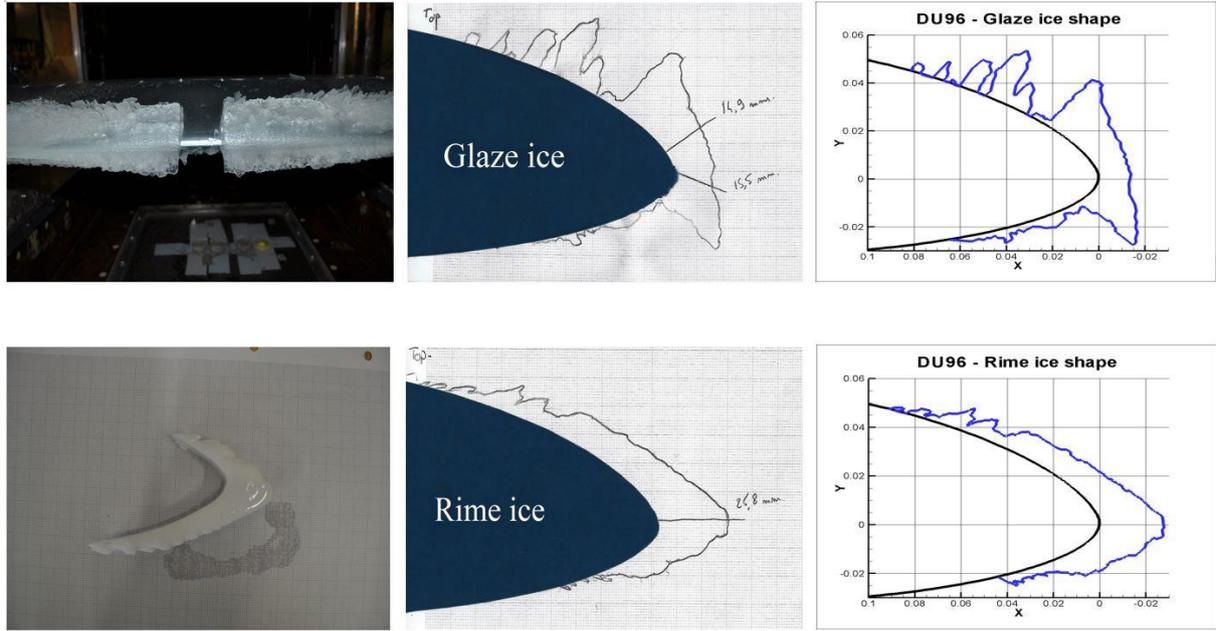


Figure 4. Experimental ice shapes for glaze and rime ice conditions.

3. Numerical Analysis

3.1 Numerical Setup

CFD based multiphase numerical simulations are carried out using ANSYS-FENSAP and FLUENT to study the airflow behaviour, droplet behaviour and to simulate the surface thermodynamics and resultant ice accretion. Aerodynamic coefficients (C_L and C_D) of clean and iced profiles are calculated and compared. Ice shapes from numerical simulations are compared with the experimental data. CFD simulations provided a detailed overview of the airflow and droplet behavior, which is not easy to study from experiments. Atmospheric ice accretion on blade profile can be numerically simulated by means of integrated thermo-fluid dynamic models, which involve the fluid flow simulation, droplet behaviour, surface thermodynamics and phase changes. Airflow behaviour is simulated by solving the nonlinear partial differential equations for the conservation of mass, momentum and energy.

$$\partial \rho_\alpha / \partial t + \vec{\nabla}(\rho_\alpha \vec{v}_\alpha) = 0 \quad (1)$$

$$\partial \rho_\alpha \vec{v}_\alpha / \partial t + \vec{\nabla}(\rho_\alpha \vec{v}_\alpha \vec{v}_\alpha) = \vec{\nabla} \cdot \sigma^{ij} + \rho_\alpha \vec{g} \quad (2)$$

$$\partial \rho_\alpha E_\alpha / \partial t + \vec{\nabla}(\rho_\alpha \vec{v}_\alpha H_\alpha) = \vec{\nabla}(\kappa_\alpha (\vec{\nabla} T_\alpha) + v_i \tau^{ij}) + \rho_\alpha \vec{g} \vec{v}_\alpha \quad (3)$$

Where ρ is the density of air, v is the velocity vector, subscript α refers to the air solution, T refers to the air static temperature in Kelvin, σ^{ij} is the stress tensor and E and H are the total energy and enthalpy, respectively. The

139 sand grain roughness for the iced surface is calculated using following Shin et al. roughness model . (Shin Jaiwon,
140 1992)

$$141 \quad [(k_s/c)/((k_s/c)_{base})]_{MVD} = \begin{cases} 1, & MVD \leq 20 \\ 1.667 - 0.0333(MVD), & MVD \geq 20 \end{cases} \quad (4)$$

142

143 Where MVD is the droplet Median Volume Diameter (in microns), whereas the corresponding value of sand-grain
144 roughness is obtained:

145

$$146 \quad k_s = 0,6839[(k_s/c)/((k_s/c)_{base})]_{LWC}[(k_s/c)/((k_s/c)_{base})]_{T_s}[(k_s/c)/((k_s/c)_{base})]_{MVD}(k_s/c)_{base}c$$

147 (5)

148 Two phase flow (*air and water droplets*) is numerically simulated using the Eulerian approach, where the super
149 cooled water droplets are assumed to be spherical. The Eulerian two phase fluid model consists of the Navier-
150 Stokes equation with the water droplets continuity and momentum equation. The water droplet drag coefficient is
151 based on the empirical correlation for the flow around the spherical droplets described by Clift et al.(R. Clift, 1978)

152

$$153 \quad \partial\alpha/\partial t + \vec{\nabla}(\alpha\vec{V}_d) = 0 \quad (6)$$

$$154 \quad \partial(\alpha\vec{V}_d)/\partial t + \vec{\nabla}(\rho_\alpha\vec{V}_dH_d) = C_D Re_d/24k \alpha(\vec{V}_a - \vec{V}_d) + \alpha(1 - \rho_a/\rho_d) 1/(Fr^2) \vec{g} \quad (7)$$

155

156 Where α is the water volume fraction, V_d is the droplet velocity, C_D is the droplet drag coefficient and Fr is the
157 Froude number. The numerical study is carried out for custom droplet distribution spectrums at MVD of 20
158 microns. Surface thermodynamics is calculated using the mass and energy conservation equations, considering the
159 heat flux due to convective and evaporative cooling, heat of fusion, viscous and kinetic heating.

160

$$161 \quad \rho_f[\partial h_f/\partial t + \vec{\nabla}(\vec{V}_f h_f)] = V_\infty LWC\beta - \dot{m}_{evap} - \dot{m}_{ice} \quad (8)$$

162

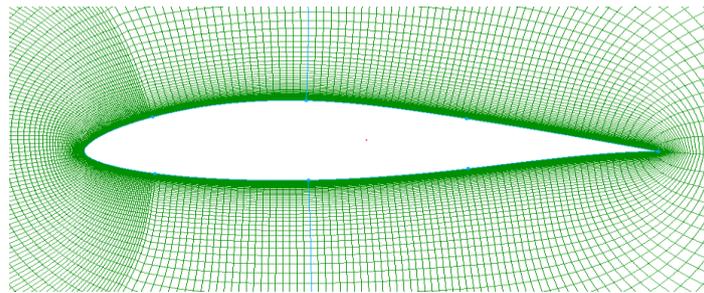
$$163 \quad \rho_f[(\partial h_f c_f \tilde{T}_f)/\partial t + \vec{\nabla}(\partial h_f c_f \tilde{T}_f)] = [c_f(\tilde{T}_\infty - \tilde{T}_f) + (\|\vec{V}_d\|^2)/2] V_\infty LWC\beta - L_{evap} \dot{m}_{evap} + (L_{fusion} -$$

164 $c_s \tilde{T}) \dot{m}_{ice} + \sigma \varepsilon (T_\infty^4 - T_f^4 - c_h(\tilde{T}_f - \tilde{T}_{ice,rec}) + Q_{anti-icing} \quad (9)$

165

166 The coefficients ρ_f , c_f , c_s , σ , ε , L_{evap} , L_{fusion} are physical properties of the fluid. The reference conditions
167 \tilde{T}_∞ , V_∞ , LWC are the airflow and droplets parameters. 3D grid is generated by extruding a single cell layer in the

168 span wise direction. ALE (Arbitrary Lagrangian Eulerian) formulation is used for the grid displacement during
 169 ice accretion, which adds the grid speed terms to the Navier-Stokes equations to account for the mesh velocity
 170 (Manual, 2010). Mesh sensitivity study was carried out using coarse, medium and fine meshes to accurately
 171 determine the boundary layer characteristics (shear stress and heat flux). For each case, the mesh was automatically
 172 displaced after each time shot to account for the ice growth without any change in mesh size. During mesh
 173 sensitivity analysis, number of mesh elements and y^+ value less than 1 for first cell layer was selected based upon
 174 the heat flux calculations, where a numerical check was imposed that the heat flux computed with the classical
 175 formulae dT/dn should be comparable with the heat flux computed with the Gresho's method. Mesh sensitivity
 176 study showed that the effect of mesh size on droplet solution was negligible, however some flow quantities
 177 including convective heat flux on the blade surface was sensitive to the mesh size, resulting in higher instantaneous
 178 ice growth in regions with higher convective heat loss. After mesh sensitivity analysis, O type structured numerical
 179 grid with y^+ value less than 1 and total 48,450 hexahedral elements was used for final simulations.



180

181

Figure 5. Numerical grid for DU96-W-180 airfoil.

182 K-omega SST turbulence model is used as a compromise between acceptable computational cost and required
 183 accuracy for simulating the turbulent flow. Sand grain roughness height for the iced surface was calculated with
 184 an empirical correlation described by Shin et al. (Shin Jaiwon, 1992). Numerical simulations are carried out at
 185 conditions specified in Table 2.

186

Table 2. Numerical setup

Chord length (m)	0.5
Angle of attack (AOA, degree)	0
Air velocity (m/s)	70 (rime), 77 (glaze)
Temperature (Celsius)	-20 (rime), -5 (glaze)
MVD (microns)	20
Droplet distribution	Custom distribution from CU (see Figure 2)

LWC (g/m³) 0.35

Icing time (mins) 15

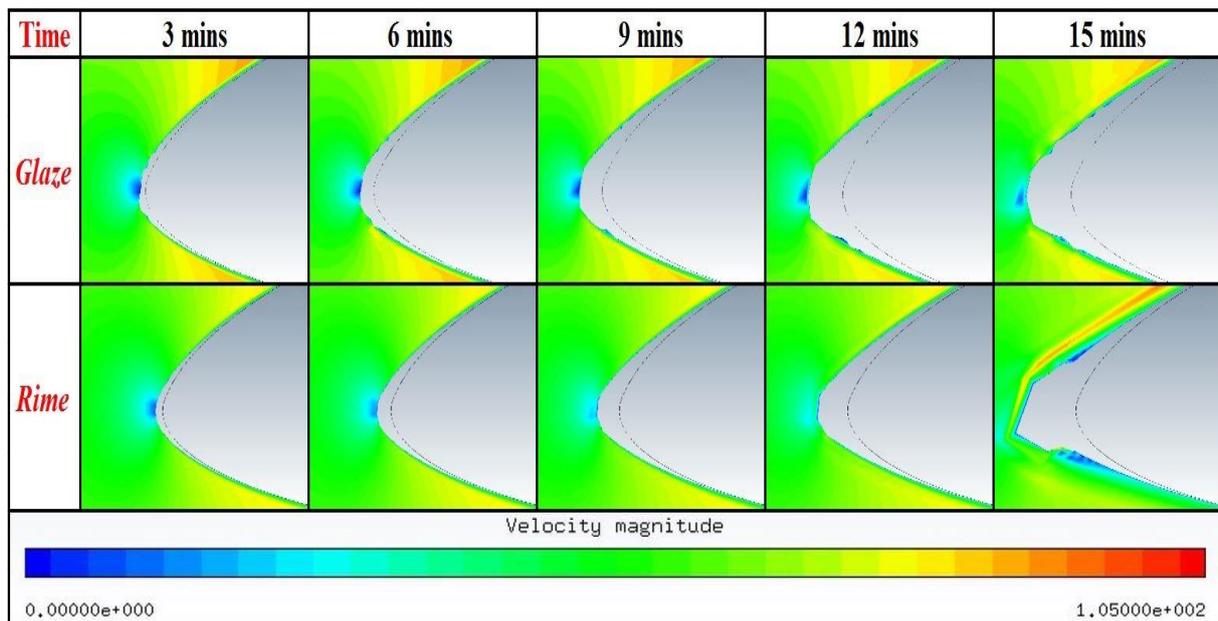
187 3.2 Numerical Results

188 Numerical simulations have been carried out to study the ice accretion physics and its effects on airflow and droplet
189 behaviour. In this numerical study, air flow and droplet behaviour is analyzed at different time intervals during ice
190 accretion, which is not easily possible during icing tunnel experiments. Quasi-steady state multi-shot numerical
191 simulations approach is used where total icing duration (15 minutes) is divided in 23 time steps. Detailed analysis
192 have been carried out to analyses the results at each time step.

193

194 3.2.1 Airflow behavior

195 Analysis of airflow behavior show a change in velocity and pressure distribution along pressure and suction sides
196 of the blade profile. A shift in the position of stagnation point is observed as ice accretes along leading edge. Not
197 much flow separation is observed in case of rime ice conditions as compared to the glaze ice conditions. Figure 6
198 shows the velocity contours along leading edge for glaze and rime ice conditions at different time steps.



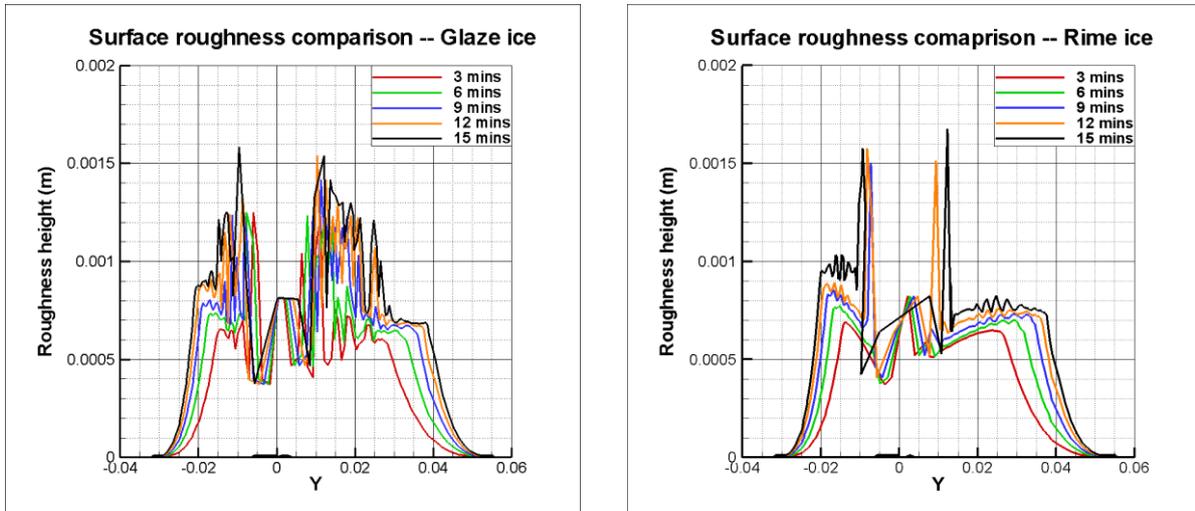
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Figure 6. Velocity magnitude at different time steps

201 During the ice accretion process, the surface roughness of the blade profile changes significantly, which effects
202 the shear stresses and heat fluxes. This leads to a change in the boundary layer thickness which influences the
203 convective heat transfer and droplet sticking efficiency along the blade surface. To better understand the effects
204 of surface roughness change on shear stress and heat flux, analysis has been carried out. Figure 7 shows the surface

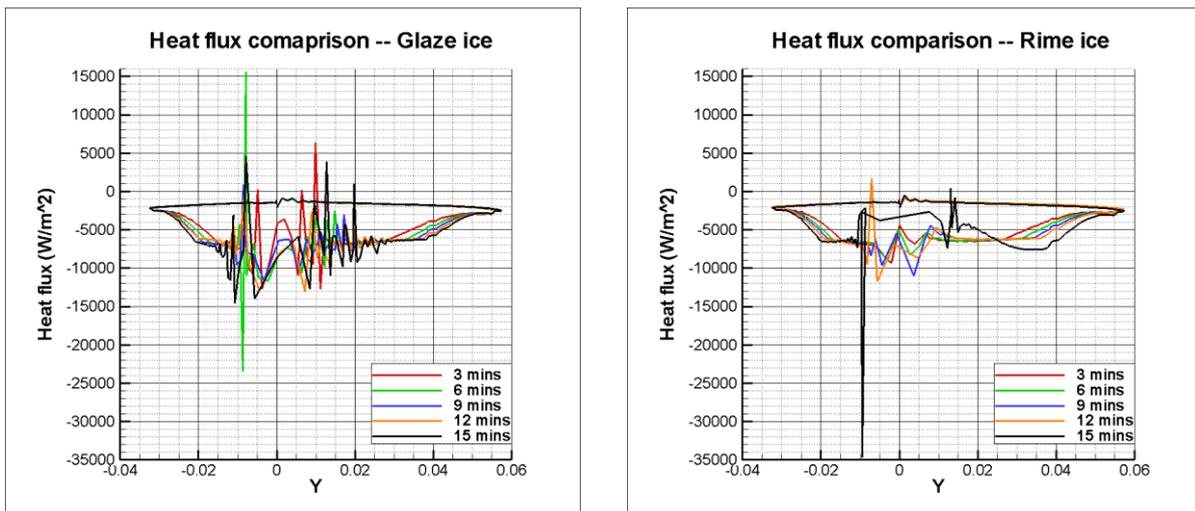
205 roughness change during the ice accretion process whereas Figure 8 shows the heat transfer comparison for both
206 rime and glaze ice conditions.



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Figure 7. Surface roughness variation during ice accretion.



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Figure 8. Heat flux variation during ice accretion.

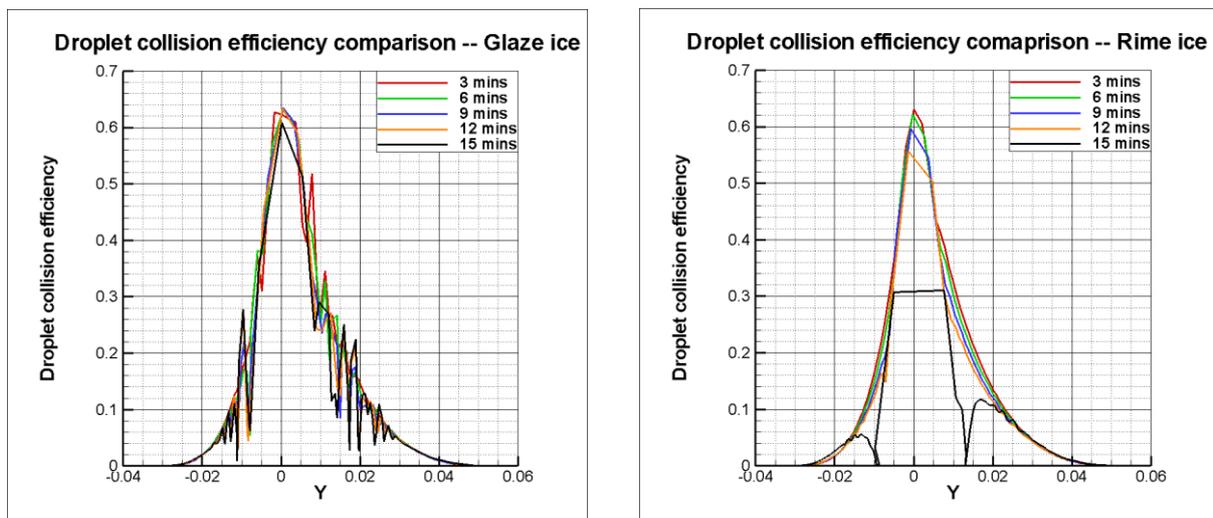
211 Analysis shows that surface roughness and heat flux increases as ice accretion process progresses. Results show
212 higher values of the surface roughness along leading edge of the blade profile, mainly due to higher droplet
213 collision which leads to more ice accretion. Leading edge sections with higher surface roughness also shows higher
214 values of surface heat transfer during ice accretion.

215

216 3.2.2 Droplet Behaviour

217 The droplet behaviour at each time step is analysed, where results show a change in droplet collision efficiency
218 with the change of blade profile shape during ice accretion process. Figure 9 shows the droplet collision efficiency

219 variation at five different time intervals for rime and glaze ice conditions. Results show that droplet collision
 220 efficiency is higher at start of the ice accretion process, but as ice accretes, a decrease in the droplet collision
 221 efficiency is observed. This change in droplet collision is mainly due to change in blade profile shape particularly
 222 at leading edge. Along leading edge significant ice accretes due to higher droplet collision efficiency, whereas
 223 water runback process also occurs during ice accretion particularly for glaze ice conditions which results in small
 224 patches/feathers of ice along other areas of the blade profile. This leads to a change in droplet collision efficiency
 225 along blade profile during ice accretion process. Figure 9 shows the droplet collision efficiency variation during
 226 ice accretion process for both rime and glaze ice conditions.



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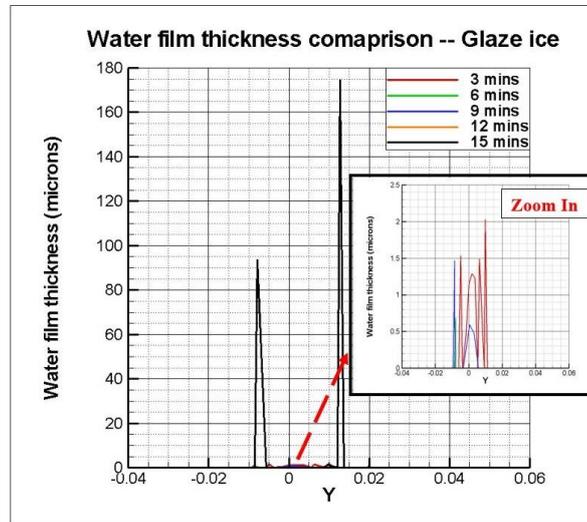
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Figure 9. Droplet collision efficiency at different time intervals.

229 3.2.3 Ice Accretion

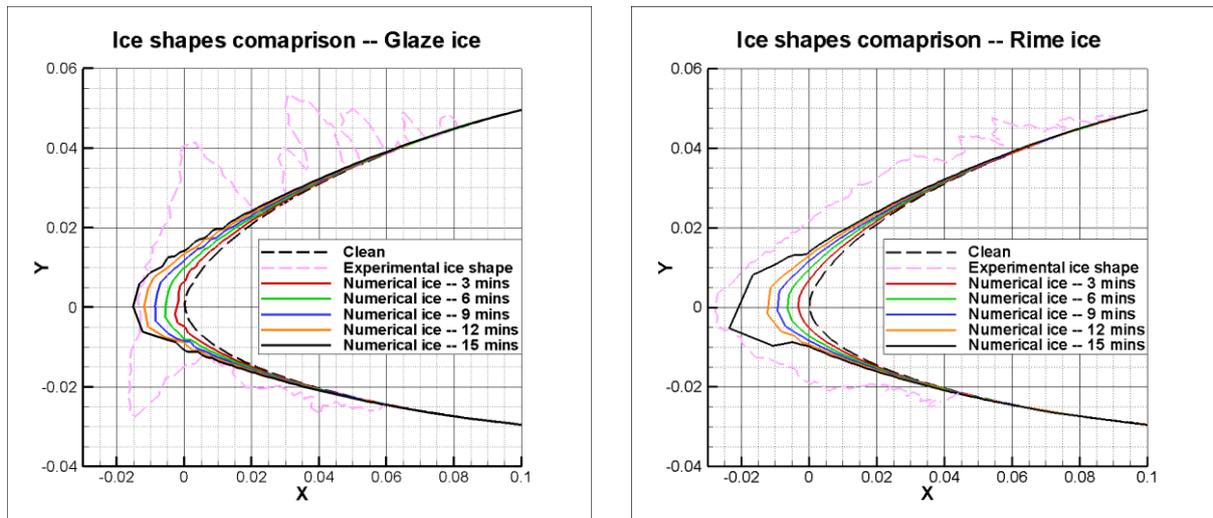
230 Figure 10 shows the water run back film thickness variation along blade profile surface during ice accretion
 231 process. Due to very low temperature for rime ice conditions, the droplet freezing fraction is 100% due to which
 232 no water run back phenomena is observed for rime ice conditions, whereas for glaze ice conditions, analysis shows
 233 water run back along profile section. Figure 11 shows the ice shapes of DU96 blade profile obtained from
 234 experimental and numerical simulations. A reasonably good agreement is found for the ice shapes particularly for
 235 the ice growth along stagnation point of the blade profile for both rime and glaze ice conditions. For glaze ice
 236 conditions, due to water run back and high aerodynamic heat flux along leading edge, complex horn shapes are
 237 obtained during the experiments. These ice horns are not captured very well in the numerically simulated ice
 238 shapes. For glaze ice conditions, when high speed water droplets collide with the airfoil surface, they do not get
 239 enough time to freeze and following droplets hit. Incoming airflow pushes these droplets further away from the
 240 stagnation line which leads to accumulation of these droplets along upper and lower sides of the airfoil as thin

241 water film. This phenomena is not easy to capture during numerical simulations. To better understand this, water
 242 film thickness was simulated at each time step, shown in Figure 10, where results show a clear trend of thick water
 243 film along leading edge.



244

245 Figure 10. Water runback film thickness for glaze ice conditions.



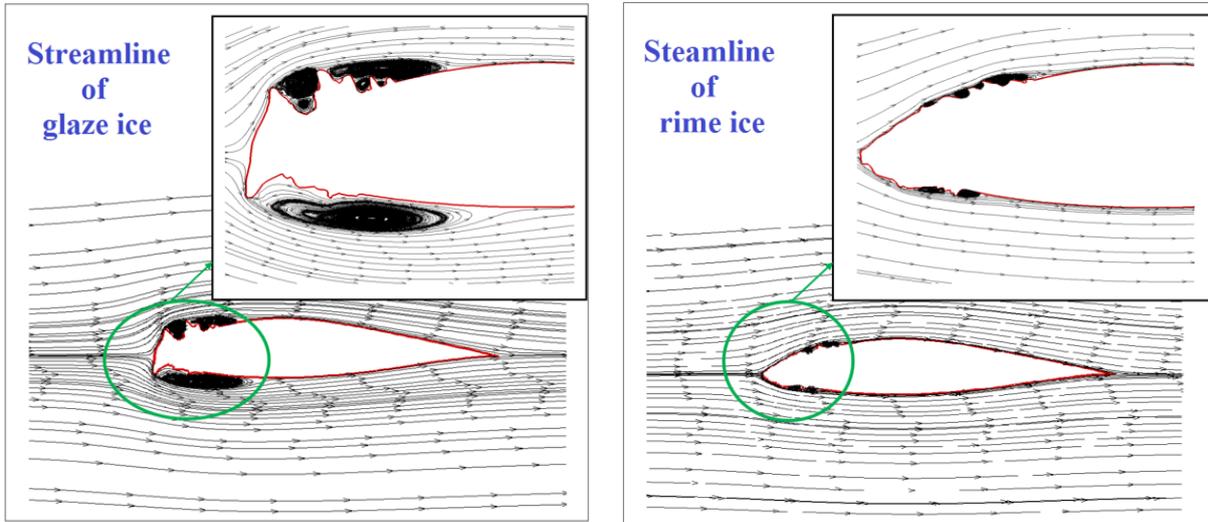
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247 Figure 11. Accreted ice shapes comparison for rime and glaze ice conditions.

248 **4. Aerodynamic Analysis of Experimental Ice Shape Profiles**

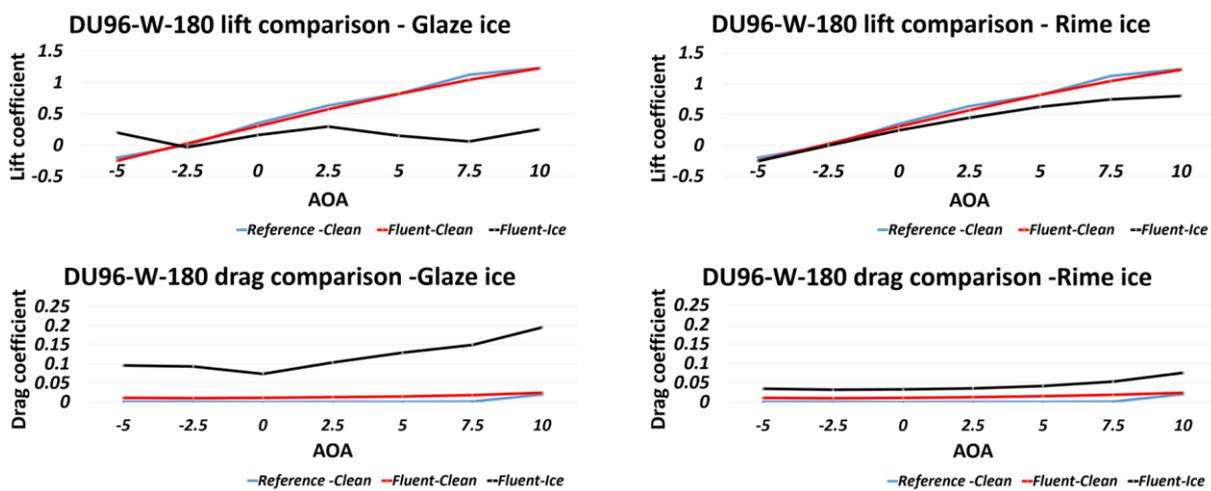
249 Airflow and droplet behaviour along blade profile surface is influenced by the accreted ice shape. During ice
 250 accretion process the shape of the blade profile surface changes significantly, which also changes the flow behavior
 251 and leads to a change in aerodynamic performance of the blade profile. This section presents a detailed numerical
 252 study to analyze the change in aerodynamic performance of the DU96. Experimental ice shapes (after 15 minutes)
 253 are used for this study, where a comparison is made between aerodynamic characteristics of the clean and

254 experimental iced blade surfaces. Figure 12 shows the velocity streamlines along the experimental ice shape
 255 obtained from rime and glaze ice conditions.



256
 257 Figure 12. Velocity streamlines along iced DU96 blade profiles.

258 To analyse the change in aerodynamics characteristics of the iced profile, a detailed parametric study has been
 259 carried out at different angles of attack. The iced profile shapes obtained from experiments are used where the
 260 flow is simulated at different AOA and comparison is made with the aerodynamic characteristics of clean DU96
 261 airfoil. These simulations are carried out using ANSYS-FLUENT. Figure 13 shows the lift and drag coefficients
 262 comparison.



263
 264 Figure 13. Comparison of lift and drag coefficients.

265 Analysis of Figure 13 shows a decrease in the aerodynamic characteristics of ice profile when it is compared with
 266 the clean profile. This change in the aerodynamic characteristics is mainly due to a change in the blade profile
 267 aerodynamic shape due to ice and surface roughness, which affects the boundary layer flow behaviour. At all

268 angles of attack, a significant increase in the drag coefficient is observed due to flow separation. As an overall,
269 higher drag coefficient is found for glaze ice shapes, which leads to higher aerodynamic performance loss.

270 **5. Conclusion**

271 This ice accretion study for DU96 airfoil showed more streamlined ice shape for dry rime ice conditions, whereas
272 complex horn type ice shape was found for case of glaze ice. This is mainly due to difference in droplet freezing
273 fraction, as due to low freezing fraction for glaze ice conditions, higher water run back and aerodynamic heat flux
274 along leading edge is observed which resulted in complex horn type ice shapes. This phenomenon is difficult to
275 capture during numerical simulations and requires further improvements in the existing numerical models of wet
276 ice accretion. Numerical results showed that blade profile surface roughness and heat flux changes significantly
277 during ice accretion process which effects the airflow and droplet behavior. The change in accreted ice shape
278 effects the airflow behavior and aerodynamics performance. Analysis show a decrease in the aerodynamic
279 characteristics of the iced airfoils when it is compared with the clean. This degradation in aerodynamic
280 performance is higher in case of glaze ice, as compared to the rime iced profile. Icing tunnel experiments provided
281 a good overview of the ice accretion process, but it's not easy and economical to study the airflow and droplet
282 behaviour during ice accretion process in experimental study. Therefore CFD based numerical techniques can be
283 a good tool to study the change in air flow and droplet behaviour during ice accretion process.

284

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290

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