

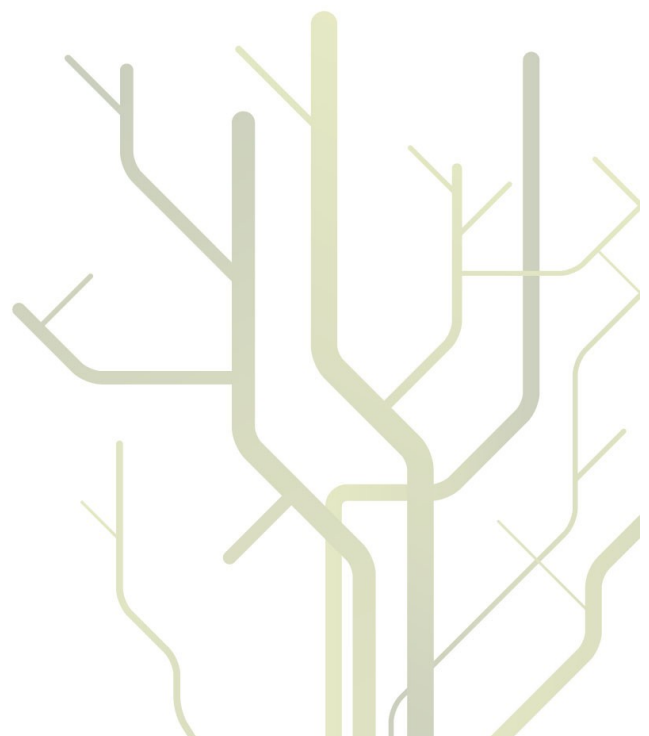
Statistical Analysis of Multilook Polarimetric Radar Images with the Mellin Transform



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A dissertation for the degree of
Philosophiae Doctor

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To my family

«The true sign of intelligence is
not knowledge but imagination.»

Albert Einstein

«Nobody creates anything. It's there,
and you just f***** grab a hold of it.»

Keith Richards

Abstract

This thesis presents methods for statistical analysis of the probability distributions used to model multilook polarimetric radar images. The methods are based on a matrix-variate version of Mellin's integral transform.

The proposed theoretical framework is referred to as Mellin kind statistics. It is an extension of a theory recently developed for single polarisation amplitude and intensity data to the complex matrix-variate case describing multilook polarimetric images. This generalisation is made possible by the rediscovery of a generalised Mellin transform, which is defined for functions of positive definite Hermitian matrices. The domain makes it suited for application to the distributions used to model the polarimetric covariance and coherency matrix.

The analysis tools include the matrix-variate Mellin kind characteristic function, which is defined with the Mellin transform in place of the conventional Fourier transform. Matrix log-moments and matrix log-cumulants are retrieved from this function. The matrix log-cumulants are used in a moment based approach to parameter estimation of the distribution parameters. The estimators make efficient use of all the statistical information in the polarimetric covariance matrix, and are superior to all known alternatives. The matrix log-cumulants are also used to construct the first known goodness-of-fit test for matrix distributions based on the multilook polarimetric product model. The algorithms are interpreted by means of a highly informative graphical visualisation tool displaying a space spanned by certain matrix log-cumulants.

It is demonstrated that the matrix-variate Mellin transform is the natural tool for analysing multilook polarimetric radar images. This conclusion is based on the simple and elegant mathematical expressions obtained, the superb statistical properties of developed estimators, as well as the intuitive interpretations offered by the Mellin kind statistics.

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Contents

Abstract	i
Acknowledgements	iii
Table of Contents	vii
List of Tables	ix
List of Figures	ix
Nomenclature	xi
List of Notation	xi
List of Acronyms	xiii
1 Introduction	1
1.1 Motivation	1
1.2 Chapter Review	3
1.3 Publication Review	4
1.4 Other Publications and Presentations	7
2 Radar Polarimetry	9
2.1 Imaging Radar	9
2.1.1 Frequency Bands	9
2.1.2 Properties of Imaging Radar	10
2.1.3 Instruments	11
2.1.4 Synthetic Aperture Radar	12
2.1.5 Polarimetric Radar	13
2.2 Data Formats	14
2.2.1 Single Look Complex Data	14
2.2.2 Scattering Matrix	15
2.2.3 Scattering Vector	15
2.2.4 Multilook Complex Data	17
2.2.5 Covariance Matrix	18
2.2.6 Coherency Matrix	18
2.3 Polarimetric Radar Statistics	19

2.3.1	Random Walk Model of Scattering	20
2.3.2	Gaussian Model	23
2.3.3	Non-Gaussian Model	24
2.3.4	Speckle	25
2.3.5	Texture	26
2.3.6	Multilook Polarimetric Product Model	27
3	The Mellin Transform	29
3.1	An Historical Note	29
3.2	Mellin Kind Statistics	30
3.3	Univariate Mellin Transform	31
3.4	Matrix-Variate Mellin Transform	32
3.5	Fundamental Properties	33
3.6	Matrix-variate Mellin Kind Statistics	36
3.6.1	Mellin Kind Characteristic Function	36
3.6.2	Mellin Kind Matrix Moments	37
3.6.3	Mellin Kind Cumulant Generating Function	38
3.6.4	Mellin Kind Matrix Cumulants	38
3.7	Multilook Polarimetric Product Model	38
3.7.1	Application of the Matrix-Variate Mellin Convolution	38
3.7.2	Mellin Kind Statistics for the Multilook Polarimetric Product Model	39
4	Probability Distributions	41
4.1	Special Functions	41
4.2	Texture Distributions	43
4.2.1	Gamma Distribution	43
4.2.2	Inverse Gamma Distribution	44
4.2.3	Fisher-Snedecor Distribution	45
4.2.4	Beta Distribution	47
4.2.5	Inverse Beta Distribution	48
4.3	Speckle Distribution	49
4.3.1	Complex Wishart Distribution	49
4.3.2	Scaled Complex Wishart Distribution	50
4.4	Compound Matrix Distributions	50
4.4.1	Matrix-Variate \mathcal{K} Distribution	51
4.4.2	Matrix-Variate \mathcal{G}^0 Distribution	52
4.4.3	Matrix-Variate \mathcal{U} Distribution	53
4.4.4	Matrix-Variate \mathcal{W} Distribution	54
4.4.5	Matrix-Variate \mathcal{M} Distribution	55
4.5	The Matrix Log-Cumulant Diagram	56

5 Paper 1:	
Estimation of the Equivalent Number of Looks in Polarimetric Synthetic Aperture Radar Imagery	59
6 Paper 2:	
Application of the Matrix-Variate Mellin Transform to Analysis of Polarimetric Radar Images	83
7 Paper 3:	
Goodness-of-Fit Tests for Multilook Polarimetric Radar Data Based on the Mellin Transform	103
8 Conclusions	127
8.1 Concluding Remarks	127
8.2 Future Research	127
A A Relaxed Wishart Model for Polarimetric SAR Data	129
B Introduction to Second Kind Statistics: Application of Log-Moments and Log-Cumulants to Analysis of Radar Images	139
Bibliography	173

List of Tables

2.1	Microwave bands with letter designation, frequency range and wavelength range.	11
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List of Figures

1.1	Portraits of Mellin, Nicolas and Mathai	2
2.1	The imaging radar principle	10
2.2	The synthetic aperture radar principle	12
2.3	Electromagnetic wave with linear polarisation	13
2.4	The principle of multilook processing	17
2.5	Surface roughness and types of scattering	19
2.6	Random walk model of strong scattering	21
2.7	Random walk model of weak scattering	22
2.8	Example of radar speckle pattern	25
3.1	Relations in matrix-variate Mellin kind statistics	37
4.1	Gamma distributed texture	44
4.2	Inverse gamma distributed texture	45
4.3	Fisher-Snedecor distributed texture	46
4.4	Beta distributed texture	47
4.5	Inverse beta distributed texture	48
4.6	\mathcal{K} distributed multilook intensity	51
4.7	\mathcal{G}^0 -distributed multilook intensity	52
4.8	\mathcal{U} -distributed multilook intensity	54
4.9	The matrix log-cumulant diagram	56

Nomenclature

List of Notation

A	amplitude (of scattering coefficient)
\mathcal{A}	set of amplitude components
\mathbf{C}	polarimetric covariance matrix
\mathbb{C}	complex plane
d	number of polarimetric channels
e	Euler's number
E	electromagnetic field component
f	probability density function
\mathcal{M}	Fisher-Snedecor distribution
g	arbitrary scalar function defined on Ω_+
\mathcal{G}^0	G^0 distribution
h	arbitrary scalar function defined on Ω_+
h_{rms}	root mean square height of rough surface
I	intensity
I_L	multilook intensity
\mathbf{I}_d	identity matrix with dimension d
j	imaginary unit
k	wave number
\mathbf{k}	Pauli basis scattering vector
\mathcal{K}	K distribution
L	equivalent/nominal number of looks
\mathbf{M}	multichannel texture matrix
\mathcal{M}	M distribution
\mathcal{M}	Mellin transform
\mathcal{M}^{-1}	inverse Mellin transform
N	number of scatterers
P_{tot}	Total measured power
r	received electromagnetic wave/field
\mathbb{R}^+	line of positive real numbers
s	complex transform variable
S	scattering coefficient

\mathbf{s}	lexicographic basis scattering vector
\mathbf{S}	scattering matrix
t	transmitted electromagnetic wave/field
T	texture variable
\mathbf{U}	unitary transformation matrix
\mathcal{U}	U distribution
\mathbf{W}	complex Wishart distributed matrix
$\widetilde{\mathbf{W}}$	scaled complex Wishart distributed matrix
\mathcal{W}	W distribution
$\mathcal{W}_d^{\mathbf{c}}$	complex Wishart distribution
$s\mathcal{W}_d^{\mathbf{c}}$	scaled complex Wishart distribution
\mathbf{X}	complex matrix in Ω_+
\mathbf{Y}	complex matrix in Ω_+
z	complex scalar
α	shape parameter
β	beta distribution
β^{-1}	inverse beta distribution
γ	gamma distribution
γ^{-1}	inverse gamma distribution
θ	phase (of scattering coefficient)
$\boldsymbol{\theta}$	vector of texture parameters
Θ	set of phase components
κ_ν	ν th-order matrix-variate Mellin kind cumulant (matrix log-cumulant)
λ	shape parameter
μ	location parameter
μ_ν	ν th-order matrix-variate Mellin kind moment (matrix log-moment)
ν	order parameter
ρ	distance between radar and target
σ^2	mean radar cross section
Σ	scale matrix of matrix distributions
ϕ	matrix-variate Mellin kind characteristic function
φ	matrix-variate Mellin kind cumulant generating function
Ω_+	cone of positive definite Hermitian matrices

List of Acronyms

AD	Anderson-Darling
ALOS	Advanced Land Observing Satellite
ASAR	Advanced Synthetic Aperture Radar
CCRS	Canadian Centre for Remote Sensing
CCT	Centre Canadien de télédétection
CDF	cumulative distribution function
CF	characteristic function
CGF	cumulant generating function
COSMO-SkyMed	CONstellation of small Satellites for the Mediterranean basin Observation
CvM	Cramér-von Mises
CW	continuous wave
EM	electromagnetic
ENL	equivalent (or effective) number of looks
ERS	European Remote Sensing satellite
FM-CW	frequency modulated continuous wave
FT	Fourier transform
GoF	goodness-of-fit
GPR	ground penetrating radar
IEEE	Institute of Electrical & Electronics Engineers
IID	independent and identically distributed
JERS	Japanese Earth Resources Satellite
JPL	Jet Propulsion Laboratory
KS	Kolmogorov-Smirnov
MKS	Mellin kind statistics
ML	maximum likelihood
MLC	matrix log-cumulant
MLC	multilook complex
MLM	matrix log-moment
MoLC	method of log-cumulants
MoMLC	method of matrix log-cumulants
MT	Mellin transform
NASA	National Aeronautics and Space Administration
PALSAR	Phased Array type L-band Synthetic Aperture Radar
PDF	probability density function
RAR	real aperture radar
RV	random variable
SAR	synthetic aperture radar
SF-CW	step frequency continuous wave
SLC	single-look complex

Chapter 1

Introduction

The first part of this chapter motivates the application of the Mellin Transform to the analysis of polarimetric radar images. The second and third part provides an extended summary of the chapters of the thesis, including three journal publications. A list of other publications produced is included at the end.

1.1 Motivation

The univariate *Mellin transform* is an integral transform named after the Finnish mathematician *Robert Hjalmar Mellin* (1854-1933), which has found many applications in mathematics, statistics, physics and engineering. It was first applied to statistical models of radar images by *Jean-Marie Nicolas*, a French professor of signal and image processing. In a milestone paper [Nicolas, 2002], he presented a new framework for statistical analysis of distributions¹ of single polarisation amplitude and intensity images. The paper was followed by a comprehensive technical report [Nicolas, 2006], which laid the foundation for a new approach to deduction of sophisticated distribution models, including their functional characterisation, the expressions for their statistical moments, and estimators of the model parameters.

In the years following the seminal publications, a large number of papers have appeared that use the analysis framework of Nicolas to solve many different radar image analysis problems. The list covers applications such as statistical modelling [Moser et al., 2006a, Moser et al., 2006b, Bombrun and Beaulieu, 2008], speckle filtering [Nicolas, 2003, Achim et al., 2006, Chen and Liu, 2008], image classification [Tison et al., 2004], image segmentation [Benboudjema et al., 2007, Galland et al., 2009], change detection [Bujor et al., 2004, Moser and Serpico, 2006, Moser and Serpico, 2009], estimation of interferometric coherence [Abdelfattah and Nicolas, 2006] and image compression [Valade and Nicolas, 2004]. Still it seems like the awareness about the strength of the method is increasing rapidly.

¹The term *distribution* is used in this thesis as equivalent to probability density function.



Figure 1.1: Three innovators who have provided the theoretical underpinning for this thesis: Hjalmar Mellin, Jean-Marie Nicolas and Arak Mathai (left to right).

During the course of my research on radar image classification, practical requirements urged me to look into certain estimation problems for parameters of multilook polarimetric radar data distributions. The solutions I obtained bore strong similarities with those achieved by Mellin transform methods, but looked like matrix-variate extensions. This was the starting point of my quest for a matrix-variate generalisation of the Mellin transform, which proved successful. It was found in shape of the so-called *M-transform*, a lesser known integral transform for functions defined on positive definite Hermitian matrices. It was proposed by the Indian/Canadian statistician *Arakaparambil Mathai Mathai* in [Mathai, 1978] and extended to complex matrices in [Mathai, 1997]. Just like the univariate Mellin transform is tailor-made for distributions of real positive random variables, the matrix-variate M-transform went hand in glove with the covariance matrix distributions used to describe multilook polarimetric radar data.

The rediscovery of the M-transform opened the door for an extension of the *Mellin kind statistics*, which is the name used for the theoretical framework of Nicolas, to multilook polarimetric radar data. The prospect of achieving results with an impact comparable to those of Nicolas was sufficient to make me redefine my Ph.D. project completely. For one reason, the matrix-variate distributions describing polarimetric covariance matrices contain special functions and complicated forms that severely limit their mathematical tractability by conventional methods. Their applicability would certainly benefit from a method which provided a new perspective on functional characterisation, computation of statistical moments, model visualisation and parameter estimation. If the complexity of these distribution cannot be handled, the only rescue is to resort to the mathematically simpler but less accurate Wishart distribution model, which in many situations restricts the quality of the model based inference. Secondly, it was envisioned that the methods based on the Mellin transform would make possible the full use of the statistical information contained in the multilook polarimetric data, including the correlations between the polarimetric channels, and not only the single polarisation intensities. Furthermore, an extension of the Mellin kind statistics would provide a general

theory treating single polarisation images as a special case. The results presented in the current thesis proves in my opinion that the change of direction was a right decision.

It was advocated above that Mellin kind statistics have a positive effect on the applicability of certain distributions that are mathematically complex, but provide better fit with real data than the simpler alternative. The distributions referred to are those arising from the *doubly stochastic product model* for multilook polarimetric radar images. It is therefore relevant to comment on the necessity of such models. For low resolution radar images, each resolution cell contains a high number of microwave scatterers. The scattered electromagnetic field is the coherent sum of contributions from all these scatterers, and the central limit theorem asserts that it can be accurately modelled by Gaussian statistics. With the steadily improving spatial resolution of operational synthetic aperture radars, the Gaussian assumption is frequently challenged and often fails, in particular for scenes of urban environment, but also for natural surfaces such as forest and sea. Concerning the polarimetric aspect, it should be expected that the technological evolution will gradually replace single polarisation radars with instruments that have increasing capabilities for polarimetry. Hence, the need for adequate polarimetric distribution models will be more and more emergent.

After having argued in terms of technicalities, a more fundamental question naturally arises: Why is statistical modelling of radar data an important research task? When we zoom out and look at the benefit for society and mankind, the importance is connected to the value of the imaging radar as a remote sensing instrument. In the context of Earth observation, the radar is distinguished by its all-weather and all-season capabilities. It performs its measurements irrespective of cloud cover and sun conditions. Together with the wide spatial coverage and relatively good temporal resolution of the image acquisitions, these properties make spaceborne radar crucial for tasks such as monitoring of rain forest degradation and deforestation, change detection in Arctic glaciers, and mapping of sea ice conditions. A common aspect of these applications is that they require observations of inaccessible areas that are impossible to cover by on-site measurements. Radar remote sensing is both a practical and a cost effective alternative. The limiting factor of radar images is their content of strong speckle or clutter, an inherent feature of the coherent imaging process which complicates the interpretation and potential for information extraction. The most efficient remedy is to analyse the images within the context of a suitable statistical model, which closes the argument.

1.2 Chapter Review

Chapter 2 provides an introduction to radar imaging, which starts at the very fundamental by discussing key properties of active microwave sensors, describing different frequency bands and different types of imaging radars. The function of synthetic aperture radars and polarimetric radars is explained. We next look at the data formats delivered by polarimetric radars, starting with single-look complex data and moving on to multilook complex data, while explaining the concept of multilooking. The chapter

ends with the presentation of a physical model of radar speckle, and an explanation of its link to the traditional statistical models for the radar measurements. We define speckle and texture, the two factors of the doubly stochastic product model, which forms the basis of our statistical analysis.

Chapter 3 is opened by a historical review of the Mellin transform, with particular focus on its use in statistical distribution theory. The underlying ideas of Nicolas' univariate Mellin kind statistics are outlined, before we present Mathai's generalised Mellin transform. The original contribution of the author starts with the derivations of fundamental properties of the matrix-variate Mellin transform. It is followed by the new definitions of matrix-variate Mellin kind statistics. Finally, the Mellin kind statistics of a general multilook polarimetric product model are revealed, while leaving the detailed derivations to Paper 2.

Chapter 4 starts with the definitions of the special functions needed in the sequel. It further introduces candidate univariate distributions that can be used to model texture, together with their Mellin kind statistics, that were derived in [Nicolas, 2006]. The scaled complex Wishart distribution is also presented as a model for fully developed speckle, unmodulated by texture. The Mellin kind statistics given for this distribution are contributed in Paper 2. With the listed texture and speckle distributions as building block, we arrive at five compound distributions for the polarimetric covariance matrix. The Mellin kind statistics derived for all distributions are new.

Chapters 5–7 contain Papers 1–3. They are described separately in the next section.

Chapter 8 gives the conclusions. It lists the main results and points out future directions of research based on the work documented in the thesis.

Appendix A is a conference paper (referenced as Paper 4 in the list of Section 1.4) which contains ideas about statistical modelling of speckle filtered multilook polarimetric radar data. It is included for completeness and availability, since the results are mentioned and referenced in the thesis.

Appendix B is a translation of [Nicolas, 2002], which is a key reference for this thesis. The paper was originally published in French, but translated by the current author for the benefit of the research reported in this thesis. It is included here, with permission from the author, for the convenience of readers who are not proficient in French.

1.3 Publication Review

The following three publications are included in this thesis as Chapters 5, 6 and 7. A summary is given for each, highlighting the original contributions of the authors. The papers appear in chronological order and document the progress of my work in terms of maturity and depth of both exposition and content. Paper 1 is least mature in the sense that the theory of Mellin kind statistics for the polarimetric case had not been formulated yet. The connection of the results to the Mellin transform was discovered

later, as documented in Paper 2. Paper 3 is a first step in the direction of applications of the theory.

Paper 1

S.N. Anfinsen, A.P. Doulgeris and T. Eltoft, “**Estimation of the Equivalent Number of Looks in Polarimetric Synthetic Aperture Radar Imagery**”, IEEE Transactions on Geoscience and Remote Sensing, vol. 47, no. 11, pp. 3795–3809, December 2009.

The paper studies estimators for the equivalent number of looks (ENL), a parameter which is found in all distributions used to model multilook radar images. Still, the literature on the topic is very sparse. The relation $L = E\{I\}^2 / \text{Var}\{I\}$, where I denotes intensity, is often mistaken for being the definition of the ENL in the single polarisation case, assuming fully developed speckle and no texture. In reality, it is just one particular way of resolving the ENL from moments relations.

In the paper, we propose two new estimators for the ENL that are adapted to multilook polarimetric radar data. The expressions are derived by examining different kinds of moments of the polarimetric covariance matrix. The first estimator is a generalisation of the expression given above. The second is found from moments of the log-determinant of the covariance matrix, and is also found to be the maximum likelihood (ML) estimator based on the Wishart distribution model for multilook polarimetric radar data. The proposed estimators are the first ones to take the full covariance matrix as input, thereby utilising all the available statistical information. This is reflected in the experimental results in terms of superior statistical properties. The ML estimator has the lowest bias and variance, and also most robust with respect to the assumption of no texture.

An approach to unsupervised estimation of the ENL is also presented, where the ML estimator is used to compute small sample estimates over the whole image, regardless of the homogeneity in the estimation window. It is shown that a robust estimate of the ENL can be extracted from a probability density function estimate of the collection of small sample estimates. This is only possible when using the novel low variance ML estimator.

There are minor differences between the published version of the paper and the version included in the thesis. Equation (23) has been corrected in the thesis version. The journal version gives an expression which is valid for real matrices, while complex matrices are considered. The paper has been reformatted to a different font size. Otherwise, the differences are mainly orthographical, due to different preferences of English style. Some symbols have also been changed to harmonise the notation of this paper with the others.

Paper 2

S.N. Anfinsen and T. Eltoft, "**Application of the Matrix-Variate Mellin Transform to Analysis of Polarimetric Radar Images**", submitted to IEEE Transactions on Geoscience and Remote Sensing.

This paper introduces the Mellin kind statistics framework for analysis of multilook polarimetric radar images. It builds on the equivalent framework derived by Nicolas for the single polarisation case, and the extension to multilook polarimetry rests upon Mathai's matrix-variate Mellin transform. The combination of these ideas, leading to a whole new set of definitions for the polarimetric case, is an original contribution.

The paper gives a thorough review of the univariate Mellin kind statistics. It explains Nicolas' idea of introducing a new kind of characteristic function for real positive random variables by replacing the Fourier transform with the Mellin transform in the definition of the conventional characteristic function. The moments and cumulants retrieved from this characteristic function are calculated on logarithmic scale, and are therefore called log-moments and log-cumulants. We also emphasise analogies between Mellin kind statistics and conventional (Fourier kind) statistics, showing that the Mellin kind statistics are the natural tools for analysis of a multiplicative signal model, just like the conventional statistics are for the familiar additive model. This care is taken to make the presentation pedagogical, but also to make Nicolas' theory available to a wider audience, since his most comprehensive derivations are only published in French [Nicolas, 2002, Nicolas, 2006].

The novel contributions include definitions of the Mellin kind characteristic function, cumulant generating function, moments and cumulants for the matrix-variate case describing multilook polarimetric radar images. We have further defined a matrix-variate Mellin convolution and correlation, and proved corresponding convolution and correlation theorems. The convolution theorem shows that the Mellin transform of the Mellin convolution of two functions decomposes as the product of Mellin transform of the individual functions. This result is needed to formulate the Mellin kind statistics under the multilook polarimetric product model, which is probably the contribution with the largest practical significance.

As an example of applications of the new theory, we have derived parameter estimators for some product model distributions for the polarimetric covariance matrix. The experimental results show that the estimators based on Mellin kind statistics are superior to all alternative estimators from the literature. The proposed estimation procedure is interpreted visually in terms of a diagram where we plot empirical matrix log-cumulants computed from data samples together with the population matrix log-cumulants of the distributions. The matrix log-cumulant diagram is a matrix-variate extension of the diagram Nicolas has used in the univariate case.

Paper 3

S.N. Anfinsen, A.P. Doulgeris and T. Eltoft, "**Goodness-of-Fit Tests for Multilook Polarimetric Radar Data Based on the Mellin Transform**", submitted to IEEE Transactions on Geoscience and Remote Sensing.

In this paper, the Mellin kind statistics framework is used to derive goodness-of-fit tests for distributions of the polarimetric covariance matrix derived under the multilook polarimetric product model. These are, to the best of our knowledge, the first formal statistical tests that have been devised for these complicated distributions. The test statistic is constructed from the matrix log-cumulants defined in Paper 2. In order to deduce sampling distributions for the test statistics, asymptotic statistics of the matrix log-cumulants have been derived. They proposed test statistic can be applied to both simple and composite hypothesis tests.

For the simple hypothesis, the sampling distribution of the test statistic is asymptotically χ^2 distributed. We demonstrate that this is a good approximation even for moderate sample sizes. For the composite test, we must resort to Monte Carlo simulations to find the sampling distribution. This approach has a higher computational cost, but produces the true sampling distribution regardless of sample size.

The simple and composite tests have been tested on simulated and real data. Assessments of the test powers show that we have found a useful method which meets the need for formal procedures of testing model fit for compound covariance matrix distributions. The matrix log-cumulant diagram introduced in Paper 2 is further promoted as an intuitive visualisation tool for interpretation of the test procedure. As a graphical aid for informal model selection and validation, it separates very well between different distributions whose differences are mainly manifested in the heavy tails.

The paper emphasises the coupling between the problems of goodness-of-fit testing and parameter estimation. A new estimation technique for parameters of the texture distributions is motivated by the proposed tests. The estimator maximises the asymptotic likelihood of the compounded matrix distribution, and is effectively a method of moment type procedure using multiple matrix log-cumulants. The estimator is tested in the experiments of Paper 2, and exhibits superior performance in terms of both bias and variance. It is also an intrinsic part of the composite test, which requires estimation of unknown distribution parameters.

1.4 Other Publications and Presentations

As first author:

1. S. N. Anfinsen, R. Jenssen and T. Eltoft, "Clustering of polarimetric SAR data with an information theoretic kernel method," presented at the *IEEE Int. Geosci. Remote Sens. Symp. (IGARSS '06)*, Denver, U.S., 31 Jul.-4 Aug. 2006, not published.

2. S. N. Anfinsen, R. Jenssen and T. Eltoft, "Spectral clustering of polarimetric SAR data with Wishart-derived distance measures", *Proc. 3rd Int. Workshop on Science and Applications of SAR Polarimetry and Polarimetric Interferometry (POLinSAR '07)*, Frascati, Italy, 22-26 Jan. 2007, ser. ESA SP-644, Mar. 2007, 8 pp.
3. S. N. Anfinsen, A. P. Doulgeris and T. Eltoft, "Estimation of the Equivalent Number of Looks in Polarimetric SAR Imagery", *Proc. IEEE Int. Geosci. Remote Sens. Symp. (IGARSS '08)*, vol. 4, Boston, U.S., 6-11 Jul. 2008, pp. 487-490.
4. S. N. Anfinsen, T. Eltoft and A. P. Doulgeris, "A relaxed Wishart model for polarimetric SAR data (POLinSAR '09)", *Proc. 4th Int. Workshop on Science and Applications of SAR Polarimetry and Polarimetric Interferometry*, Frascati, Italy, 26-30 Jan. 2009, ser. ESA SP-668, Apr. 2009, 8 pp.
5. S. N. Anfinsen and T. Eltoft, "Moment-based goodness-of-fit tests for polarimetric radar data", presented at the *Proc. IEEE Int. Geosci. Remote Sens. Symp. (IGARSS '09)*, Cape Town, South Africa, 12-17 Jul. 2009, not published.
6. S. N. Anfinsen and T. Eltoft, "Analysis of multilook polarimetric data with the matrix-variate Mellin transform", *Proc. 8th Eur. Conf. Synthetic Aperture Radar (EUSAR 2010)*, Aachen, Germany, 7-10 Jun. 2010, in press.

As coauthor:

6. A. P. Doulgeris, S. N. Anfinsen and T. Eltoft, "Analysis of non-Gaussian PolSAR data", *Proc. IEEE Int. Geosci. Remote Sens. Symp. (IGARSS '08)*, Barcelona, Spain, 23-27 Jul. 2008, pp. 160-163.
7. A. P. Doulgeris, S. N. Anfinsen and T. Eltoft, "Classification with a non-Gaussian model for PolSAR data", *IEEE Trans. Geosci. Remote Sens.*, vol. 46, no. 10, pp. 2999-3009, Oct. 2008.
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Chapter 2

Radar Polarimetry

This chapter gives an overview of the fundamental properties of radar imaging, with emphasis on the potential of spaceborne polarimetric synthetic aperture radar. The foundation is laid for the subsequent theoretical developments by the definition of the multilook polarimetric product model. The connections to an underlying physical model of microwave scattering are explained.

2.1 Imaging Radar

An *imaging radar*¹ illuminates the target scene with directional pulses of electromagnetic (EM) energy, measures the backscattered energy and the round-trip time, and uses this information to form an image. The radar operates in the microwave region of the EM spectrum, and the backscatter depends on dielectric and geometrical properties (roughness and shape) of the target. If both amplitude and phase of the backscattered wave is measured, the image will constitute a two-dimensional map of the EM *scattering coefficient*, which is defined in Section 2.2.2. The scattering coefficient can be measured with different combinations of polarisations at the transmitter and receiver. The use of multiple polarisations gives rise to multidimensional image data known as polarimetric radar images. The principles of radar imaging are shown in Figure 2.1.

2.1.1 Frequency Bands

Radars are classified by the frequency band their emitted pulse belongs to. The frequency bands most commonly used by spaceborne earth observation radars are L-band (Seasat, JERS-1 and ALOS PALSAR), C-band (ERS-1, ERS-2, Envisat ASAR, Radarsat-1 and Radarsat-2) and X-band (TerraSAR-X and COSMO-SkyMed). P-band² radar has

¹Radar was originally an acronym for the method of *radio detection and ranging*, but has been assimilated as a standard word in most languages, and now refers to the instrument itself.

²The P-band is not defined in the IEEE standard reported in Table 3, but refers to a band which lies partially in the UHF-band and partially in the VHF-band, according to this designation.

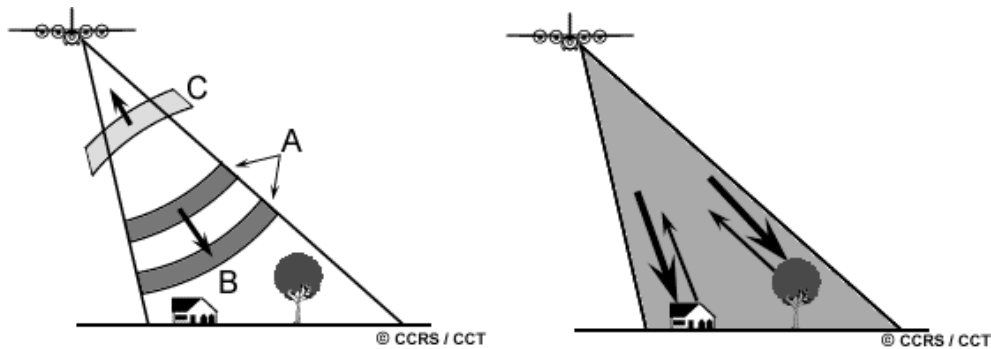


Figure 2.1: An imaging radar is an active microwave instrument which transmits electromagnetic pulses (A), focused by the antenna into a beam (B), and receives a portion of the reflected energy (C) backscattered from various objects.

only been mounted on airborne research missions, but a satellite mission named BIOMASS is planned. The letter designation applied to different frequency ranges is shown in Table 3, whose source is the IEEE Standard 521-2000(R2009) [IEEE, 2009].

The penetration depth of the microwave into a target medium, such as vegetation, glacier, sea ice and soil, increases with wavelength. It also depends on the moisture level of the medium. Thus, the usefulness of the different frequency bands vary with application. For instance, C-band is generally preferred for mapping of sea ice, because it provides the best contrast between sea and ice (although the contrast also depends largely on polarisation and incidence angle). L-band is preferred for studies of many types of vegetation, for which the wave penetrates the vegetation canopy and reaches the ground or surface level. P-band is required to penetrate rain forest, and thus to obtain meaningful estimates of biomass for this biotope, since the measurements saturate at shorter wavelengths, resulting in underestimation. The frequency band also determines the scale of roughness which interacts with the radar wave and influences the measurements. The backscattered energy is sensitive to surface curvature and roughness at length scales near the radar wavelength.

2.1.2 Properties of Imaging Radar

The major advantage of active microwave instruments is that they work independently of sunlight conditions and cloud cover. Unlike optical sensors, they operate equally well nighttime as daytime, and the attenuation of the signal by clouds and water vapour is negligible in most bands. L-band radars may experience disturbance by ionospheric Faraday rotation under certain conditions [Freeman and Saatchi, 2004], while X-band and Ku-band backscatter is sensitive to precipitation, which has been successfully retrieved from spaceborne SAR observations [Marzano and Weinman, 2008].

Radar systems give access to different parameters compared to optical systems. Their

³Frequencies from 216-450 MHz are sometimes called P-band.

Table 2.1: Microwave bands with letter designation, frequency range and wavelength range.

Band	Frequency range	wavelength range [cm]
HF	3-30 MHz	10-100 m
VHF ³	30-300 MHz	1-10 m
UHF ³	300-1000 MHz	30-100 cm
L	1-2 GHz	15-30 cm
S	2-4 GHz	7.5-15 cm
C	4-8 GHz	3.75-7.5 cm
X	8-12 GHz	2.5-3.75 cm
Ku	12-18 GHz	16.7-25 mm
K	18-27 GHz	11.1-16.7 mm
Ka	27-40 GHz	7.5-11.1 mm
V	40-75 GHz	4.0-7.5 mm
W	75-110 GHz	2.7-4.0 mm
mm	110-170 GHz	1.8-2.7 mm

measurements can be related to surface roughness, humidity and geometrical properties, as previously discussed. They cannot be used to retrieve biophysical parameters that require access to radiances, reflectances and brightness temperatures. Nevertheless, many mapping and classification products can be obtained with similar quality as for optical instruments, and the independence of solar illumination and cloud cover assures temporal consistence of the service. This property is especially attractive for monitoring of high latitude areas, where Arctic/Antarctic winter and frequent cloud cover limits the usefulness of optical sensors. Persistent cloud cover also severely restrict the capacity of optical monitoring of tropical rain forest, whereas radar instruments deliver consistent results, which is a vital requirement for operational services.

2.1.3 Instruments

Different types of radar can be classified as imaging radars. A ground-penetrating radar (GPR) is an active microwave instrument used to image the subsurface, which is often implemented as a continuous wave (CW) radar⁴. It is applied to a variety of media, including rock, soil, ice, snow, fresh water and man-made structures. The data are processed into a radargram which shows the depth profile and indicates boundaries between layers with different dielectric constant. A weather radar is normally a ground-based Doppler radar capable of locating precipitation, calculating its intensity

⁴A continuous wave radar transmits continuous waves instead of pulses. Range measurements are enabled by use of e.g. a frequency modulated (FM-CW) or step-frequency (SF-CW) transmitted wave.

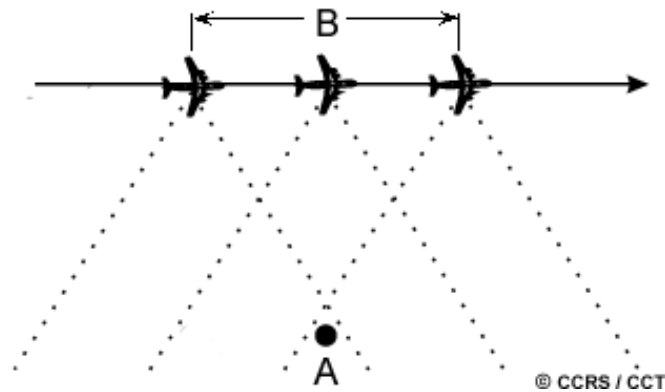


Figure 2.2: The SAR principle: A target (A) is illuminated by several pulses of the radar beam. The backscattered echoes of each pulse is recorded. The length of the synthesized antenna (B) is the the distance between the points where the target enters and leaves the radar beam.

and velocity, and identifying its type. It scans a volume of air around the radar station, and images are produced as different cross-sections of the scanned volume. In addition, we have real aperture radar (RAR) and *synthetic aperture radar* (SAR), whose function is explained in the next section. These can both be ground-based or airborne. SAR data are also widely available from spaceborne instruments that provide regular global coverage. These instruments are therefore important and reliable sources of information for various monitoring programs.

The methods presented in this thesis are aimed at multilook polarimetric data, a data format which is defined in Section 2.2.4. The most obvious subject for the proposed analysis tools is polarimetric SAR data. However, polarimetric GPRs and weather radars exist [Langley et al., 2009, Galetti et al., 2008], and their data can also be processed into the same format, even though other representations are more common. The developments of this thesis are presented as a generic theory for polarimetric radar data, since the theoretical framework can in principle be applied to any kind of multilook polarimetric radar data.

2.1.4 Synthetic Aperture Radar

The synthetic aperture processing technique [Oliver and Quegan, 2004, Cumming and Wong, 2005, Massonnet and Souyris, 2008] must be credited for the availability of high resolution radar images captured by spaceborne platforms. The difference between a RAR and a SAR sensor lies in the image resolution in the azimuth, or along-track direction. The azimuth resolution of a RAR is determined by the beamwidth of the antenna, which is limited by practical constraints. A SAR obtains multiple measurements of the scene at different azimuth angles, and uses advanced signal processing to combine these into an image with improved azimuth resolution. The effect is the same as using

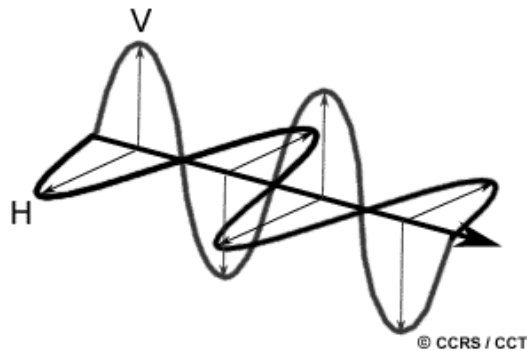


Figure 2.3: Electromagnetic wave with linear polarisation.

an aperture whose size is larger than the actual, hence the term *synthetic aperture*. The finest azimuth resolution achieved by the most recent spaceborne SAR instruments is 1 meter for TerraSAR-X and COSMO-SkyMED and 3 metres for Radarsat-2 and ALOS PALSAR. The SAR principle is explained by Figure 2.2.

2.1.5 Polarimetric Radar

An EM wave consists of electric and magnetic field components that oscillate in phase perpendicular to each other and perpendicular to the direction of energy propagation (see Figure 2.3). The *polarisation* of an EM wave describes the orientation of its oscillations. A fully polarimetric radar simultaneously transmits microwave pulses with two different orthogonal polarisations, it measures the electric field components at two orthogonal polarisations, and resolves the scattering coefficients for all four combinations of transmit and receive polarisation. Any polarisation can be synthesised as a superposition of two orthogonal polarisation. The fully polarimetric measurements thus constitute a complete description of the scattering characteristics of the resolution cell, which can be analysed for an arbitrary polarisation by a simple transformation of the orthogonal basis.

SAR instruments commonly use linear polarisations, where the electrical field is oriented in a single direction, normally horizontally and vertically. Another option is to use circular or elliptical polarisations, where the electric field rotates rightwards or leftwards in the direction of propagation. The polarisations at the transmitter and receiver need not be the same, and can be chosen to optimise system performance or in accordance with given restrictions or requirements.

In *radar polarimetry*, we use the complex scattering coefficients measured at different polarisations to characterise the target. A polarimetric radar extends the capabilities of normal single polarisation radar, not only because it provides multichannel data. Most importantly, it provides a strong link to the physics of the scattering process and allows us to identify distinct scattering mechanisms. Coherent scattering from point sources can be resolved as scattering from elementary geometrical objects, such

as spheres, cylinders, dipoles, diplanes, dihedrals and trihedrals. Incoherent scattering can be resolved as surface (single bounce) scattering, double bounce scattering and volume scattering. The methods used are known as *polarimetric decompositions*, and a rich literature has emerged on the topic [Mott, 2007, Lee and Pottier, 2009, Cloude, 2010].

2.2 Data Formats

This section describes the data formats encountered for polarimetric radar data. We start from a mathematical description of the most elementary measurable in polarimetric radar imaging, the matrix holding the scattering coefficients of all polarimetric channels. We then describe the multilooking process and transformation of the data into the intensity domain.

2.2.1 Single Look Complex Data

As described in the previous section, the fully polarimetric SAR instrument separately transmits orthogonally polarised microwaves pulses, and measures orthogonal components of the received signal. For each pixel, the measurements result in a matrix of scattering coefficients. These are complex-valued, dimensionless numbers that describe the transformation of the transmitted (incoming) EM field to the received (backscattered) EM field for all combinations of transmit and receive polarisation.

The transformation can be expressed as

$$\begin{bmatrix} E_x^r \\ E_y^r \end{bmatrix} = \frac{e^{jk\rho}}{\rho} \begin{bmatrix} S_{xx} & S_{xy} \\ S_{yx} & S_{yy} \end{bmatrix} \begin{bmatrix} E_x^t \\ E_y^t \end{bmatrix} \quad (2.1)$$

where $j = \sqrt{-1}$ is the imaginary unit, k denotes wavenumber and ρ is the distance between radar and target. The subscript of the EM field component E_i^j , where $i \in \{x, y\}$, refers to the polarisation it is associated with. The superscript of E_i^j , where $j \in \{r, t\}$, indicates if it is the transmitted or received field component. The orthogonal polarisations are denoted x and y for generality, although it has been assumed that the same polarisations are used at the transmitter and the receiver. The scattering coefficients S_{ij} , $i, j \in \{x, y\}$, are subscripted with the associated receive and transmit polarisation, in that order.

The scattering coefficients are complex-valued, and they comprise what is known in radar imaging terminology as a *look*. The SAR processor may split the full synthetic aperture into several subapertures. This is done by splitting the Doppler bandwidth into sub-bands and extracting the band-limited signal. The portion of the SAR signal associated with one subaperture or subband represents an individual look of the scene. The scattering coefficients represent one such look after the image has been focused. Data structures that contain the scattering coefficients are therefore referred to as *single-look complex* (SLC) data. The single-look term is used also for sensors without the synthetic aperture capacity, even though the look extraction aspect is missing.

2.2.2 Scattering Matrix

The choice of polarisations is from now on restricted to the linear pair, that is, the horizontal or vertical polarisation. This is the most commonly used set of orthogonal polarisations for SAR systems, and the restriction can be done without loss of generality for the methods subsequently derived. The exception is for interpretations of scattering mechanisms, that rely explicitly on the choice of polarisation basis.

We extract from (2.1) the *scattering matrix*

$$\mathbf{S} = \begin{bmatrix} S_{hh} & S_{hv} \\ S_{vh} & S_{vv} \end{bmatrix} \in \mathbb{C}^{2 \times 2}, \quad (2.2)$$

where the subscripts h and v denote horizontal and vertical polarisation, respectively. This matrix is also known as the *Sinclair matrix*, and holds all scattering coefficients measured by the *fully polarimetric radar*.

2.2.3 Scattering Vector

Lexicographic Basis

The scattering vector is simply the vectorised version of the scattering matrix, defined as

$$\mathbf{s} = \text{vec}(\mathbf{S}^T) = \begin{bmatrix} S_{hh} \\ S_{hv} \\ S_{vh} \\ S_{vv} \end{bmatrix} \in \mathbb{C}^{4 \times 1}, \quad (2.3)$$

where $\text{vec}(\cdot)$ is the column stacking vectorisation operator. The vector elements can also be seen as coefficients of the lexicographic decomposition of the scattering matrix:

$$\mathbf{S} = S_{hh} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + S_{hv} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} + S_{vh} \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} + S_{vv} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}. \quad (2.4)$$

This vector is therefore known as the *lexicographic basis scattering vector*.

Pauli Basis

Another representation of the scattering vector is obtained by a linear transformation of the lexicographic basis vector. The *Pauli basis scattering vector* is obtained as

$$\mathbf{k} = \mathbf{U} \mathbf{s} = \frac{1}{\sqrt{2}} \begin{bmatrix} S_{hh} + S_{vv} \\ S_{hh} - S_{vv} \\ S_{hv} + S_{vh} \\ j(S_{hv} - S_{vh}) \end{bmatrix} \in \mathbb{C}^{4 \times 1}, \quad (2.5)$$

where \mathbf{U} is the unitary transformation matrix

$$\mathbf{U} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & 1 & 0 \\ 0 & j & -j & 0 \end{bmatrix}. \quad (2.6)$$

The vector elements are the coefficients in the Pauli decomposition of the scattering matrix, given by

$$\begin{aligned} \mathbf{S} = & \frac{S_{hh} + S_{vv}}{\sqrt{2}} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \frac{S_{hh} - S_{vv}}{\sqrt{2}} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \\ & + \frac{S_{hv} + S_{vh}}{\sqrt{2}} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} + \frac{j(S_{hv} - S_{vh})}{\sqrt{2}} \begin{bmatrix} 0 & -j \\ j & 0 \end{bmatrix}. \end{aligned} \quad (2.7)$$

The basis of this decomposition contains the three 2×2 Pauli matrices, that were originally introduced in quantum mechanics to describe the spin of a spin $1/2$ particle in three spatial directions. The fourth basis matrix is the 2×2 identity matrix, which is associated with the first element of \mathbf{k} .

The advantage of the Pauli basis scattering vector is that it provides physical interpretations of its elements in terms of elementary scattering mechanisms. The respective elements of \mathbf{k} , denoted $\{k_1, \dots, k_4\}$, can be related to: single or odd-bounce scattering from a plane surface (k_1), diplane scattering (double-bounce or even-bounce) from corners with a relative orientation of 0° (k_2) and 45° (k_3), and the residue of antisymmetric components (k_4) [Lee and Pottier, 2009, Cloude, 2010].

Reduced Dimension Scattering Vectors

Let d be the dimension of \mathbf{s} , which is equivalent to the number of polarimetric channels. It will be referred to as the polarimetric dimension. The polarimetric dimension can be reduced compared to the quadrature polarimetric case of $d = 4$. If only one polarisation is used at the transmitter (or receiver), then only a subset of the scattering coefficients can be measured, and we obtain *dual polarisation* data with $d = 2$. Single polarisation data are obviously also a special case, with $d = 1$.

The case of $d = 3$ is encountered when we assume *reciprocity* of the target, defined as follows: For natural terrain it can safely be assumed that the cross-polarised channels are approximately equal: $S_{hv} \simeq S_{vh}$. These measurements can then be averaged to reduce uncertainty, and we obtain the reduced scattering vector

$$\mathbf{s} = \begin{bmatrix} S_{hh} \\ (S_{hv} + S_{vh})/\sqrt{2} \\ S_{vv} \end{bmatrix} \in \mathbf{C}^{3 \times 1}. \quad (2.8)$$

The term $\sqrt{2}$ ensures that the total measured power:

$$P_{\text{tot}} = |S_{hh}|^2 + |S_{hv}|^2 + |S_{vh}|^2 + |S_{vv}|^2$$

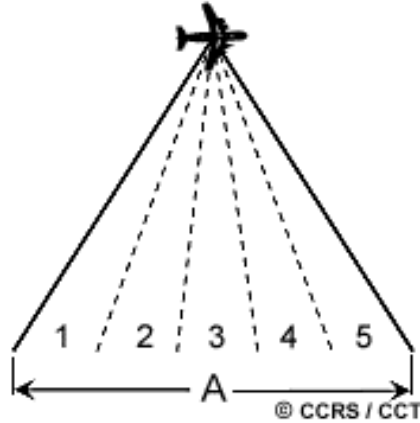


Figure 2.4: In multilook processing, the radar beam (A) is divided into several sub-beams, each providing an independent (in an ideal case) look at the illuminated scene.

is maintained regardless of a change of basis. The reciprocity assumption breaks down in urban environments, or generally when the target geometry is such that the amount of cross-polarised power depends on the radar look angle, which is typical for man-made targets and structures with non-random orientation.

The equivalent version of the Pauli basis scattering vector is

$$\mathbf{k} = \frac{1}{\sqrt{2}} \begin{bmatrix} S_{hh} + S_{vv} \\ S_{hh} - S_{vv} \\ S_{hv} + S_{vh} \end{bmatrix} \in \mathbb{C}^{3 \times 1}. \quad (2.9)$$

2.2.4 Multilook Complex Data

The *look* term was defined in Section 2.2.1 as a portion of the SAR signal recorded by a part of the synthetic aperture, known as a subaperture. The subaperture signal can be extracted from the total signal by filtering in the frequency domain. Multiple looks can be summed incoherently to produce a multilook image, an averaging operation known as *multilooking* [Cumming and Wong, 2005, Massonnet and Souyris, 2008]. Multilooking can also be done in the spatial domain, after the image has been focused. This is done by computing the mean value of a group of adjacent pixels, and must also be done incoherently, since the single-look data are complex-valued. The multilook principle is illustrated by Figure 2.4.

Multilooking of single polarisation radar data produces amplitude or intensity data that are real-valued. All phase information is discarded. In the polarimetric case, the multilooking process creates complex data which preserves information on the mean phase difference between the polarimetric channels. The format is known as *multilook complex* (MLC) data.

We remark that the averaged looks are correlated. For multilooking in the frequency domain, the correlation occurs because the filters used to split the Doppler bandwidth

into subbands have a slight overlap. In spatial domain multilooking, we average neighbour pixels that are correlated because they share a certain amount of information from the focusing process, due to the radar point spread function [Rignot and Chellappa, 1993]. This has an impact on the statistical modelling of MLC data. The exact form of the PDF for correlated data does not have a simple closed form expression [Goodman, 1975, Rignot and Chellappa, 1993, Gierull and Sikaneta, 2002]. The practical approach has been to derive the distribution as if the multilook samples were statistically independent, and replace the *nominal number of looks*, equivalent to the actual number of samples, with an *equivalent number of looks* (ENL). The ENL is must be estimated from the data, which is done by equating certain empirical sample moments with the corresponding theoretical population moments under the assumed statistical distribution model. This is the topic of Paper 1 in Chapter 5.

2.2.5 Covariance Matrix

Assume that L looks are available, in the form of the scattering vector sample $\{\mathbf{s}_\ell\}_{\ell=1}^L$. We refer to L as the nominal number of looks. The lexicographic basis scattering vector is multilooked by computing its sample covariance matrix, under the assumption that the $\{\mathbf{s}_\ell\}$ are zero mean, a condition discussed in Section 2.3.1. The multilooking operation is formulated as

$$\mathbf{C} = \frac{1}{L} \sum_{\ell=1}^L \mathbf{s}_\ell \mathbf{s}_\ell^H = \begin{bmatrix} \langle S_{hh} S_{hh}^* \rangle & \langle S_{hh} S_{hv}^* \rangle & \langle S_{hh} S_{vh}^* \rangle & \langle S_{hh} S_{vv}^* \rangle \\ \langle S_{hv} S_{hh}^* \rangle & \langle S_{hv} S_{hv}^* \rangle & \langle S_{hv} S_{vh}^* \rangle & \langle S_{hv} S_{vv}^* \rangle \\ \langle S_{vh} S_{hh}^* \rangle & \langle S_{vh} S_{hv}^* \rangle & \langle S_{vh} S_{vh}^* \rangle & \langle S_{vh} S_{vv}^* \rangle \\ \langle S_{vv} S_{hh}^* \rangle & \langle S_{vv} S_{hv}^* \rangle & \langle S_{vv} S_{vh}^* \rangle & \langle S_{vv} S_{vv}^* \rangle \end{bmatrix}, \quad (2.10)$$

where $(\cdot)^*$ and $(\cdot)^H$ are the complex conjugation and Hermitian transposition operators, respectively, and $\langle \cdot \rangle$ denotes a sample mean over all single-look measurements. This produces the *polarimetric covariance matrix* \mathbf{C} , which is positive definite and Hermitian symmetric. It has the real-valued intensities of the polarimetric channels on the main diagonal, and their complex covariances off the diagonal. The matrix \mathbf{C} is the measurable when working with multilook polarimetric radar data, and it is the subject of the subsequent analysis.

2.2.6 Coherency Matrix

The covariance matrix \mathbf{C} can be replaced by the *polarimetric coherency matrix* $\mathbf{\Omega}$, which is the term used for the sample covariance matrix of the Pauli basis scattering vector. It is computed from the sample $\{\mathbf{k}_\ell\}_{\ell=1}^L$ by

$$\mathbf{\Omega} = \frac{1}{L} \sum_{\ell=1}^L \mathbf{k}_\ell \mathbf{k}_\ell^H = \mathbf{UCU}. \quad (2.11)$$

The coherency matrix is often the preferred representation, because of its the physical interpretation. In the quadrature polarimetric case, the first three elements on the diagonal are the intensities:

$$[\Omega]_{11} = \langle |S_{hh} + S_{vv}|^2 \rangle, \quad (2.12)$$

$$[\Omega]_{22} = \langle |S_{hh} - S_{vv}|^2 \rangle, \quad (2.13)$$

$$[\Omega]_{33} = \langle |S_{hv} + S_{vh}|^2 \rangle, \quad (2.14)$$

that can be interpreted in terms of the same elementary scattering mechanisms as the elements of \mathbf{k} . Incoherent polarimetric decompositions, such as the Freeman decomposition [Freeman and Durden, 1998] and Yamaguchi decomposition [Yamaguchi et al., 2005], go even further in extracting entities that can be directly related to surface scattering, double-bounce scattering and volume scattering and their relative proportion.

For the methods developed in this thesis, it is of no concern whether we use the covariance matrix or the coherency matrix. The results are equally valid and useful for both data formats.

2.3 Polarimetric Radar Statistics

In this section we review statistical models for the polarimetric radar measurements. We specifically look at probability density functions (PDFs) for the different data formats based on a physical description of the scattering process. The EM field measured at the radar is a superposition in the far field of coherent microwave components, each produced by a reflection from a unique surface element called a *scatterer*. If the resolution cell contains only a limited number of scatterers, whose position could be accurately determined, then a deterministic description is in theory possible. In practice, we must

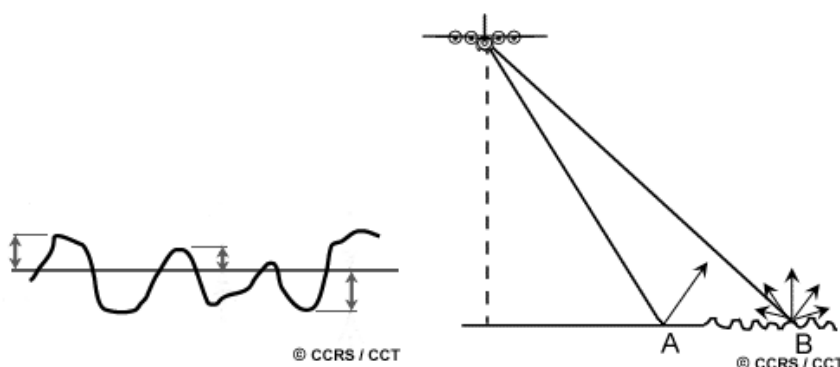


Figure 2.5: The roughness of the surface (left) determines the type of scattering produced (right). A smooth surface (A) causes specular reflection, while a rough surface scatters diffusely (B). The scattering can also be characterised as coherent (A) and incoherent (B).

resort to a stochastic model to describe the EM field.

A *rough surface* is a surface with height variation on a scale larger than the wavelength used to illuminate it [Delignon et al., 2001, Lee, 2005]⁵. In a radar remote sensing context, many natural surfaces can be considered as rough. The scatterers vary in number and geometry according to the surface type and the spatial resolution of the observation, but the number is generally assumed to be high. The nature of the scattering depends on the relative strength of the reflections from the scatterers and the distribution of the phase shift they induce. The next sections establish a mathematical description of the scattering process and definitions that characterise the scattering as *strong* or *weak*. Surface roughness and different types of scattering is shown in Figure 2.5.

At this point, it is pertinent to define some other terms. In physics, *coherence* is a property of two or more waves that are in phase both temporally and spatially. More generally, it describes the correlation between all physical quantities of the wave [Born and Wolf, 1999, Glickman, 2000]. This is the property which enables stationary interference. *Incoherent scattering* is defined as the scattering produced when an incident wave encounters scatterers that cause the scattered EM field to exhibit random variations in phase and amplitude due to lack of coherence. *Coherent scattering*, on the other hand, produces a deterministic scattered EM field. This happens when the incident wave is scattered by a fixed point target or a distributed target with scatterers whose relative position is fixed [Glickman, 2000].

2.3.1 Random Walk Model of Scattering

The scattering process is often described by a random walk model [Goodman, 2007, Lopès et al., 2008]. Let the EM field measured at the sensor be a sum of the field components reflected by N scatterers in the resolution cell. The scattering coefficient S of a general polarimetric channel thus represents the total scattering experienced by the microwave after interaction with N scatterers. This can be written as the sum

$$S = \sum_{k=1}^N S^{(k)} = \sum_{k=1}^N A^{(k)} e^{j\theta^{(k)}}. \quad (2.15)$$

Here $S^{(k)}$ is the scattering coefficient associated with the k th scatterer, whose polar decomposition yields the amplitude component $A^{(k)} = |S^{(k)}|$ and phase component $\theta^{(k)} = \angle S^{(k)}$. These represent the attenuation and phase shift imposed on the incident wave by the k th scatterer. We initially make three fundamental assumptions:

1. The amplitudes $\{A^{(k)}\}_{k=1}^N$ are independent and identically distributed (IID) random variables. So are the phases $\{\theta^{(k)}\}_{k=1}^N$.

⁵Surface roughness can be measured by the product of the wavenumber k and the root mean square height h_{rms} of the roughness. Moderate roughness starts at $kh_{\text{rms}} > 1$ and high roughness at $kh_{\text{rms}} > 5$.

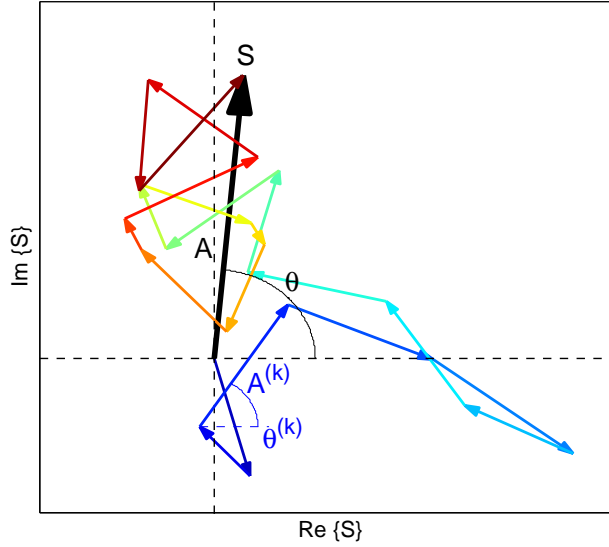


Figure 2.6: Random walk model of strong scattering. The scattering coefficient S with amplitude A and phase θ is shown as a phasor in the complex plane (black arrow). It is decomposed as the coherent sum of the scattering coefficients $S^{(k)}$ corresponding to N individual scatterers with amplitude $A^{(k)}$ and phase $\theta^{(k)}$. The accumulated phasor sum is indicated by the coloured arrows.

2. The $\{A^{(k)}\}$ and the $\{\theta^{(k)}\}$ are independent.
3. The $\{\theta^{(k)}\}$ are uniformly distributed over all angles.

Figure 2.6 illustrates equation (2.15) under the three assumptions listed above. It pictures the coherent summation as a two-dimensional random walk in the complex plane, with the in-phase component along the first axis and the quadrature phase along the second. We shall use the figure as a starting point for a discussion of different scattering regimes.

Strong Scattering

Figure 2.6 shows an example of a random walk with $N = 20$ steps, each representing the reflection by an individual scatterer. The coloured arrows show the accumulated coherent sum, which ends up as the total scattering coefficient S , shown as the thick black arrow annotated with amplitude A and phase θ . The function of the colour coding is just to show the progress of the vector summation of the scattering coefficient components $S^{(k)}$ with amplitude $A^{(k)}$ and phase $\theta^{(k)}$.

We have assumed that the phase components are uniformly distributed over all angles, denoted as: $\theta^{(k)} \sim U[0, 2\pi]$. This implies that the angle of the coherent sum is also uniform over the same interval: $\theta \sim U[0, 2\pi]$. This condition defines the *strong scattering* regime [Barakat, 1986, Jakeman and Tough, 1987]. It means that the random walk has

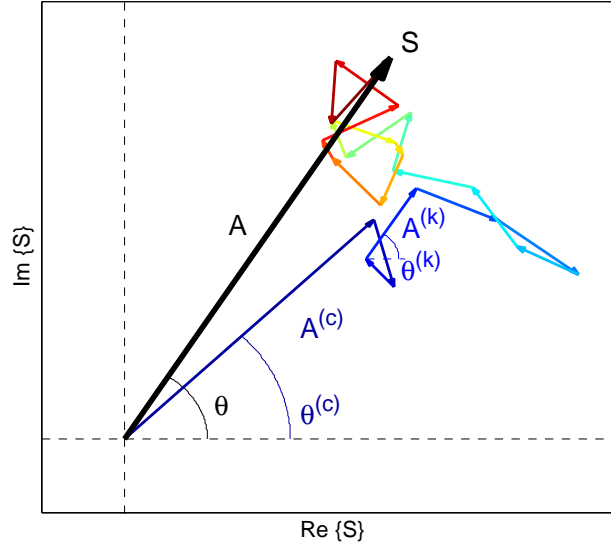


Figure 2.7: Random walk model of weak scattering. The scattering coefficient S with amplitude A and phase θ is shown as a phasor in the complex plane (black arrow). It is decomposed as the sum of a dominating coherent term S_c , with amplitude $A^{(c)}$ and phase $\theta^{(c)}$, and an incoherent term equal to the sum of random phasors described in Figure 2.6, contributed by N distributed scatterers.

an equal chance of ending up in any angular sector of the complex plane, and moreover that the coherent sum is a random variable with zero mean, even though the amplitude is positive. The zero mean of S is sometimes expressed as vanishing of the coherent, or specular, component [Ruffing and Fleischer, 1985]. This is an equivalent condition for strong scattering. A scattering process with the zero mean property produces *fully developed speckle*, a term which is explained in more detail in Section 2.3.4.

Weak Scattering

A non-uniform distribution of the phase components, written as $\theta^{(k)} \approx U[0, 2\pi]$, results in a non-uniform distribution of θ and a non-zero mean of S , which thus has a non-vanishing incoherent or specular component. These are the characteristics of a *weak scattering* process, which produces a *partially developed speckle* pattern [Ruffing and Fleischer, 1985, Barakat, 1986]. The weak scattering regime is illustrated by Figure 2.7. Compared to Figure 2.6 and Equation (2.15), a dominant scatterer has been added, meaning a scatterer whose reflection is much stronger than the others. The random walk model has simply been translated in the complex plane by the first vector component, which represents the dominant scatterer.

The described weak scattering process can be expressed as

$$S = S^{(c)} + S^{(i)} = A^{(c)} e^{j\theta^{(c)}} + \sum_{k=1}^N A^{(k)} e^{j\theta^{(k)}} = A e^{j\theta}, \quad (2.16)$$

where S is a sum of a coherent term, $S^{(c)}$, and an incoherent term, $S^{(i)}$. The coherent term has deterministic amplitude $A^{(c)}$ and phase $\theta^{(c)}$, that can be modelled with degenerate distributions⁶. The incoherent term represents the stochastic model of strong scattering from (2.15), and contains the same random amplitude and phase components. We thus have a set of amplitude components, $\mathcal{A} = \{A^{(k)}, A^{(1)}, \dots, A^{(N)}\}$, and a set of phase components, $\Theta = \{\theta^{(c)}, \theta^{(1)}, \dots, \theta^{(N)}\}$. The distribution of Θ becomes non-uniform, since $\theta^{(c)}$ is a deterministic component. In the described case, the IID assumptions are also violated. However, weak scattering can occur without the presence of dominant scatterers, if only the phase component distribution is non-uniform. See for instance [Jakeman and Tough, 1987], where the $\{\theta^{(k)}\}$ are assumed to follow a von Mises distribution.

2.3.2 Gaussian Model

Single-look Complex Data

The simplest statistical model for the radar return is based on the random walk model in (2.15) under all three assumptions listed in Section 2.3.1. Further assume that the number of scatterers, N , is large and constant. It follows from the central limit theorem that the scattering coefficient S converges in distribution to a complex Gaussian random variable with zero mean and variance σ^2 as $N \rightarrow \infty$, where σ^2 is the mean radar cross section of the scattering medium. The PDF of S is given by

$$f_S(S; \sigma^2) = \frac{1}{\pi\sigma^2} \exp\left(-\frac{S^2}{\sigma^2}\right). \quad (2.17)$$

In the polarimetric case, the zero mean scattering vector \mathbf{s} follows a multivariate circular complex Gaussian distribution with covariance matrix $\mathbf{\Sigma} = \text{E}\{\mathbf{s}\mathbf{s}^H\}$. The covariance matrix contains the complex cross-correlation between scattering coefficients at different polarimetric channels. Let S_i and S_k be the scattering coefficients of two polarimetric channels. They are decomposed into their real and complex parts by: $S_i = \text{Re}(S_i) + j\text{Im}(S_i)$. The *circularity* property of complex Gaussian distribution is defined by the correlations [van den Bos, 1994, Goodman, 2007]

$$\text{E}\{\text{Re}(S_i) \text{Re}(S_k)\} = \text{E}\{\text{Im}(S_i) \text{Im}(S_k)\}, \quad (2.18)$$

$$\text{E}\{\text{Re}(S_i) \text{Im}(S_k)\} = -\text{E}\{\text{Re}(S_k) \text{Im}(S_i)\}, \quad (2.19)$$

that follow from the assumptions on the $\{A_k\}$ and the $\{\theta_k\}$ in the random walk model. The cross-correlations also imply that

$$\text{E}\{S_i S_k\} = 0. \quad (2.20)$$

The PDF of \mathbf{s} is [van den Bos, 1994]

$$f_{\mathbf{s}}(\mathbf{s}; \mathbf{\Sigma}) = \frac{1}{\pi^d |\mathbf{\Sigma}|} \exp(-\mathbf{s}^H \mathbf{\Sigma}^{-1} \mathbf{s}). \quad (2.21)$$

⁶A deterministic value x_0 can be modelled as a random variable with the Dirac delta function, $\delta(x-x_0)$, as its PDF. This is referred to as a *degenerate* distribution.

The derivations for the lexicographic basis and Pauli basis cases are analogous. The following presentation is therefore restricted to results for the former, without loss of generality.

We see that the random walk model of strong scattering leads to complex Gaussian distributions for SLC data. We will now present distributions for MLC data, but note that the term *Gaussian model* is often used somewhat imprecisely about other data formats, whose distributions are derived directly from the complex Gaussian distributions of the SLC data by appropriate transformations of the random variates. Thus, a Rayleigh distributed amplitude: $A = |S^2|$, an exponentially distributed single-look intensity: $I = A^2$, and a gamma distributed multilook intensity: $I_L = \langle I \rangle = \langle A^2 \rangle$ are considered to belong to the Gaussian model. The derivation of all these distributions are shown in [Oliver and Quegan, 2004]. In line with this usage of terminology, the scaled complex Wishart distribution, presented in the next section, can be considered as the multilook polarimetric extension of the Gaussian model [Doulgeris et al., 2008, Vasile et al., 2009].

Multilook Complex Data

Assume that we have a multilook sample of L independent scattering vectors that are multivariate circular complex Gaussian with dimension d , and that $L \geq d$. The sample covariance matrix \mathbf{C} computed from this sample will then be non-singular, and belongs to the cone of positive definite complex Hermitian matrices, denoted as Ω_+ . It follows what we refer to as the *scaled complex Wishart distribution*. The PDF of \mathbf{C} under the Gaussian model is given by

$$f_{\mathbf{C}}(\mathbf{C}; L, \Sigma) = \frac{L^{Ld}}{\Gamma_d(L)} \frac{|\mathbf{C}|^{L-d}}{|\Sigma|^L} \exp(\text{tr}(-L\Sigma^{-1}\mathbf{C})), \quad (2.22)$$

which is defined on Ω_+ , where $|\cdot|$ is the determinant operator, $\text{tr}(\cdot)$ is the trace operator and $\Gamma_d(L)$ is the multivariate gamma function of the complex kind, defined in (4.3). The true complex Wishart distribution [Goodman, 1963] describes the matrix variable $\mathbf{W} = L\mathbf{C}$, hence (2.22) follows by a linear transformation.

The distribution parameters are the shape parameter L and the scale matrix $\Sigma = E\{\mathbf{C}\}$. In radar statistics, L is recognised as the number of looks, with reference to the multilooking process. In other contexts, it is often referred to as the degrees of freedom. The equivalent number of looks, discussed in 2.2.4, should be inserted for L to account for correlation between the looks. The substituted parameter is still subject to the condition of non-singularity.

2.3.3 Non-Gaussian Model

It has been shown experimentally that the Gaussian model, with its bundle of equivalent distributions for different data formats, presents a good fit to real radar image data



Figure 2.8: Example of the appearance of the speckle pattern in a radar image.

when the scene is homogeneous, with low to moderate roughness and a high number of scatterers. Nevertheless, there is abundant empirical evidence that real data deviate from the model too, especially for images of urban environment, but also for natural terrain, such as rough sea and forest in general (See e.g. [Jakeman and Pusey, 1976, Oliver, 2000, Tison et al., 2004]). This is generally explained by the notion of *texture*, thought of as variations in the mean radar reflectivity between pixels with the same thematic content, which is not accounted for in the Gaussian model. The texture term is defined and discussed in more detail in Section 2.3.5. The Gaussian model only encompasses statistical variation attributed to *speckle*, the interference pattern produced by the coherent sum in the random walk model, which is the topic of the next section.

Several distributions have been proposed for single polarisation amplitude and intensity, that imply non-Gaussian statistics for the scattering coefficient. The Weibull distribution and the log-normal distribution are two of the most popular examples [Oliver and Quegan, 2004]. Even though neither of them bear links to physical modelling of the scattering process, they have been shown to provide reasonably good fit to real data covering rough surfaces. However, they have not yet been extended to matrix-variate statistics and cannot be used to model the polarimetric covariance matrix. This has only been achieved for distributions derived from the doubly stochastic product model, which models the contribution of speckle and texture as independent random variates. The multilook polarimetric version of this model is presented in Section 2.3.6.

2.3.4 Speckle

Speckle is often referred to as noise, but this strictly is a misnomer. Although the appearance of the speckle patterns is granular, noise-like, and strongly limits the interpretabil-

ity of the radar image, the interference causing it is an inherent property of the measurement process that is common to all kind of coherent imaging. The same effect is found in laser images, and much of the fundamental research utilised in statistical analysis of radar images is taken from the field of laser optics (see e.g. [Goodman, 1975, Jakeman and Pusey, 1976]). Speckle is also an artifact of B-mode ultrasound sonography, where a linear array of acoustic transducers simultaneously scans a plane through an object that can be viewed as a two-dimensional image (see e.g. [Wagner et al., 1987, Eltoft, 2006]). In all these cases, the speckle pattern occurs because the measured field results from an incoherent scattering process. An example of a radar speckle is shown in Figure 2.8.

Radar measurements are commonly described by a multiplicative signal model⁷. The first factor in the product is the mean radar reflectivity, defined as the mean fraction of electromagnetic energy reflected by the scattering medium. The other factor is a random variable which models the variation due to speckle. The main characteristic of the product model is that the statistical variation is proportional to the reflectivity. This has a severely distorting effect on the image analysis, considering that the information we normally want to infer is related to the mean reflectivity. In the *doubly stochastic product model*, we also treat the mean reflectivity as a random variable, to obtain a more realistic description of the radar measurements.

The distinction between fully developed and partially developed speckle was made in Section 2.3.1. Fully developed speckle is described by the Gaussian model. Partially developed speckle is modelled by the Rician distribution for the single polarisation amplitude [Goodman, 2000]. The doubly stochastic product model for polarimetric SLC data has been extended to allow for a coherent component in [Eltoft et al., 2006], but has yet to be amended for MLC data, at least when it comes to deriving a closed form expression for the PDF.

2.3.5 Texture

In the field of image processing, the *texture* term is commonly used when referring to the deterministic or stochastic structure of an image region, characterised in terms of the spatial arrangement and directional alignment of the pixel intensities (in gray-scale or colour). As noted in [Tuceryan and Jain, 1994], «we recognize texture when we see it but it is very difficult to define». They list a number of proposed universal definitions, but it may prove more productive to search for a description of texture which is appropriate for the application at hand. In radar imaging, for instance, the texture term has indeed received its own meaning.

When we talk about texture in the context of radar images and the distributions of their data, it normally refers to variations in the mean radar reflectivity, as discussed in Section 2.3.4. We can model the mean reflectivity as a random variable, which is then referred to as a *texture variable*. The texture variable may also absorb other types

⁷According to commonly used terminology, which was argued against in the above, the radar images are subject to multiplicative noise.

of variation, such as the effect of inhomogeneity or mixed targets within a resolution cell. For instance, urban areas exhibit large variations that are badly described by the Gaussian model. This variation can be attributed to the mixture of different scattering media, just as much as the spatial variation of reflectivity within a single medium.

In the doubly stochastic product model, we assign a distribution model to the texture variable and estimate the parameters of the distribution. Textural features, defined in a multitude of ways, contain valuable information for delineation and recognition of image segments and classes in many pattern recognition problems. Classical texture features have also been used in analysis of radar images [Clausi and Yue, 2004, De Grandi et al., 2009]. A less travelled path is to exploit the parameters of the texture distributions, that contain radar specific textural information, as features in radar image analysis tasks, such as segmentation and classification [Oliver, 2000].

The distinction between *texture modulated speckle* and *pure speckle* is finally introduced, where the latter refers to the variation pattern created by the interference phenomenon alone. By these terms, we establish terminology to express clearly whether texture is affecting the speckle pattern and has been included in the signal model.

2.3.6 Multilook Polarimetric Product Model

Let $\widetilde{\mathbf{W}}$ be a scaled complex Wishart distributed random matrix, written as $\widetilde{\mathbf{W}} \sim s\mathcal{W}_d^c(L, \Sigma)$. The doubly stochastic product model for multilook polarimetric data is given by

$$\mathbf{C} = T \cdot \widetilde{\mathbf{W}}. \quad (2.23)$$

The variation of the polarimetric covariance matrix \mathbf{C} is decomposed into the contribution of texture and fully developed speckle, where the scalar random variable $T \in \mathbb{R}^+$ represents texture and the random matrix $\widetilde{\mathbf{W}} \in \Omega_+$ represents speckle. The texture variable must be strictly positive, and is normally assumed to have unit mean, hence all scale information is put into $\widetilde{\mathbf{W}}$ (or $|\widetilde{\mathbf{W}}|$). This is the convention used in this thesis. The alternative solution is to normalise $\widetilde{\mathbf{W}}$, as it is done in [Doulgeris and Eltoft, 2010]. This yields a matrix variable which contains only information about polarimetry, in form of covariance structure between the channels, while all scale information is put into the texture variable.

By using a scalar random texture variable, it is implicitly assumed that the textural variation is the same in all polarimetric channels. This is not necessarily supported by real data, and is a limitation of the model. It has been suggested that one should at least use distinct texture variables for co-polarised and cross-polarised channels. An extended doubly stochastic product model with multichannel texture has been proposed in [Yu, 1998, Zou et al., 2000] on the form

$$\mathbf{C} = \mathbf{T} \cdot \widetilde{\mathbf{W}}, \quad (2.24)$$

with the random diagonal multi-texture matrix defined as

$$\mathbf{T} = \text{diag}(T_1, T_2, \dots, T_d) = \begin{bmatrix} T_1 & 0 & \cdots & 0 \\ 0 & T_2 & & \vdots \\ \vdots & & \ddots & 0 \\ 0 & \cdots & 0 & T_d \end{bmatrix}. \quad (2.25)$$

Another objection to the doubly stochastic product model is that it is not adequate for co-polarised channels. Analysis of real data has shown that there is an additive term superimposed on the product of texture and fully developed speckle [Séry and Lopès, 1997, López-Martínez and Fàbregas, 2005]. The extensions to multi-texture and a mixed multiplicative additive signal model are not considered in this thesis, but left as topics of future research.

The distribution of \mathbf{C} depends on the chosen distribution for T . For a general texture distribution $f_T(T; \boldsymbol{\theta})$, whose parameters are stored in the texture parameter vector $\boldsymbol{\theta}$, the PDF is given by

$$f_{\mathbf{C}}(\mathbf{C}; L, \boldsymbol{\Sigma}, \boldsymbol{\theta}) = \int_0^\infty f_{\mathbf{C}|T}(\mathbf{C}|T; L, \boldsymbol{\Sigma}) f_T(T; \boldsymbol{\theta}) dT. \quad (2.26)$$

It will be seen that this integral is a matrix-variate Mellin convolution, as defined in Chapter 3. The distribution of $\mathbf{C}|T$ is recognised as the scaled complex Wishart distribution from (2.22). Closed form expressions of $f_{\mathbf{C}}(\mathbf{C}; L, \boldsymbol{\Sigma}, \boldsymbol{\theta})$ has been obtained for a limited number of texture distributions. For gamma distributed texture, the matrix-variate \mathcal{K} distribution was derived in [Lee et al., 1994]. This distribution has shown the same merits as the \mathcal{K} distributions for single polarisation amplitude and intensity (See [Oliver and Quegan, 2004]). It is versatile and provides good fit for natural surfaces with a certain degree of heterogeneity, such as rough sea and forest. For inverse gamma distributed texture, the matrix-variate \mathcal{G}^0 distribution was derived in [Freitas et al., 2005]. It has its strength for extremely heterogeneous surfaces, such as urban areas. Both the mentioned distributions are special cases of the matrix-variate \mathcal{U} distribution, derived in [Bombrun and Beaulieu, 2008] from Fisher-Snedecor distributed texture. Because it has two texture parameters, one more than the \mathcal{K} and \mathcal{G}^0 distributions, it covers many types of surfaces between the heterogeneous and the extremely heterogeneous case. The added flexibility in terms of modelling capability comes at the cost of more complicated and less accurate parameter estimation, as demonstrated in the paper of Chapter 6.

The mathematical details of the described distributions for the polarimetric covariance matrix is given in Chapter 4. For each distribution the PDF is presented, together with the matrix-variate Mellin kind statistics, that are defined in the next chapter.

Chapter 3

The Mellin Transform

This chapter gives an historical overview which motivates the application of the Mellin transform to statistical analysis of radar data distributions. The univariate and matrix-variate Mellin transform are defined, and the fundamental properties of the latter derived. Some of the main contributions of the thesis are revealed as the framework of matrix-variate Mellin kind statistics is presented.

3.1 An Historical Note

The Mellin transform first appeared in a memoir by the German mathematician Bernard Riemann (1826-1866) about the zeta function, but it is named after the Finnish mathematician *Robert Hjalmar Mellin* (1854-1933), who gave a systematic treatment of the transform and its inverse [Butzer and Jansche, 1997]. In contrast to the Fourier transform and the Laplace transforms, it arose in a mathematical context. Since then, it has also found numerous applications in many areas of physics, statistics and engineering, as reviewed in [Bertrand et al., 2000].

The Mellin transform has been applied to a wealth of problems in analytic combinatorics and analysis of algorithms [Flajolet et al., 1995, Szpankowski, 2001]. It is a key component in systematic methods for evaluation of integrals [Marichev, 1982, Fikioris, 2006, Fikioris, 2007]. In signal processing, it has been applied to time-frequency analysis on logarithmic scale, and its scale invariance property has been utilised to construct affine transformations [Bertrand et al., 1990, Ovarlez et al., 1992, Cohen, 1993, Ruth and Gilbert, 1994, Nelson, 1995, Kaiser, 1996].

The application of the Mellin transform to analysis of matrix-variate radar data builds on results in probability and the theory of statistical distributions. Benjamin Epstein was first to note that the Mellin transform is the «natural analytical tool to use in studying the distribution of products and quotients of independent random variables» [Epstein, 1948], after its utility in analysis of multivariate problems was indicated in [Nair, 1939]. This spurred a series of papers deriving distributions of product and quotients in one or more variables [Dolan, 1964, Springer and Thompson, 1966, Lom-

nicki, 1967, Springer and Thompson, 1970]. Products of dependent random variables were covered in [Subrahmaniam, 1970]. A good overview of the early literature is given in [Cook, 1981].

It will be seen in the following that the Mellin transform is related to the Fourier transform and the bilateral Laplace transform applied to logarithmically transformed data. It is an efficient tool to derive logarithmic moments and cumulants for products and quotients of random variables, and their matrix-variate generalisations, from which we obtain estimators for the distribution parameters with excellent statistical properties. Such logarithmic statistics and parameter estimators have been derived earlier in the univariate case, without explicit reference to the Mellin transform [Stacy, 1962, Stacy and Mihram, 1965, Hoekman, 1991, Kreithen and Hogan, 1991, Blacknell, 1994]. It was *Jean-Marie Nicolas* who developed these ideas into a systematic theory on logarithmic statistics, characterisation of data radar distributions, and estimation of their parameters, with the Mellin transform as the cornerstone [Nicolas, 2002, Nicolas, 2006]. The framework offered by Nicolas has triggered much research activity and many new applications to analysis of radar data. An extensive reference list is given in Paper 2 of Chapter 6.

3.2 Mellin Kind Statistics

Nicolas developed a framework for analysis of random variables defined on \mathbb{R}^+ , which he called *second kind statistics*. The results presented in this thesis extend the framework to random matrices defined on Ω_+ , and renames it as *Mellin kind statistics*, which is the term consistently used from this point. Another appropriate name used by some authors is log-statistics.

The starting point for Nicolas is the definition of the *Mellin kind characteristic function*. His ingenious trick is to replace the Fourier transform in the definition of the conventional characteristic function with the Mellin transform. The *Mellin kind moments* retrieved from the new characteristic function are seen to be moments computed on logarithmic scale. The *Mellin type cumulant generating function* is then defined as the logarithm of the Mellin type characteristic function. From this function we can retrieve *Mellin type cumulants*, that have some very appealing properties. When the target of analysis can be modelled as a product of random variates, the Mellin type cumulant will separate the contribution of the factors [Nicolas, 2002, Nicolas, 2006].

Analysis of radar data on logarithmic scale is not a new and revolutionary concept. Radar images have long been visualised on a logarithmic decibel scale to make more efficient use of the dynamic range of pixel intensities and improve the contrast. The homomorphic speckle filter uses a logarithmic transformation to transform the product model into an additive model, which is easier to handle with traditional signal processing methods [Franceschetti et al., 1995, Solbø and Eltoft, 2004]. The separability induced by the Mellin transform therefore comes as no great surprise. The ease of the mathematical derivations, and the simple expressions obtained for the Mellin kind statistics

of common radar distributions (that is, their characteristic function, moments, cumulant generating function and cumulants) is a more striking result. They underline the appropriateness of the approach and resonate with the words of Epstein: The Mellin transform is truly a natural tool for analysing products of random variables, and thus for multilook radar image data under the product model.

The mathematical details of Nicolas' univariate theory are omitted here. A thorough review is given in the paper of Chapter 6, which stresses the analogy between the derivations for the univariate case describing multilook single polarisation data and the matrix-variate case representing multilook polarimetric data. The paper also repeats the traditional *Fourier kind statistics* and highlights their intrinsic link with the additive signal model. Here is only given a repetition of the univariate Mellin transform, before we proceed with the definition of the matrix-variate Mellin kind statistics.

3.3 Univariate Mellin Transform

Let $g(x)$ be a real-valued function of the variable x . The univariate Mellin transform of $g(x)$ is given by

$$\mathcal{M}\{g(x)\}(s) = \int_0^{\infty} x^{s-1} g(x) dx = \mathfrak{G}(s), \quad (3.1)$$

where $s \in \mathbb{C}$ is the complex transform variable. Further let the Fourier transform be defined as

$$\mathcal{F}\{g(x)\}(\xi) = \int_{-\infty}^{+\infty} e^{-2\pi j x \xi} g(x) dx \quad (3.2)$$

and the bilateral or two-sided Laplace transform as

$$\mathcal{L}_B\{g(x)\}(s) = \int_{-\infty}^{+\infty} e^{-sx} g(x) dx. \quad (3.3)$$

The relations between the transforms are shown to be [Bertrand et al., 2000]

$$\mathcal{M}\{g(x)\}(\sigma + j2\pi\xi) = \mathcal{F}\{e^{-\sigma x} g(e^{-x})\}(\xi), \quad (3.4)$$

$$\mathcal{M}\{g(x)\}(s) = \mathcal{L}_B\{g(e^{-x})\}(s) \quad (3.5)$$

by virtue of the substitution of $s = \sigma + j2\pi\xi$. The inverse relations are

$$\mathcal{F}\{g(x)\}(\xi) = \mathcal{M}\{x^{-\sigma} g(e^{-\ln y})\}(s), \quad (3.6)$$

$$\mathcal{L}_B\{g(x)\}(s) = \mathcal{M}\{g(-\ln x)\}(s). \quad (3.7)$$

From Lebesgue's dominated convergence theorem [Bartle, 1995], we know that the ordinary unilateral Laplace transform converges absolutely in the semi-plane $\text{Re}(s) > b$ for some constant b , and possibly at points on the line $\text{Re}(s) = b$. The intersection of the regions of convergence for the left-sided and the right-sided Laplace transform form

the region of converge for the bilateral Laplace transform. This becomes the region of holomorphy, also known as the analytic strip, and takes the form $b < \Re(s) < c$, for some real constants b and c , possibly including the boundary lines. When transforming a PDF defined on \mathbb{R}^+ , the unilateral (right-sided) and bilateral Laplace transforms are identical, and the analytic strip becomes the semi-plane $\Re(s) > b$.

It can be seen directly from (3.5) that the analytic strip of the Mellin transform is the same as for the bilateral Laplace transform. The boundaries depends on the transformed function, and the strip may even extend to the whole complex plane. The inverse Mellin transform is defined by [Bertrand et al., 2000]

$$\mathcal{M}^{-1}\{\mathcal{G}(s)\} = \frac{1}{2\pi j} \int_{c-j\infty}^{c+j\infty} x^{-s} \mathcal{G}(s) ds. \quad (3.8)$$

The notation for the integral limits implies that this is a line integral taken over a vertical line $\Re(s) = c$ in the complex plane, which must lie within the analytic strip.

Comprehensive accounts on the univariate Mellin transform, including lists of fundamental properties, examples of important transforms and tables of basic integrals, are found in [Poularikas, 1999, Bertrand et al., 2000, Szpankowski, 2001, Debnath and Bhatta, 2007]. An extension of the transform to complex variables is given in [Kotlarski, 1965] and to multivariate variables in [Mathur and Krishna, 1977]. The extension to the matrix-variate case is the next topic.

3.4 Matrix-Variate Mellin Transform

A generalised transform for functions of real matrices, named the M-transform, was defined in [Mathai, 1978]. Mathai also referred to it as the generalised Mellin transform in [Mathai, 1981]. The extension to functions of complex matrices was presented in [Mathai, 1997]. In the following, let $g(\mathbf{X})$ be a real-valued scalar function defined on Ω_+ , and let g be symmetric in the sense: $g(\mathbf{X}\mathbf{Y}) = g(\mathbf{Y}\mathbf{X}) = g(\mathbf{X}^{1/2}\mathbf{Y}\mathbf{X}^{1/2}) = g(\mathbf{Y}^{1/2}\mathbf{X}\mathbf{Y}^{1/2})$, where $\mathbf{X}, \mathbf{Y} \in \Omega_+$. The matrix square roots defined by $\mathbf{X} = \mathbf{X}^{1/2}\mathbf{X}^{1/2}$ and $\mathbf{Y} = \mathbf{Y}^{1/2}\mathbf{Y}^{1/2}$ are therefore guaranteed to exist. Whenever the integral exists, the complex matrix-variate Mellin transform of $g(\mathbf{X})$ is defined by

$$\mathcal{M}\{g(\mathbf{X})\}(s) = \int_{\Omega_+} |\mathbf{X}|^{s-d} g(\mathbf{X}) d\mathbf{X} = G(s). \quad (3.9)$$

We note that $\mathcal{M}\{g(\mathbf{X})\}(s)$ is a function of a complex scalar transform variable s , whereas $g(\mathbf{X})$ is defined on a matrix space, thus the transform is not unique and has no inverse. The symmetry requirement restricts the functions (3.9) can be applied to, but does not pose any problem for the distributions of the polarimetric covariance matrix, derived from the doubly stochastic product model. We may therefore use the transform to define Mellin kind statistics for the complex matrix-variate case.

The univariate Mellin transform has a convolution theorem and a correlation theorem. To derive the analogue relations for the matrix-variate Mellin transform, we need some definitions. Let $g(\mathbf{X})$ and $h(\mathbf{X})$ be two real-valued scalar functions defined on Ω_+ , symmetric by the definition given above. The *matrix-variate Mellin convolution* is introduced by the definition

$$\begin{aligned} (g \hat{\star} h)(\mathbf{X}) &= \int_{\Omega_+} |\mathbf{Y}|^{-d} g(\mathbf{Y}^{-\frac{1}{2}} \mathbf{X} \mathbf{Y}^{-\frac{1}{2}}) h(\mathbf{Y}) d\mathbf{Y} \\ &= \int_{\Omega_+} |\mathbf{Y}|^{-d} h(\mathbf{Y}^{-\frac{1}{2}} \mathbf{X} \mathbf{Y}^{-\frac{1}{2}}) g(\mathbf{Y}) d\mathbf{Y}, \end{aligned} \quad (3.10)$$

which is an associative and commutative operation. The *matrix-variate Mellin correlation* is also introduced, defined by

$$(g \hat{\otimes} h)(\mathbf{X}) = \int_{\Omega_+} |\mathbf{Y}|^d g(\mathbf{Y}^{\frac{1}{2}} \mathbf{X} \mathbf{Y}^{\frac{1}{2}}) h(\mathbf{Y}) d\mathbf{Y}. \quad (3.11)$$

This operation is neither associative nor commutative. It reduces to the univariate Mellin correlation defined in [Nicolas, 2002, Nicolas, 2006].

3.5 Fundamental Properties

It is now time to deduce some fundamental properties of the matrix-variate Mellin transform, in analogy with the fundamental properties that have been derived for the univariate transform.

Property 1 (Scaling by nonsingular matrix of constants):

$$\mathcal{M}\{g(\mathbf{A}\mathbf{X})\}(s) = |\mathbf{A}|^{-s} G(s). \quad (3.12)$$

Proof 1: The transformation $\mathbf{Y} = \mathbf{A}\mathbf{X}$ has Jacobian determinant $|J(\mathbf{X} \rightarrow \mathbf{Y})| = |\mathbf{A}|^d$ [Mathai, 1997]. Thus,

$$\begin{aligned} \mathcal{M}\{g(\mathbf{A}\mathbf{X})\}(s) &= \int_{\Omega_+} |\mathbf{X}|^{s-d} g(\mathbf{A}\mathbf{X}) d\mathbf{X} \\ &= \int_{\Omega_+} (|\mathbf{Y}|/|\mathbf{A}|)^{s-d} g(\mathbf{Y}) |\mathbf{A}|^{-d} d\mathbf{Y} = |\mathbf{A}|^{-s} \int_{\Omega_+} |\mathbf{Y}|^{s-d} g(\mathbf{Y}) d\mathbf{Y}, \end{aligned}$$

which is equivalent to (3.12). ■

Property 2 (Scaling by scalar constant):

$$\mathcal{M}\{g(a\mathbf{X})\}(s) = a^{-ds} G(s). \quad (3.13)$$

Proof 2: This is a special case of (3.12) with $\mathbf{A} = a\mathbf{I}_d$, where \mathbf{I}_d is the $d \times d$ identity matrix. Equation (3.13) follows from $|\mathbf{A}| = a^d$. ■

Property 3 (Multiplication by $|\mathbf{X}|^a$):

$$\mathcal{M}\{|\mathbf{X}|^a g(\mathbf{X})\}(s) = G(s + a). \quad (3.14)$$

Proof 3: The proof is trivial. Take the transform of $|\mathbf{X}|^a g(\mathbf{X})$:

$$\mathcal{M}\{|\mathbf{X}|^a g(\mathbf{X})\}(s) = \int_{\Omega_+} |\mathbf{X}|^{s-d} |\mathbf{X}|^a g(\mathbf{X}) d\mathbf{X} = \int_{\Omega_+} |\mathbf{X}|^{(s+a)-d} g(\mathbf{X}) d\mathbf{X}$$

and identify this as (3.14). ■

Property 4 (Raising the independent variable to an integer power):

$$\mathcal{M}\{g(\mathbf{X}^a)\}(s) = \frac{1}{a^d} G\left(\frac{s}{a}\right). \quad (3.15)$$

Proof 4: We use the transformation $\mathbf{Y} = \mathbf{X}^a = \overbrace{\mathbf{X} \cdots \mathbf{X}}^a$ with differential relation $d\mathbf{Y} = a^d |\mathbf{X}|^{(a-1)d} d\mathbf{X}$ to show that

$$\begin{aligned} \mathcal{M}\{g(\mathbf{X}^a)\}(s) &= \int_{\Omega_+} |\mathbf{X}|^{s-d} g(\mathbf{X}^a) d\mathbf{X} \\ &= \int_{\Omega_+} |\mathbf{Y}|^{\frac{s-d}{a}} g(\mathbf{Y}) \left(a^{-d} |\mathbf{Y}|^{-\frac{(a-1)d}{a}} d\mathbf{Y}\right) \\ &= \frac{1}{a^d} \int_{\Omega_+} |\mathbf{Y}|^{\frac{s}{a}-d} d\mathbf{Y} = \frac{1}{a^d} G\left(\frac{s}{a}\right). \quad \blacksquare \end{aligned}$$

Property 5 (Inverting the independent variable):

$$\mathcal{M}\{g(\mathbf{X}^{-1})\}(s) = (-1)^d G(-s). \quad (3.16)$$

Proof 5: This is a special case of (3.15) with $a = -1$. ■

Property 6 (Inverting and multiplying by the independent variable):

$$\mathcal{M}\{|\mathbf{X}|^{-1} g(\mathbf{X}^{-1})\}(s) = (-1)^d G(1 - s). \quad (3.17)$$

Proof 6: Use the transformation $\mathbf{Y} = \mathbf{X}^{-1}$ with differential relation $d\mathbf{Y} = (-1)^d |\mathbf{X}|^{-2d}$, which yields

$$\begin{aligned} \mathcal{M}\{|\mathbf{X}|^{-1}g(\mathbf{X}^{-1})\}(s) &= \int_{\Omega_+} |\mathbf{X}|^{s-d} |\mathbf{X}|^{-1} g(\mathbf{X}^{-1}) d\mathbf{X} \\ &= \int_{\Omega_+} |\mathbf{Y}|^{-(s-d-1)} g(\mathbf{Y}) ((-1)^{-d} |\mathbf{Y}|^{-2d} d\mathbf{Y}) \\ &= (-1)^d \int_{\Omega_+} |\mathbf{Y}|^{1-s-d} g(\mathbf{Y}) d\mathbf{Y} = (-1)^d G(1-s). \quad \blacksquare \end{aligned}$$

Property 7 (Multiplication by $\ln |\mathbf{X}|$):

$$\mathcal{M}\{\ln |\mathbf{X}|g(\mathbf{X})\}(s) = \frac{d}{ds} G(s). \quad (3.18)$$

Proof 7: By using the result

$$\frac{d}{ds} |\mathbf{X}|^{s-d} = \frac{d}{ds} e^{\ln |\mathbf{X}|(s-d)} = |\mathbf{X}|^{s-d} \ln |\mathbf{X}|$$

and Leibniz's integral rule, we show that

$$\begin{aligned} \frac{d}{ds} G(s) &= \int_{\Omega_+} \left(\frac{d}{ds} |\mathbf{X}|^{s-d} \right) g(\mathbf{X}) d\mathbf{X} \\ &= \int_{\Omega_+} |\mathbf{X}|^{s-d} \ln |\mathbf{X}| g(\mathbf{X}) d\mathbf{X} = \mathcal{M}\{\ln |\mathbf{X}|g(\mathbf{X})\}. \quad \blacksquare \end{aligned}$$

Property 8 (Multiplication by a power of $\ln |\mathbf{X}|$):

$$\mathcal{M}\{(\ln |\mathbf{X}|)^\nu g(\mathbf{X})\}(s) = \frac{d^\nu}{ds^\nu} G(s). \quad (3.19)$$

Proof 8: We use the result

$$\frac{d^\nu}{ds^\nu} |\mathbf{X}|^{s-d} = |\mathbf{X}|^{s-d} (\ln |\mathbf{X}|)^\nu$$

and follow the approach of Proof 7. \blacksquare

Property 9 (Mellin convolution):

$$\mathcal{M}\{(g \hat{\star} h)(\mathbf{X})\}(s) = \mathcal{M}\{g(\mathbf{X})\}(s) \cdot \mathcal{M}\{h(\mathbf{X})\}(s). \quad (3.20)$$

Proof 9: Substitute $\mathbf{Z} = \mathbf{X}\mathbf{Y}$ and note that $\mathbf{Z} \in \Omega_+$ follows. We further have $\mathbf{X} = \mathbf{Y}^{-1/2}\mathbf{Z}\mathbf{Y}^{-1/2}$ and $d\mathbf{X} = d\mathbf{Z}/|\mathbf{Y}|^d$ [Mathai, 1997], which leads to

$$\begin{aligned} & \mathcal{M}\{g(\mathbf{X})\}(s) \cdot \mathcal{M}\{h(\mathbf{Y})\}(s) \\ &= \int_{\Omega_+} (|\mathbf{Z}|/|\mathbf{Y}|)^{s-d} g(\mathbf{Y}^{-\frac{1}{2}}\mathbf{Z}\mathbf{Y}^{-\frac{1}{2}}) |\mathbf{Y}|^{-d} d\mathbf{Z} \int_{\Omega_+} |\mathbf{Y}|^{s-d} h(\mathbf{Y}) d\mathbf{Y} \\ &= \int_{\Omega_+} |\mathbf{Z}|^{s-d} \left[\int_{\Omega_+} |\mathbf{Y}|^{-d} g(\mathbf{Y}^{-\frac{1}{2}}\mathbf{Z}\mathbf{Y}^{-\frac{1}{2}}) h(\mathbf{Y}) d\mathbf{Y} \right] d\mathbf{Z} \\ &= \mathcal{M}\{(g \hat{\star} h)(\mathbf{Z})\}(s), \end{aligned}$$

since the term in the square brackets can be identified as the Mellin convolution from (3.10). ■

Property 10 (Mellin correlation):

$$\mathcal{M}\{(g \hat{\otimes} h)(\mathbf{X})\}(s) = \mathcal{M}\{g(\mathbf{X})\}(s) \cdot \mathcal{M}\{h(\mathbf{X})\}(2d - s). \quad (3.21)$$

Proof 10: Use the substitution $\mathbf{X} = \mathbf{Y}\mathbf{Z}$ with differential relation $d\mathbf{X} = |\mathbf{Y}|^d d\mathbf{Z}$ to find

$$\begin{aligned} & \mathcal{M}\{g(\mathbf{X})\}(s) \cdot \mathcal{M}\{h(\mathbf{Y})\}(2d - s) \\ &= \int_{\Omega_+} (|\mathbf{Y}\mathbf{Z}|)^{s-d} g(\mathbf{Y}\mathbf{Z}) |\mathbf{Y}|^d d\mathbf{Z} \int_{\Omega_+} |\mathbf{Y}|^{d-s} h(\mathbf{Y}) d\mathbf{Y} \\ &= \int_{\Omega_+} |\mathbf{Z}|^{s-d} \left[\int_{\Omega_+} |\mathbf{Y}|^{-d} g(\mathbf{Y}\mathbf{Z}) h(\mathbf{Y}) d\mathbf{Y} \right] d\mathbf{Z} \\ &= \mathcal{M}\{(g \hat{\otimes} h)(\mathbf{X})\}(s), \end{aligned}$$

after the term in the square brackets is identified as the matrix-variate Mellin correlation defined in (3.11). ■

3.6 Matrix-variate Mellin Kind Statistics

The matrix-variate Mellin transform is defined on Ω_+ , and can therefore be applied to PDFs that have the same domain. It is this property that makes it relevant to the statistical analysis of distributions for the polarimetric covariance matrix. From now on, we assume that \mathbf{X} is a random matrix described by $f_{\mathbf{X}}(\mathbf{X})$, which is defined on Ω_+ .

3.6.1 Mellin Kind Characteristic Function

The Mellin kind characteristic function of a complex random matrix \mathbf{X} is defined as

$$\phi_{\mathbf{X}}(s) = \mathbb{E}\{|\mathbf{X}|^{s-d}\} = \mathcal{M}\{f_{\mathbf{X}}(\mathbf{X})\}(s). \quad (3.22)$$

when \mathbf{X} and its PDF, $p_{\mathbf{X}}(\mathbf{X})$, satisfy all requirements of (3.9). Note that the moments of the determinant $|\mathbf{X}|$ can be retrieved from the Mellin kind characteristic function as $\mathbb{E}\{|\mathbf{X}|^{\nu}\} = \phi_{\mathbf{X}}(\nu + d)$.

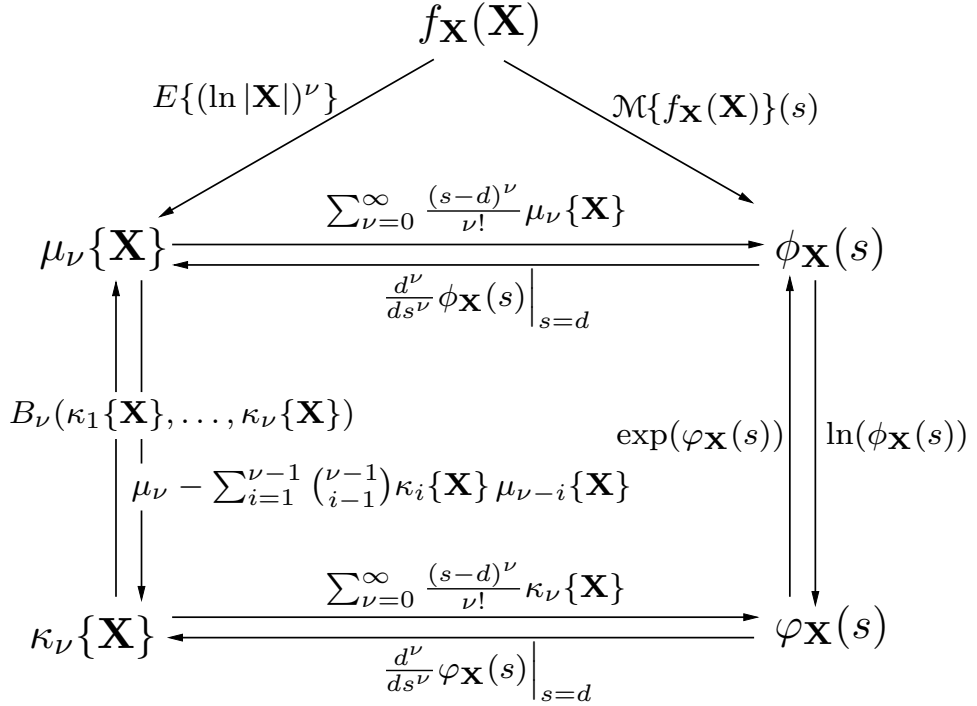


Figure 3.1: Relations between the components of matrix-variate Mellin kind statistics. The function $B_{\nu}(\cdot)$ is the complete Bell polynomial, defined in Paper 2 (Chapter 6).

3.6.2 Mellin Kind Matrix Moments

The ν th-order Mellin kind moment of \mathbf{X} is

$$\mu_{\nu}\{\mathbf{X}\} = \frac{d^{\nu}}{ds^{\nu}} \phi_{\mathbf{X}}(s) \Big|_{s=d}. \quad (3.23)$$

If all Mellin kind matrix moments exist, the Mellin kind characteristic function can be written as the power series expansion

$$\begin{aligned} \phi_{\mathbf{X}}(s) &= \int_{\Omega_+} e^{(s-d) \ln |\mathbf{X}|} f_{\mathbf{X}}(\mathbf{X}) d\mathbf{X} \\ &= \sum_{\nu=0}^{\infty} \frac{(s-d)^{\nu}}{\nu!} \mu_{\nu}\{\mathbf{X}\} \end{aligned} \quad (3.24)$$

in terms of the $\mu_{\nu}\{\mathbf{X}\}$. The derivation of (3.24) reveals that

$$\mu_{\nu}\{\mathbf{X}\} = E\{(\ln |\mathbf{X}|)^{\nu}\} = \int_{\Omega_+} (\ln |\mathbf{X}|)^{\nu} f_{\mathbf{X}}(\mathbf{X}) d\mathbf{X} \quad (3.25)$$

which justifies the denotation of $\mu_{\nu}\{\mathbf{X}\}$ as a matrix log-moment.

3.6.3 Mellin Kind Cumulant Generating Function

The Mellin kind cumulant generating function of \mathbf{X} is defined as

$$\varphi_{\mathbf{X}}(s) = \ln \phi_{\mathbf{X}}(s). \quad (3.26)$$

3.6.4 Mellin Kind Matrix Cumulants

The ν th-order Mellin kind cumulant of \mathbf{X} is

$$\kappa_{\nu}\{\mathbf{X}\} = \left. \frac{d^{\nu}}{ds^{\nu}} \varphi_{\mathbf{X}}(s) \right|_{s=d}. \quad (3.27)$$

When all Mellin kind matrix moments exist, the Mellin kind cumulant generating function can be expanded as

$$\varphi_{\mathbf{X}}(s) = \ln \phi_{\mathbf{X}}(s) = \sum_{\nu=0}^{\infty} \frac{(s-d)^{\nu}}{\nu!} \kappa_{\nu}\{\mathbf{X}\} \quad (3.28)$$

in terms of the $\kappa_{\nu}\{\mathbf{X}\}$, that are also called matrix log-cumulants.

The relations between the the functions and moments defined under the matrix-variate Mellin kind statistics framework are shown in Figure 3.1.

3.7 Multilook Polarimetric Product Model

This section shows how the matrix-variate Mellin kind statistics are applied to the doubly stochastic product model for multivariate polarimetric data. We recollect the model as

$$\mathbf{C} = \mathbf{T} \cdot \widetilde{\mathbf{W}}, \quad (3.29)$$

where \mathbf{C} is the polarimetric covariance matrix, decomposed in terms of the random matrices $\widetilde{\mathbf{W}} \sim s\mathcal{W}_d^{\mathbf{C}}(L, \Sigma)$ and \mathbf{T} , respectively modelling fully developed speckle and texture. The PDF of \mathbf{T} , $f_{\mathbf{T}}(\mathbf{T}; \boldsymbol{\theta})$, is left unspecified for the moment. We here assume that \mathbf{T} has equal diagonal entries, $\mathbf{T} = \text{diag}(T, \dots, T) = T \cdot \mathbf{I}_d$, thereby restricting the texture model to a single random variable representing all polarimetric channels. The advantage of writing the scalar texture variable as a matrix will become apparent.

3.7.1 Application of the Matrix-Variate Mellin Convolution

The distribution of \mathbf{C} can now be written as

$$f_{\mathbf{C}}(\mathbf{C}; L, \Sigma, \boldsymbol{\theta}) = \int_{\Omega_+} f_{\mathbf{C}|\mathbf{T}}(\mathbf{C}|\mathbf{T}; L, \Sigma) f_{\mathbf{T}}(\mathbf{T}; \boldsymbol{\theta}) d\mathbf{C}, \quad (3.30)$$

parametrised by the equivalent number of looks L , scale matrix Σ , and the texture parameter vector $\boldsymbol{\theta}$. The integral is identified as the matrix-variate Mellin convolution, defined in (3.10). From Property 9 of the matrix-variate Mellin transform, it follows that

$$\mathcal{M}\{f_{\mathbf{C}}(\mathbf{C}; L, \Sigma, \boldsymbol{\theta})\} = \mathcal{M}\{f_{\widetilde{\mathbf{W}}}(\widetilde{\mathbf{W}}; L, \Sigma)\} \cdot \mathcal{M}\{f_{\mathbf{T}}(\mathbf{T}; \boldsymbol{\theta})\}. \quad (3.31)$$

The implication in terms of the matrix-variate Mellin statistics defined in the previous section is

$$\phi_{\mathbf{C}}(s; L, \Sigma, \boldsymbol{\theta}) = \phi_{\widetilde{\mathbf{W}}}(s; L, \Sigma) \cdot \phi_{\mathbf{T}}(s; \boldsymbol{\theta}), \quad (3.32)$$

$$\varphi_{\mathbf{C}}(s; L, \Sigma, \boldsymbol{\theta}) = \varphi_{\widetilde{\mathbf{W}}}(s; L, \Sigma) + \varphi_{\mathbf{T}}(s; \boldsymbol{\theta}), \quad (3.33)$$

$$\kappa_{\nu}\{\mathbf{C}; L, \Sigma, \boldsymbol{\theta}\} = \kappa_{\nu}\{\widetilde{\mathbf{W}}; L, \Sigma\} + \kappa_{\nu}\{\mathbf{T}; \boldsymbol{\theta}\}. \quad (3.34)$$

These respective equations present the Mellin kind characteristic function, the Mellin kind cumulant generating function and the Mellin kind cumulants for the multivariate polarimetric product model. The Mellin kind moments can be retrieved from the cumulants by the relation

$$\kappa_{\nu}\{\cdot\} = \mu_{\nu}\{\cdot\} - \sum_{i=1}^{\nu-1} \binom{\nu-1}{i-1} \kappa_i\{\cdot\} \mu_{\nu-i}\{\cdot\}, \quad (3.35)$$

which is valid both in the univariate and matrix-variate case, and for moments and cumulants on linear or logarithmic scale.

3.7.2 Mellin Kind Statistics for the Multilook Polarimetric Product Model

The texture variable was written as a matrix so that the matrix-variate convolution property could be used directly. The Mellin kind statistics of \mathbf{T} must still be resolved, and they have been in the paper of Chapter 6. It is found that

$$\phi_{\mathbf{T}}(s; \boldsymbol{\theta}) = \phi_T(d(s-d) + 1; \boldsymbol{\theta}), \quad (3.36)$$

where $\phi_T(s)$ is the univariate Mellin kind characteristic function of the scalar variable T , derived in [Nicolas, 2002, Nicolas, 2006] as

$$\phi_T(s) = \int_0^{\infty} T^{s-1} f_T(T; \boldsymbol{\theta}) dT = \mathbb{E}\{T^{s-1}; \boldsymbol{\theta}\}. \quad (3.37)$$

The cumulant generating function for \mathbf{T} follows readily as

$$\varphi_{\mathbf{T}}(s; \boldsymbol{\theta}) = \varphi_T(d(s-d) + 1; \boldsymbol{\theta}), \quad (3.38)$$

with $\varphi_T(s; \boldsymbol{\theta}) = \ln \phi_T(s; \boldsymbol{\theta})$. The Mellin kind cumulants are retrieved as

$$\kappa_{\nu}\{\mathbf{T}\} = d^{\nu} \kappa_{\nu}\{T\}, \quad (3.39)$$

where the $\kappa_\nu\{T; \boldsymbol{\theta}\}$ are the ν th-order univariate Mellin kind cumulants defined in [Nicolas, 2002, Nicolas, 2006]. We thus arrive at the following Mellin kind statistics for the multivariate polarimetric product model:

$$\phi_{\mathbf{C}}(s; L, \boldsymbol{\Sigma}, \boldsymbol{\theta}) = \phi_{\widetilde{\mathbf{W}}}(s; L, \boldsymbol{\Sigma}) \cdot \phi_T(d(s-d) + 1; \boldsymbol{\theta}), \quad (3.40)$$

$$\varphi_{\mathbf{C}}(s; L, \boldsymbol{\Sigma}, \boldsymbol{\theta}) = \varphi_{\widetilde{\mathbf{W}}}(s; L, \boldsymbol{\Sigma}) + \varphi_T(d(s-d) + 1; \boldsymbol{\theta}), \quad (3.41)$$

$$\kappa_\nu\{\mathbf{C}; L, \boldsymbol{\Sigma}, \boldsymbol{\theta}\} = \kappa_\nu\{\widetilde{\mathbf{W}}; L, \boldsymbol{\Sigma}\} + d^\nu \kappa_\nu\{T; \boldsymbol{\theta}\}. \quad (3.42)$$

This is a general formulation, which requires input for specific choices of the texture variable distribution, $f_T(T; \boldsymbol{\theta})$. A set of candidate distributions for the univariate texture is reviewed in the next chapter. Explicit expressions for the Mellin kind statistics of the speckle distribution, $f_{\widetilde{\mathbf{W}}}(\widetilde{\mathbf{W}}; L, \boldsymbol{\Sigma})$ are also given, such that Equations (3.40)–(3.42) can be evaluated.

Chapter 4

Probability Distributions

This chapter presents specific distributions for multilook polarimetric radar data, and evaluate their Mellin kind statistics. The mathematical expressions contain many special functions, that are defined initially. We then review a number of univariate distributions that are useful models for the texture random variable. Expressions are given for their PDF, Mellin kind characteristic function, and Mellin kind cumulants (log-cumulants). These have been derived previously in [Nicolas, 2006]. The scaled complex Wishart distribution used to model speckle is then presented, together with its Mellin kind statistics, that have been derived in Paper 2 of Chapter 6. Finally, we present the compound distributions modelling the texture modulated speckle of the polarimetric covariance matrix. The Mellin kind statistics of the texture distributions and the speckle distribution are combined into novel expressions for the multilook polarimetric case.

4.1 Special Functions

The expression of the PDFs, Mellin kind characteristic functions, moments and cumulants presented in this chapter require a number of special functions. These are defined below.

Gamma Function: The gamma function, often credited explicitly to Euler by name, is a generalisation of the factorial function to non-integer numbers. It is defined for $z \in \mathbb{C}$ with $\Re\{z\} > 0$ by [Weisstein, 2010c]

$$\Gamma(z) = \int_0^{\infty} u^{z-1} e^{-u} du \quad (4.1)$$

and satisfies the relation

$$\Gamma(z + 1) = z\Gamma(z) . \quad (4.2)$$

Multivariate Gamma Function of the Complex Kind: The multivariate gamma function of the complex kind is a generalisation of the gamma function. It is defined for $s \in \mathbb{C}$ by [Goodman, 1963]

$$\Gamma_d(s) = \int_{\Omega_+} \exp(-\operatorname{tr}(\mathbf{Z})) |\mathbf{Z}|^{s-d} d\mathbf{Z} = \pi^{d(d-1)/2} \prod_{i=0}^{d-1} \Gamma(s-i) \quad (4.3)$$

where Ω_+ is the cone of positive definite complex Hermitian matrices, $\operatorname{tr}(\cdot)$ is the trace operator, and the matrix $\mathbf{Z} \in \Omega_+$ has dimensions $d \times d$. We also note that $\Gamma_d(z)$ is the matrix-variate Mellin transform of $\exp(\operatorname{tr}(\mathbf{Z}))$.

Polygamma Function: The polygamma function of order ν is the $\nu + 1$ th derivative of the logarithm of the gamma function. It is defined by [Weisstein, 2010e]

$$\begin{aligned} \psi^{(\nu)}(z) &= (-1)^{(\nu+1)} \int_0^\infty \frac{u^\nu e^{-zu}}{1-e^{-u}} du \\ &= \frac{d^\nu}{dz^\nu} \psi^{(0)} = \frac{d^{\nu+1}}{dz^{\nu+1}} \ln \Gamma(z) \end{aligned} \quad (4.4)$$

where $\psi^{(0)}(z) = \Gamma'(z)/\Gamma(z)$ is the digamma function. A recurrence relation is given by

$$\psi^{(\nu)}(z+1) = \psi^{(\nu)}(z) + (-1)^\nu \nu! z^{-(\nu+1)}. \quad (4.5)$$

Multivariate Polygamma Function of the Complex Kind: The multivariate polygamma function of the complex kind is introduced in this thesis as an intuitive and simplifying generalisation of the ordinary polygamma function (see the previous paragraph). It is defined as

$$\psi_d^{(\nu)}(z) = \sum_{i=0}^{d-1} \psi^{(\nu)}(z-i) = \frac{d^{\nu+1}}{dz^{\nu+1}} \ln \Gamma_d(z). \quad (4.6)$$

The multivariate digamma function satisfies (see derivations in Paper 2, Chapter 6)

$$\psi_d^{(0)}(z) = \frac{\Gamma'_d(z)}{\Gamma_d(z)}. \quad (4.7)$$

Beta Function: The beta function, also called the Euler integral of the first kind, is defined for $\{x, y\} \in \mathbb{C}$ with $\operatorname{Re}\{x\} > 0$ and $\operatorname{Re}\{y\} > 0$ by [Weisstein, 2010a]

$$B(x, y) = \int_0^1 u^{x-1} (1-u)^{y-1} du = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}. \quad (4.8)$$

Modified Bessel function of the Second Kind: The modified Bessel function of the second kind, also known as the Bessel K function, is defined for order $\nu \in \mathbb{R}$ and argument $z \in \mathbb{C}$ by [Weisstein, 2010d]

$$K_\nu(z) = \frac{\Gamma(\nu + 1/2)(2z)^\nu}{\sqrt{\pi}} \int_0^\infty \frac{\cos u}{(u^2 + z^2)^{\nu+1/2}} du. \quad (4.9)$$

Confluent Hypergeometric Function of the Second Kind: The confluent hypergeometric function of the second kind, also known as the Kummer U function, is defined for arguments $\{a, b, z\} \in \mathbb{C}$ with $\Re\{a\}, \Re\{z\} > 0$ by [Weisstein, 2010b]

$$U(a, b, z) = \frac{1}{\Gamma(a)} \int_0^\infty e^{-zu} u^{a-1} (1+u)^{b-a-1} du. \quad (4.10)$$

4.2 Texture Distributions

According to the discussion of texture in Section 2.3.5, a *texture distribution* describes a real random variable T , which is strictly positive and unit mean ($T \in \mathbb{R}^+$, $E\{T\} = 1$). We here present some candidate distributions and expressions derived from them with full parametrisation. That is, they contain a location parameter μ , which allows for a mean value $\mu \neq 1$. This makes it easier to recognise the distribution from the literature, while the normalised expressions used in the context of the product model are easily obtained by the substitution of $\mu = 1$.

The texture distributions presented in the following were all treated in [Nicolas, 2006], where the PDF, Mellin kind characteristic function, ν th-order moment and ν th-order log-cumulant were given for each of them. The parametrisation used here is slightly different from the one Nicolas uses, in insisting that μ should be identical to the mean for all distributions. Departures from Nicolas' parametrisation are remarked for the distributions concerned. The advantage of this approach is that all expressions for the normalised distribution are obtained simply by setting $\mu = 1$. The shape parameters α and λ remain unchanged. On the other hand, the parametrisation of Nicolas is more elegant when showing how the composite distributions (i.e., the Fisher-Snedecor, beta, and inverse beta distribution) are composed by Mellin correlation, convolution, and inverse convolution of gamma distributions and inverse gamma distributions (see [Nicolas, 2006] for details).

4.2.1 Gamma Distribution

A gamma distributed random variable is denoted $X \sim \gamma(\mu, \alpha)$ with location parameter $\mu > 0$ and shape parameter $\alpha > 0$. A gamma distributed texture variable is denoted $T \sim \bar{\gamma}(\alpha) = \gamma(1, \alpha)$, due to the normalisation of $\mu = 1$.

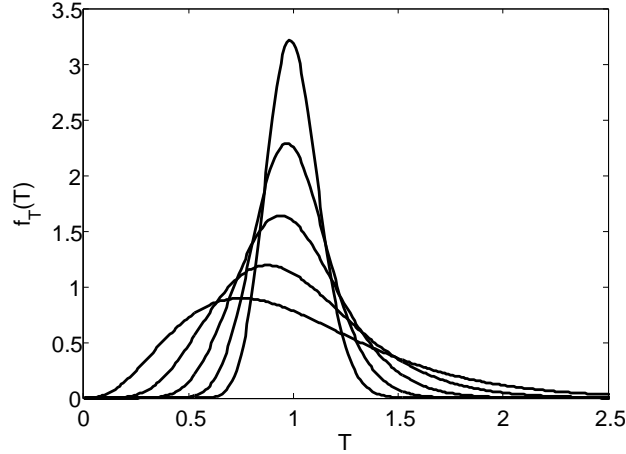


Figure 4.1: Probability density function $f_T(T; \alpha)$ of the unit mean gamma distributed texture variable T shown for shape parameters $\alpha = \{4, 8, 16, 32, 64\}$. The peakedness increases with α .

Probability density function:

$$f_X(x; \mu, \alpha) = \frac{1}{\Gamma(\alpha)} \left(\frac{\alpha}{\mu}\right) \left(\frac{\alpha x}{\mu}\right)^{\alpha-1} e^{-\frac{\alpha x}{\mu}}; \quad x \geq 0. \quad (4.11)$$

Mellin kind characteristic function:

$$\phi_X(s; \mu, \alpha) = \left(\frac{\mu}{\alpha}\right)^{s-1} \frac{\Gamma(\alpha + s - 1)}{\Gamma(\alpha)}. \quad (4.12)$$

Moments:

$$m_\nu\{X; \mu, \alpha\} = \left(\frac{\mu}{\alpha}\right)^\nu \frac{\Gamma(\alpha + \nu)}{\Gamma(\alpha)}. \quad (4.13)$$

Log-cumulants:

$$\kappa_\nu\{X; \mu, \alpha\} = \begin{cases} \psi^{(0)}(\alpha) + \ln\left(\frac{\mu}{\alpha}\right) & \text{for } \nu = 1 \\ \psi^{(\nu-1)}(\alpha) & \text{for } \nu > 1 \end{cases}. \quad (4.14)$$

4.2.2 Inverse Gamma Distribution

An inverse gamma distributed random variable is denoted $X \sim \gamma^{-1}(\mu, \lambda)$ with location parameter $\mu > 0$ and shape parameter $\lambda > 0$. An inverse gamma distributed texture variable is denoted $T \sim \bar{\gamma}^{-1}(\lambda) = \gamma^{-1}(1, \lambda)$.

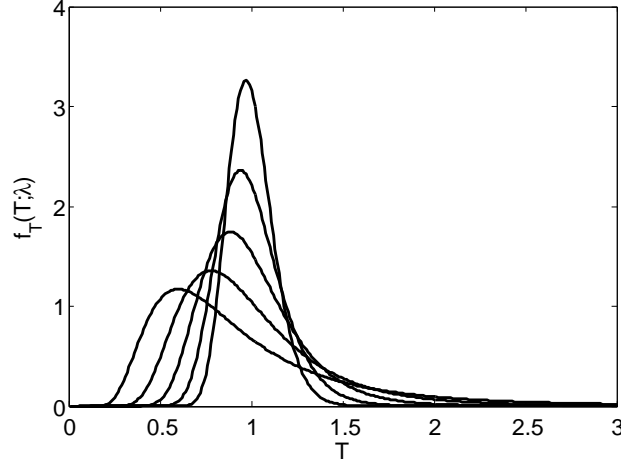


Figure 4.2: Probability density function $f_T(T; \lambda)$ of the unit mean inverse gamma distributed texture variable T shown for shape parameters $\lambda = \{4, 8, 16, 32, 64\}$. The peakedness increases with λ .

Probability density function:

$$f_X(x; \mu, \lambda) = \frac{1}{\Gamma(\lambda)} \frac{1}{(\lambda - 1)\mu} \left(\frac{(\lambda - 1)\mu}{x} \right)^{\lambda+1} e^{-\frac{(\lambda-1)\mu}{x}}; \quad x > 0. \quad (4.15)$$

Mellin kind characteristic function:

$$\phi_X(s; \mu, \lambda) = ((\lambda - 1)\mu)^{s-1} \frac{\Gamma(\lambda + 1 - s)}{\Gamma(\lambda)}. \quad (4.16)$$

Moments:

$$m_\nu\{X; \mu, \lambda\} = ((\lambda - 1)\mu)^\nu \frac{\Gamma(\lambda - \nu)}{\Gamma(\lambda)}. \quad (4.17)$$

Log-cumulants:

$$\kappa_\nu\{X; \mu, \lambda\} = \begin{cases} -\psi^{(0)}(\lambda) + \ln((\lambda - 1)\mu) & \text{for } \nu = 1 \\ (-1)^\nu \psi^{(\nu-1)}(\lambda) & \text{for } \nu > 1 \end{cases}. \quad (4.18)$$

The parametrisation of [Nicolas, 2006] is obtained by substituting $\mu = \left(\frac{\lambda}{\lambda-1}\right) \mu'$.

4.2.3 Fisher-Snedecor Distribution

A Fisher-Snedecor distributed random variable is denoted $X \sim \mathcal{F}(\mu, \alpha, \lambda)$ with location parameter $\mu > 0$ and shape parameters $\alpha > 0$ and $\lambda > 0$. A Fisher-Snedecor distributed texture variable is denoted $T \sim \bar{\mathcal{F}}(\alpha, \lambda) = \mathcal{F}(1, \alpha, \lambda)$.

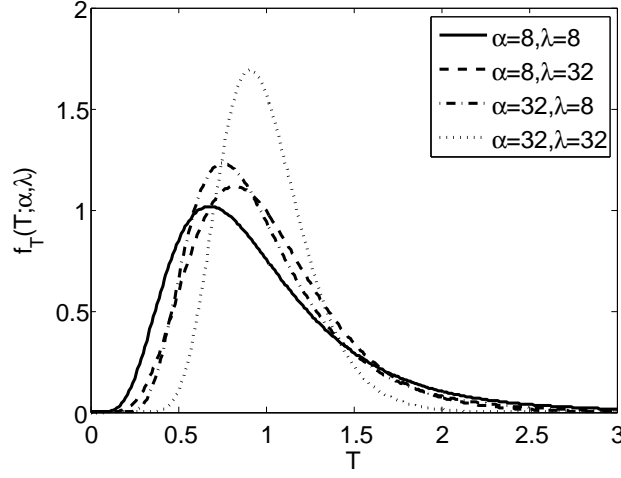


Figure 4.3: Probability density function $f_T(T; \alpha, \lambda)$ of the unit mean Fisher-Snedecor distributed texture variable T shown for different pairs of shape parameters (α, λ) .

Probability density function:

$$f_X(x; \mu, \alpha, \lambda) = \frac{1}{B(\alpha, \lambda)} \frac{\alpha}{(\lambda - 1)\mu} \frac{\left(\frac{\alpha x}{(\lambda - 1)\mu}\right)^{\alpha - 1}}{\left(\frac{\alpha x}{(\lambda - 1)\mu} + 1\right)^{\alpha + \lambda}}; \quad x \geq 0. \quad (4.19)$$

Mellin kind characteristic function:

$$\phi_X(s; \mu, \alpha, \lambda) = \left(\frac{(\lambda - 1)\mu}{\alpha}\right)^{s - 1} \frac{\Gamma(\alpha + s - 1)}{\Gamma(\alpha)} \frac{\Gamma(\lambda + 1 - s)}{\Gamma(\lambda)}. \quad (4.20)$$

Moments:

$$m_\nu\{X; \mu, \alpha, \lambda\} = \left(\frac{(\lambda - 1)\mu}{\alpha}\right)^\nu \frac{\Gamma(\alpha + \nu)}{\Gamma(\alpha)} \frac{\Gamma(\lambda - \nu)}{\Gamma(\lambda)}. \quad (4.21)$$

Log-cumulants:

$$\kappa_\nu\{X; \mu, \alpha, \lambda\} = \begin{cases} \psi^{(0)}(\alpha) - \psi^{(0)}(\lambda) + \ln\left(\frac{(\lambda - 1)\mu}{\alpha}\right) & \text{for } \nu = 1 \\ \psi^{(\nu - 1)}(\alpha) + (-1)^\nu \psi^{(\nu - 1)}(\lambda) & \text{for } \nu > 1 \end{cases}. \quad (4.22)$$

The parametrisation of [Nicolas, 2006] is obtained by substituting $\mu = \left(\frac{\lambda}{\lambda - 1}\right) \mu'$. The Fisher-Snedecor distribution is also called Fisher distribution. It can be seen as a generalised F distribution with an additional location parameter.

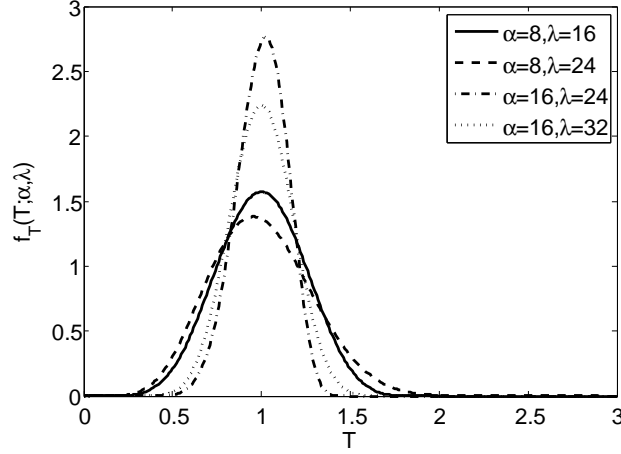


Figure 4.4: Probability density function $f_T(T; \alpha, \lambda)$ of the unit mean beta distributed texture variable T shown for different pairs of shape parameters (α, λ) .

4.2.4 Beta Distribution

A beta distributed random variable is denoted $X \sim \beta(\mu, \alpha, \lambda)$ with location parameter $\mu > 0$ and shape parameters $\alpha > 0$ and $\lambda > 0$. A beta distributed texture variable is denoted $T \sim \bar{\beta}(\alpha, \lambda) = \beta(1, \alpha, \lambda)$.

Probability density function:

$$f_X(x; \mu, \alpha, \lambda) = \frac{1}{B(\alpha, \lambda - \alpha)} \frac{\alpha}{\lambda \mu} \left(\frac{\alpha x}{\lambda \mu} \right)^{\alpha-1} \left(1 - \frac{\alpha x}{\lambda \mu} \right)^{\lambda-\alpha-1}; \quad 0 \leq x \leq \frac{\lambda \mu}{\alpha}. \quad (4.23)$$

Mellin kind characteristic function:

$$\phi_X(s; \mu, \alpha, \lambda) = \left(\frac{\lambda \mu}{\alpha} \right)^{s-1} \frac{\Gamma(\alpha + s - 1)}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda + s - 1)}. \quad (4.24)$$

Moments:

$$m_\nu\{X; \mu, \alpha, \lambda\} = \left(\frac{\lambda \mu}{\alpha} \right)^\nu \frac{\Gamma(\alpha + \nu)}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda + \nu)}. \quad (4.25)$$

Log-cumulants:

$$\kappa_\nu\{X; \mu, \alpha, \lambda\} = \begin{cases} \psi^{(0)}(\alpha) - \psi^{(0)}(\lambda) + \ln \left(\frac{\lambda \mu}{\alpha} \right) & \text{for } \nu = 1 \\ \psi^{(\nu-1)}(\alpha) - \psi^{(\nu-1)}(\lambda) & \text{for } \nu > 1 \end{cases}. \quad (4.26)$$

The beta distribution is often defined with two shape parameters only. The definition given here is a generalisation of the two parameter version, with an additional location parameter.

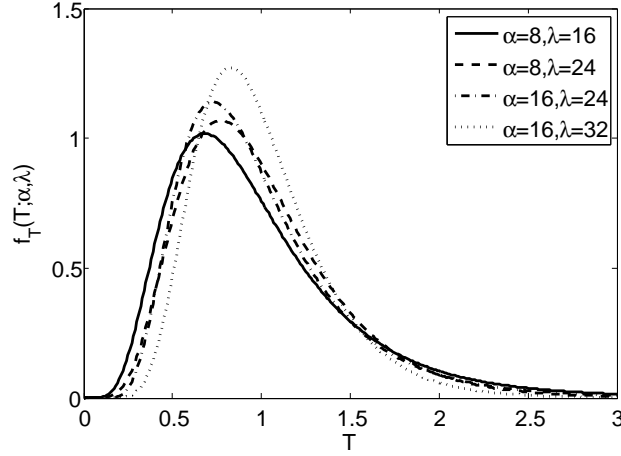


Figure 4.5: Probability density function $f_T(T; \alpha, \lambda)$ of the unit mean inverse beta distributed texture variable T shown for different pairs of shape parameters (α, λ) .

4.2.5 Inverse Beta Distribution

An inverse beta distributed random variable is denoted $X \sim \beta^{-1}(\mu, \alpha, \lambda)$ with location parameter $\mu > 0$ and shape parameters $\alpha > 0$ and $\lambda > 0$. An inverse beta distributed texture variable is denoted $T \sim \tilde{\beta}^{-1}(\alpha, \lambda) = \beta^{-1}(1, \alpha, \lambda)$.

Probability density function:

$$f_X(x; \mu, \alpha, \lambda) = \frac{1}{B(\alpha, \lambda - \alpha)} \frac{(\lambda - 1)}{(\alpha - 1)\mu} \frac{\left(\frac{(\lambda-1)x}{(\alpha-1)\mu} - 1\right)^{\lambda-\alpha-1}}{\left(\frac{(\lambda-1)x}{(\alpha-1)\mu}\right)^\lambda}; \quad x \geq \frac{(\alpha - 1)\mu}{\lambda - 1}, \quad \lambda \geq \alpha. \quad (4.27)$$

Mellin kind characteristic function:

$$\phi_X(s; \mu, \alpha, \lambda) = \left(\frac{(\alpha - 1)\mu}{\lambda - 1}\right)^{s-1} \frac{\Gamma(\alpha + 1 - s)}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda + 1 - s)}. \quad (4.28)$$

Moments:

$$m_\nu\{X; \mu, \alpha, \lambda\} = \left(\frac{(\alpha - 1)\mu}{\lambda - 1}\right)^\nu \frac{\Gamma(\alpha - \nu)}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda - \nu)}. \quad (4.29)$$

Log-cumulants:

$$\kappa_\nu\{X; \mu, \alpha, \lambda\} = \begin{cases} -\psi^{(0)}(\alpha) + \psi^{(0)}(\lambda) + \ln\left(\frac{(\alpha-1)\mu}{(\lambda-1)}\right) & \text{for } \nu = 1 \\ (-1)^\nu \psi^{(\nu-1)}(\alpha) + (-1)^{\nu-1} \psi^{(\nu-1)}(\lambda) & \text{for } \nu > 1 \end{cases}. \quad (4.30)$$

The parametrisation of [Nicolas, 2006] is obtained by substituting $\mu = \left(\frac{\alpha-1}{\alpha}\right) \left(\frac{\lambda}{\lambda-1}\right) \mu'$. The inverse beta distribution is also known as the beta distribution of the second kind, or the beta prime distribution. It is often defined with two shape parameters only. The definition given here is a generalisation of the two parameter version, with an additional location parameter.

4.3 Speckle Distribution

The term *speckle distribution* is here used to denote a distribution which models the randomness of the radar signal due to the interference phenomenon only. That is, it describes pure speckle. When it is compounded with a texture distribution, we obtain under the doubly stochastic product model a distribution which describes texture modulated speckle. The Weibull and log-normal distributions for single polarisation amplitude and intensity data [Oliver and Quegan, 2004] are examples of distributions that describe texture modulated speckle without explicitly modelling the texture through a texture variable. Matrix-variate counterparts of these distributions have not been derived, as far as the author is aware of. There are no known distributions that describe partially developed speckle in the multilook polarimetric case either. We are left with the scaled complex Wishart distribution, which models pure and fully developed speckle, but will also present the matrix-variate Mellin kind statistics of the true complex Wishart distribution.

4.3.1 Complex Wishart Distribution

The complex extension of the Wishart distribution, defined on Ω_+ , was first reported in [Goodman, 1963]. The distribution describes $\mathbf{W} = L\widetilde{\mathbf{W}}$ and the complex Wishart distributed matrix is denoted $\mathbf{W} \sim \mathcal{W}_d^c(L, \Sigma)$, where the equivalent number of looks, L , is a shape parameter and $\Sigma = \text{E}\{\mathbf{W}\}/L$ is a scale matrix. The nonsingular distribution is obtained for $L \geq d$.

Probability density function:

$$f_{\mathbf{W}}(\mathbf{W}; L, \Sigma) = \frac{|\mathbf{C}|^{L-d}}{\Gamma_d(L)|\Sigma|^L} \text{etr}(-\Sigma^{-1}\mathbf{C}). \quad (4.31)$$

Mellin kind characteristic function:

$$\phi_{\mathbf{W}}(s; L, \Sigma) = \frac{\Gamma_d(L+s-d)}{\Gamma_d(L)} |\Sigma|^{s-d}. \quad (4.32)$$

Mellin kind cumulant generating function:

$$\varphi_{\mathbf{W}}(s; L, \Sigma) = \ln \Gamma_d(L+s-d) - \ln \Gamma_d(L) + (s-d) \ln |\Sigma|. \quad (4.33)$$

Matrix log-cumulants:

$$\kappa_\nu\{W; L, \Sigma\} = \begin{cases} \psi_d^{(0)}(L) + \ln |\Sigma| & \text{for } \nu = 1, \\ \psi_d^{(\nu-1)}(L) & \text{for } \nu > 1. \end{cases} \quad (4.34)$$

The derivation of the Mellin kind statistics is shown in Paper 2 of Chapter 6.

4.3.2 Scaled Complex Wishart Distribution

The observable covariance matrix in pure and fully developed speckle is $\widetilde{\mathbf{W}}$, which follows the scaled complex Wishart distribution. This is denoted $\widetilde{\mathbf{W}} \sim s\mathcal{W}_d^{\mathbb{C}}(L, \Sigma)$. The PDF is derived from (4.31) as $f_{\widetilde{\mathbf{W}}}(\widetilde{\mathbf{W}}) = f_{\mathbf{W}}(L\widetilde{\mathbf{W}})|J_{\mathbf{W} \rightarrow \widetilde{\mathbf{W}}}|$ by using the transformation $\widetilde{\mathbf{W}} = \mathbf{W}/L$ with Jacobian determinant $|J_{\mathbf{W} \rightarrow \widetilde{\mathbf{W}}}| = L^{d^2}$.

Probability density function:

$$f_{\mathbf{C}}(\mathbf{C}; L, \Sigma) = \frac{L^{Ld}}{\Gamma_d(L)} \frac{|\mathbf{C}|^{L-d}}{|\Sigma|^L} \text{etr}(-L\Sigma^{-1}\mathbf{C}). \quad (4.35)$$

Mellin kind characteristic function:

$$\phi_{\mathbf{W}}(s; L, \Sigma) = \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} \left(\frac{|\Sigma|}{L^d} \right)^{s-d}. \quad (4.36)$$

Mellin kind cumulant generating function:

$$\varphi_{\mathbf{W}}(s; L, \Sigma) = \ln \Gamma_d(L + s - d) - \ln \Gamma_d(L) + (s - d)(\ln |\Sigma| - d \ln L). \quad (4.37)$$

Matrix log-cumulants:

$$\kappa_\nu\{W; L, \Sigma\} = \begin{cases} \psi_d^{(0)}(L) + \ln |\Sigma| - d \ln L & \text{for } \nu = 1, \\ \psi_d^{(\nu-1)}(L) & \text{for } \nu > 1. \end{cases} \quad (4.38)$$

The derivation of the Mellin kind statistics is shown in Paper 2 of Chapter 6.

4.4 Compound Matrix Distributions

Compound matrix distributions that model texture modulated speckle are obtained from the Mellin convolution in (2.26). Five distributions are reported, that combine the scaled complex Wishart distribution with each of the texture distribution presented in Section 4.2 through the multilook polarimetric product model. For all of them, the PDF is listed together with the matrix-variate Mellin kind characteristic function, cumulant generating function and cumulants. The Mellin kind statistics are derived by evaluating Equations (3.40)–(3.42) with the expressions given in Sections 4.2 and 4.3.2.

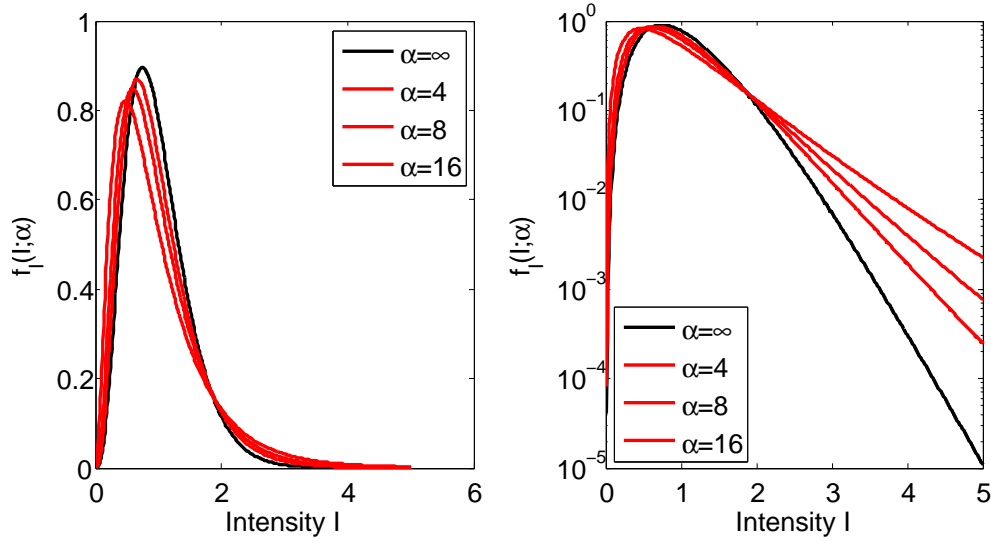


Figure 4.6: PDF $f_I(I; \alpha)$ of the \mathcal{K} distributed multilook intensity I shown for shape parameters $\alpha = \{4, 6, 8, \infty\}$ at linear (left) and logarithmic (right) scale. The asymptotic case $\alpha = \infty$ is equivalent to gamma distributed intensity, corresponding to a Wishart distributed polarimetric covariance matrix \mathbf{C} . The other parameters are $L = 4$ and $\langle I \rangle = 1$.

4.4.1 Matrix-Variate \mathcal{K} Distribution

The matrix-variate \mathcal{K} distribution was derived in [Lee et al., 1994]. The name originates from the Bessel K function, which appears in the PDF expression. It results from a Mellin convolution of the scaled complex Wishart distribution with the normalised (unit mean) gamma distribution. A \mathcal{K} distributed covariance matrix is denoted $\mathbf{C} \sim \mathcal{K}(L, \Sigma, \alpha)$, where the parameters are inherited from the speckle and texture distribution.

Probability density function:

$$f_{\mathbf{C}}(\mathbf{C}; L, \Sigma, \alpha) = \frac{|\mathbf{C}|^{L-d} 2(L\alpha)^{\frac{\alpha+Ld}{2}}}{|\Sigma|^L \Gamma_d(L)\Gamma(\alpha)} (\text{tr}(\Sigma^{-1}\mathbf{C}))^{\frac{\alpha-Ld}{2}} K_{\alpha-Ld}(2\sqrt{L\alpha \text{tr}(\Sigma^{-1}\mathbf{C})}). \quad (4.39)$$

Mellin kind characteristic function:

$$\phi_{\mathbf{C}}(s; L, \Sigma, \alpha) = \frac{\Gamma_d(L+s-d)}{\Gamma_d(L)} \frac{\Gamma(\alpha+d(s-d))}{\Gamma(\alpha)} \left(\frac{|\Sigma|}{(\alpha L)^d} \right)^{s-d}. \quad (4.40)$$

Mellin kind cumulant generating function:

$$\varphi_{\mathbf{C}}(s; L, \Sigma, \alpha) = \ln \frac{\Gamma_d(L+s-d)}{\Gamma_d(L)} + \ln \frac{\Gamma(\alpha+d(s-d))}{\Gamma(\alpha)} + (s-d)(\ln |\Sigma| - d \ln(\alpha L)). \quad (4.41)$$

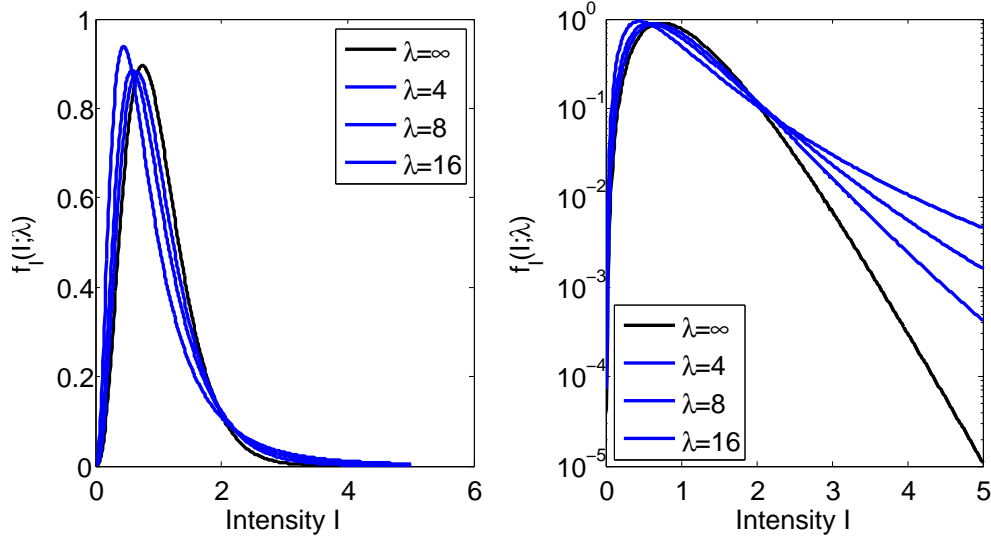


Figure 4.7: PDF $f_I(I; \lambda)$ of the \mathcal{G}^0 distributed multilook intensity I shown for shape parameters $\lambda = \{4, 6, 8, \infty\}$ at linear (left) and logarithmic (right) scale. The asymptotic case $\lambda = \infty$ is equivalent to gamma distributed intensity, corresponding to a Wishart distributed polarimetric covariance matrix \mathbf{C} . The other parameters are $L=4$ and $\langle I \rangle = 1$.

Matrix log-cumulants:

$$\kappa_\nu\{\mathbf{C}; L, \Sigma, \alpha\} = \begin{cases} \psi_d^{(0)}(L) + \ln |\Sigma| + d \left(\psi^{(0)}(\alpha) - \ln(\alpha L) \right) & \text{for } \nu = 1, \\ \psi_d^{(\nu-1)}(L) + d^\nu \psi^{(\nu-1)}(\alpha) & \text{for } \nu > 1. \end{cases} \quad (4.42)$$

Matrix-variate distributions are difficult to visualise, and the solution is to plot instead the the marginal density of the single polarisation intensities found as diagonal elements in the polarimetric covariance matrix. Figure 4.6 shows examples of univariate \mathcal{K} distributions of intensity for different values of the texture parameter α .

4.4.2 Matrix-Variate \mathcal{G}^0 Distribution

The matrix-variate \mathcal{G}^0 distribution was derived in [Freitas et al., 2005]. It results from a Mellin convolution of the scaled complex Wishart distribution with the normalised inverse gamma distribution. A \mathcal{G}^0 distributed covariance matrix is denoted $\mathbf{C} \sim \mathcal{G}^0(L, \Sigma, \lambda)$, where the parameters are inherited from the speckle and texture distribution.

Probability density function:

$$f_{\mathbf{C}}(\mathbf{C}; L, \Sigma) = \frac{L^{Ld}}{\Gamma_d(L)} \frac{|\mathbf{C}|^{L-d} \Gamma(Ld + \lambda) (\lambda - 1)^\lambda}{|\Sigma|^L \Gamma(\lambda)} (L \operatorname{tr}(\Sigma^{-1} \mathbf{C}) + \lambda - 1)^{-\lambda - Ld}. \quad (4.43)$$

Mellin kind characteristic function:

$$\phi_{\mathbf{C}}(s; L, \Sigma, \lambda) = \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} \frac{\Gamma(\lambda - d(s - d))}{\Gamma(\lambda)} \left(|\Sigma| \left(\frac{\lambda - 1}{L} \right)^d \right)^{s-d}. \quad (4.44)$$

Mellin kind cumulant generating function:

$$\begin{aligned} \varphi_{\mathbf{C}}(s; L, \Sigma, \lambda) = & \ln \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} + \ln \frac{\Gamma(\lambda - d(s - d))}{\Gamma(\lambda)} \\ & + (s - d) \left(\ln |\Sigma| + d(\ln(\lambda - 1) - \ln L) \right). \end{aligned} \quad (4.45)$$

Matrix log-cumulants:

$$\kappa_{\nu}\{\mathbf{C}; L, \Sigma, \lambda\} = \begin{cases} \psi_d^{(0)}(L) + \ln |\Sigma| + d \left(\ln \left(\frac{\lambda - 1}{L} \right) - \psi^{(0)}(\lambda) \right) & \text{for } \nu = 1, \\ \psi_d^{(\nu-1)}(L) + (-d)^{\nu} \psi^{(\nu-1)}(\lambda) & \text{for } \nu > 1. \end{cases} \quad (4.46)$$

Figure 4.7 shows examples of univariate \mathcal{G}^0 distributions of intensity for different values of the texture parameter λ .

4.4.3 Matrix-Variate \mathcal{U} Distribution

The matrix-variate \mathcal{U} distribution was derived in [Bombrun and Beaulieu, 2008]. The name originates from the Kummer U function, which appears in the PDF expression. It results from a Mellin convolution of the scaled complex Wishart distribution with the normalised Fisher-Snedecor distribution. A \mathcal{U} distributed covariance matrix is denoted $\mathbf{C} \sim \mathcal{U}(L, \Sigma, \alpha, \lambda)$, where the parameters are inherited from the speckle and texture distribution.

Probability density function:

$$\begin{aligned} f_{\mathbf{C}}(\mathbf{C}; L, \Sigma, \alpha, \lambda) = & \frac{L^{Ld}}{\Gamma_d(L)} \frac{|\mathbf{C}|^{L-d}}{|\Sigma|^L} \frac{\Gamma(\alpha + \lambda) \Gamma(Ld + \lambda)}{\Gamma(\alpha) \Gamma(\lambda)} \left(\frac{\alpha}{\lambda - 1} \right) \\ & \times U \left(Ld + \lambda, Ld - \alpha + 1, L \operatorname{tr}(\Sigma^{-1} \mathbf{C}) \frac{\alpha}{(\lambda - 1)} \right). \end{aligned} \quad (4.47)$$

Mellin kind characteristic function:

$$\begin{aligned} \phi_{\mathbf{C}}(s; L, \Sigma, \alpha, \lambda) = & \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} \frac{\Gamma(\alpha + d(s - d))}{\Gamma(\alpha)} \frac{\Gamma(\lambda - d(s - d))}{\Gamma(\lambda)} \\ & \times \left(|\Sigma| \left(\frac{\lambda - 1}{\alpha L} \right)^d \right)^{s-d}. \end{aligned} \quad (4.48)$$

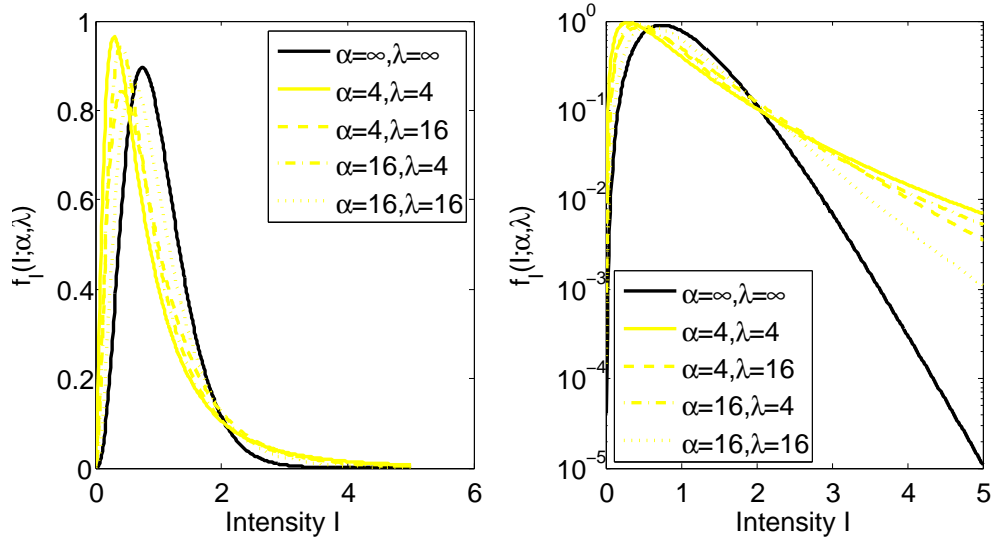


Figure 4.8: PDF $f_I(I; \lambda)$ of the \mathcal{U} distributed multilook intensity I shown for different pairs of shape parameters, (α, λ) , at linear (left) and logarithmic (right) scale. The asymptotic case $(\alpha = \infty, \lambda = \infty)$ is equivalent to gamma distributed I , corresponding to a scaled Wishart distributed \mathbf{C} . The asymptotic cases $\alpha \rightarrow \infty$ and $\lambda \rightarrow \infty$ correspond to \mathcal{K} and \mathcal{G}^0 distributed I and \mathbf{C} , respectively. The other parameters are $L = 4$ and $\langle I \rangle = 1$.

Mellin kind cumulant generating function:

$$\begin{aligned} \varphi_{\mathbf{C}}(s; L, \Sigma, \alpha, \lambda) = & \ln \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} + \ln \frac{\Gamma(\alpha + d(s - d))}{\Gamma(\alpha)} + \ln \frac{\Gamma(\lambda - d(s - d))}{\Gamma(\lambda)} \\ & + (s - d) \left(\ln |\Sigma| + d(\ln(\lambda - 1) - \ln \alpha - \ln L) \right). \end{aligned} \quad (4.49)$$

Matrix log-cumulants:

$$\begin{aligned} \kappa_1\{\mathbf{C}; L, \Sigma, \alpha, \lambda\} = & \psi_d^{(0)}(L) + \ln |\Sigma| + d \left(\psi^{(0)}(\alpha) - \psi^{(0)}(\lambda) + \ln \left(\frac{\lambda - 1}{\alpha L} \right) \right), \\ \kappa_{\nu > 1}\{\mathbf{C}; L, \Sigma, \alpha, \lambda\} = & \psi_d^{(\nu-1)}(L) + d^\nu \left(\psi^{(\nu-1)}(\alpha) + (-1)^\nu \psi^{(\nu-1)}(\lambda) \right). \end{aligned} \quad (4.50)$$

Figure 4.8 shows examples of univariate \mathcal{U} distributions of intensity for different pairs of the texture parameters α and λ .

4.4.4 Matrix-Variate \mathcal{W} Distribution

The PDF of the matrix-variate \mathcal{W} distribution has not yet been derived, but it is found that it will contain a special function known as the Whittaker \mathcal{W} function, hence the proposed name. The distribution results from a Mellin convolution of the scaled complex

Wishart distribution with the normalised beta distribution, and the Mellin statistics are easily deduced from the existing results. A \mathcal{W} distributed covariance matrix is denoted $\mathbf{C} \sim \mathcal{W}(L, \mathbf{\Sigma}, \alpha, \lambda)$, where the parameters are inherited from the speckle and texture distribution.

Mellin kind characteristic function:

$$\phi_{\mathbf{C}}(s; L, \mathbf{\Sigma}, \alpha, \lambda) = \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} \frac{\Gamma(\alpha + d(s - d))}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda + d(s - d))} \left(|\mathbf{\Sigma}| \left(\frac{\lambda}{\alpha L} \right)^d \right)^{s-d}. \quad (4.51)$$

Mellin kind cumulant generating function:

$$\varphi_{\mathbf{C}}(s; L, \mathbf{\Sigma}, \alpha, \lambda) = \ln \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} + \ln \frac{\Gamma(\alpha + d(s - d))}{\Gamma(\alpha)} + \ln \frac{\Gamma(\lambda)}{\Gamma(\lambda + d(s - d))} + (s - d) \left(\ln |\mathbf{\Sigma}| + d(\ln \lambda - \ln \alpha - \ln L) \right). \quad (4.52)$$

Matrix log-cumulants:

$$\begin{aligned} \kappa_1\{\mathbf{C}; L, \mathbf{\Sigma}, \alpha, \lambda\} &= \psi_d^{(0)}(L) + \ln |\mathbf{\Sigma}| + d \left(\psi^{(0)}(\alpha) - \psi^{(0)}(\lambda) + \ln \left(\frac{\lambda}{\alpha L} \right) \right) \\ \kappa_{\nu>1}\{\mathbf{C}; L, \mathbf{\Sigma}, \alpha, \lambda\} &= \psi_d^{(\nu-1)}(L) + d^\nu \left(\psi^{(\nu-1)}(\alpha) - \psi^{(\nu-1)}(\lambda) \right). \end{aligned} \quad (4.53)$$

4.4.5 Matrix-Variate \mathcal{M} Distribution

The PDF of the matrix-variate \mathcal{M} distribution has not yet been derived, but it is found that it will contain a special function known as the Whittaker \mathcal{M} function, hence the proposed name. The distribution results from a Mellin convolution of the scaled complex Wishart distribution with the normalised inverse beta distribution, and the Mellin statistics are easily deduced from the existing results. A \mathcal{M} distributed covariance matrix is denoted $\mathbf{C} \sim \mathcal{M}(L, \mathbf{\Sigma}, \alpha, \lambda)$, where the parameters are inherited from the speckle and texture distribution.

Mellin kind characteristic function:

$$\phi_{\mathbf{C}}(s; L, \mathbf{\Sigma}, \alpha, \lambda) = \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} \frac{\Gamma(\alpha - d(s - d))}{\Gamma(\alpha)} \frac{\Gamma(\lambda)}{\Gamma(\lambda - d(s - d))} \left(|\mathbf{\Sigma}| \left(\frac{(\alpha - 1)}{(\lambda - 1)L} \right)^d \right)^{s-d}. \quad (4.54)$$

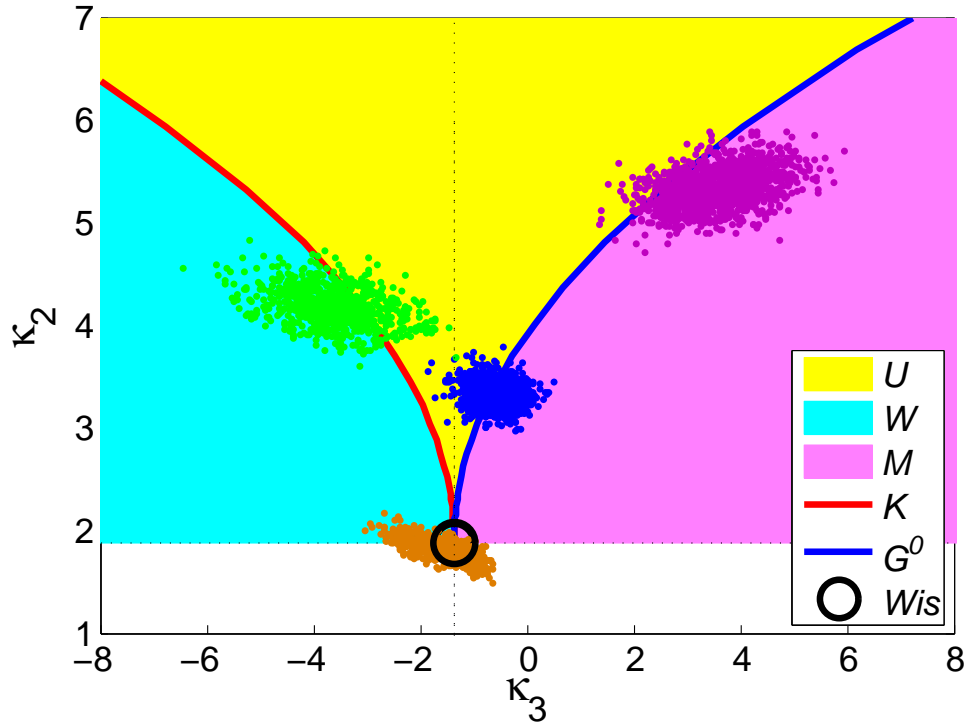


Figure 4.9: The matrix log-cumulant diagram

Mellin kind cumulant generating function:

$$\begin{aligned} \varphi_{\mathbf{C}}(s; L, \Sigma, \alpha, \lambda) = & \ln \frac{\Gamma_d(L + s - d)}{\Gamma_d(L)} + \ln \frac{\Gamma(\alpha - d(s - d))}{\Gamma(\alpha)} + \ln \frac{\Gamma(\lambda)}{\Gamma(\lambda - d(s - d))} \\ & + (s - d) \left(\ln |\Sigma| + d(\ln(\alpha - 1) - \ln(\lambda - 1) - \ln L) \right). \end{aligned} \quad (4.55)$$

Matrix log-cumulants:

$$\begin{aligned} \kappa_1\{\mathbf{C}; L, \Sigma, \alpha, \lambda\} = & \psi_d^{(0)}(L) + \ln |\Sigma| + d \left(\psi^{(0)}(\lambda) - \psi^{(0)}(\alpha) + \ln \left(\frac{(\alpha - 1)}{(\lambda - 1)L} \right) \right) \\ \kappa_{\nu > 1}\{\mathbf{C}; L, \Sigma, \alpha, \lambda\} = & \psi_d^{(\nu-1)}(L) + (-d)^\nu \left(\psi^{(\nu-1)}(\alpha) - \psi^{(\nu-1)}(\lambda) \right). \end{aligned} \quad (4.56)$$

4.5 The Matrix Log-Cumulant Diagram

We now introduce the matrix log-cumulant diagram, which is shown in Figure 4.9. This is a visualisation tool which allows us to both characterise the different matrix distributions and compare data to the distribution models. The diagram promotes intuition about the applications of the Mellin kind statistics framework, and is used in Papers 2

and 3 (Chapters 6 and 7) to provide geometrical interpretations of the proposed estimation procedures and goodness-of-fit tests.

In the diagram we plot: (i) manifolds spanned by the theoretical matrix log-cumulants that can be attained under given models, and (ii) points that represent the empirical sample matrix log-cumulants computed from data samples. The dimension of a given manifold equals the number of parameters in the texture distribution associated with that particular model. The Wishart distribution has no texture parameters, and is represented by a black point (zero-dimensional manifold). The matrix-variate \mathcal{K} and \mathcal{G}^0 distributions have one texture parameter, and are represented by the red and blue curves (one-dimensional manifolds), respectively. The \mathcal{U} , \mathcal{W} and \mathcal{M} distributions have two texture parameters, and are visualised in the diagram as the respective yellow, cyan and magenta surfaces.

The diagram is spanned by the second and third-order matrix log-cumulants, and is a direct extension of the log-cumulants proposed in [Nicolas, 2002, Nicolas, 2006] for the univariate case. It was shown in (4.38) that under the multilook polarimetric product model, matrix log-cumulants of order $\nu \geq 2$ are independent of the speckle distribution scale matrix Σ . They depend only on L , which is considered a constant, and the texture parameters. By plotting the third-order matrix log-cumulant against the second-order matrix log-cumulant (a convention introduced by Nicolas), we obtain a diagram which shows the solitary impact of the texture parameters upon the models. Thus, it promotes insight about how we can select between the different compound matrix distributions (including the Wishart distribution). We can also assess the fit between data and models by looking the distances between them in matrix log-cumulant space. Finally, the diagram visualises how texture parameters are estimated from data by projecting sample matrix log-cumulants onto the manifolds of population matrix log-cumulants.

The clusters of sample matrix log-cumulants plotted in Figure 4.9 represent targets that have been selected from different images acquired by the airborne AIRSAR L-band sensor previously operated by NASA/JPL. The samples represent forest (green cluster), sea (blue cluster), urban area (purple cluster) and cropland (brown cluster), and indicate which model is suitable in each case. The multiple points in each cluster were obtained by bootstrap sampling the covariance matrix samples. This is a way of showing the variance of the sample matrix log-cumulants. Note also that the matrix distributions presented in this chapter cover the whole matrix log-cumulant plane, remembering that the lower semi-plane represents texture distributions with negative variances.

Chapter 5

Paper 1:

Estimation of the Equivalent Number of Looks in Polarimetric Synthetic Aperture Radar Imagery

Chapter 6

Paper 2:

Application of the Matrix-Variate Mellin Transform to Analysis of Polarimetric Radar Images

Chapter 7

Paper 3:

Goodness-of-Fit Tests for Multilook
Polarimetric Radar Data Based on the
Mellin Transform

Chapter 8

Conclusions

This chapter gives concluding remarks and outlines some directions of future research.

8.1 Concluding Remarks

In this thesis, it has been shown that the Mellin kind statistics framework is a natural tool for analysis of the matrix distributions derived under the multilook polarimetric product model. The simple and elegant mathematical expression we obtain, the performance of the parameter estimators, and the interpretability of the results are taken as proof of this statement.

Some excellent parameter estimators have been derived by the method of matrix log-cumulants, and a pioneering goodness-of-fit test has been constructed. However, by looking at the long list of applications that have grown out of Nicolas' univariate Mellin kind statistics (See the second paragraph of Section 1.1), it can be expected that more algorithms for analysis of multilook polarimetric images will follow. We have already identified some target areas and applications.

8.2 Future Research

The first paper where the contours of the matrix-variate Mellin kind statistics framework could be seen was [Anfinsen et al., 2009]. This paper has been included in Appendix A. It discusses statistical modelling of speckle filtered multilook polarimetric images and demonstrates that the filtering alters the data such that new models are required. This is a topic which is worth pursuing. By looking at the characteristics of the speckle filtered data in matrix log-cumulant space, the matrix-variate \mathcal{W} and \mathcal{M} distributions are launched as potential models for filtered speckle in heterogeneous and extremely heterogeneous areas. Another task would be to find statistical models for partially developed speckle, and the Mellin kind statistics could possibly be helpful in the characterisation of such a model.

An interesting observation is that the log-determinant transformation of the polarimetric covariance matrix compresses the matrix data to a scalar band. By storing this band only, we can still compute matrix log-moments and matrix log-cumulants and make inferences based on them. This compression could be utilised to make fast algorithms for different image analysis tasks that still maintain high performance, since the matrix log-cumulants have a high content of statistical information.

It was mentioned in Section 2.3.5 that the texture parameters of the radar image distributions can be used as textural features in various image analysis problems. The matrix log-cumulants would contain the same information, but are cheaper to compute and not confined to a specific distribution model. This is an area that we want to explore further. Compound matrix distributions have already been applied to classification [Doulgeris et al., 2008] and segmentation [Bombrun and Beaulieu, 2008, Harant et al., 2009, Vasile et al., 2009]. In polarimetric change detection on the other hand, the test statistics proposed in [Conradsen et al., 2003, Kersten et al., 2005] are developed under the assumption of Wishart statistics. Extensions could possibly be achieved by means of Mellin kind statistics.

As a final remark, we note again that the log-determinant compression induced by the matrix-variate Mellin transform reduces the data dimension to one. The statistical information contained in the polarimetric covariance matrix could possibly be better preserved by using the multivariate Mellin transform from [Mathur and Krishna, 1977], even though the off-diagonal complex correlations will be discarded in such an approach. Thus, there is a chance we could use the Mellin transform to squeeze even more information out of polarimetric radar images.

Appendix A

A Relaxed Wishart Model for Polarimetric SAR Data

A RELAXED WISHART MODEL FOR POLARIMETRIC SAR DATA

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ABSTRACT

In this paper we demonstrate that simple yet flexible modelling of multilook polarimetric synthetic aperture radar (PolSAR) data can be obtained through a relaxation of the Wishart model. The degrees of freedom of the complex Wishart distribution is treated as a spatially nonstationary parameter, which is allowed to vary between thematic classes and segments of the PolSAR scene.

Key words: synthetic aperture radar; polarimetry; statistical modelling; Wishart distribution.

1. INTRODUCTION

The Wishart distribution is the de facto statistical model for multilook PolSAR data. It is based on the assumption that the complex scattering coefficients are jointly circular Gaussian. However, this is only satisfied for homogeneous areas with fully developed speckle and no texture, which renders the model inadequate in many cases. Improved modelling is achieved by using more complex models that account for texture, such as the polarimetric \mathcal{G} distribution family [1], with the polarimetric \mathcal{K} distribution [2] as a special case. These models allow for better adaption to data whose distribution is heavy-tailed and non-Gaussian, but this comes at the cost of higher mathematical complexity.

The comparatively higher mathematical tractability of the Wishart distribution motivates us to pursue a relaxed Wishart model as an alternative. In the context of multilook PolSAR data, the degrees of freedom of the Wishart distribution is interpreted as the equivalent number of looks, a constant, global value that quantifies the effective number of data samples averaged in the multilooking process. In contrast, we treat it as a free parameter, which varies between, and possibly also within, classes and segments of the PolSAR scene. This reflects the highly variable degree of smoothing imposed on the data by nonlinear speckle filters. The choice can also be justified by looking at the degrees of freedom as a shape parameter of the distribution, which is determined not only by the degree of averaging, but also by texture. Thus, the influence of multilooking, speckle filtering, and texture is

assimilated into one parameter, which can be estimated efficiently with a recently proposed estimator [3, 4]

Sec. 2 reviews some existing density models for multilook PolSAR data and proposes the relaxed Wishart distribution as an alternative. In Sec. 3 we derive certain matrix moments that are used to illustrate the adaptivity of the different density models, and as a new domain for visual goodness-of-fit assessment. Sec. 4 presents experiments with airborne NASA/JPL AIRSAR data, and in Sec. 5 we give our conclusions.

2. STATISTICAL MODELLING

The full-polarimetric SAR instrument separately transmits orthogonally polarised microwave pulses, and measures orthogonal components of the received signal. For each pixel, the measurements result in a matrix of scattering coefficients. These are complex-valued, dimensionless numbers that describe the transformation of the transmitted (incoming) electromagnetic (EM) field to the received (backscattered) EM field for all combinations of transmit and receive polarisation.

The transformation can be expressed as

$$\begin{bmatrix} E_h^r \\ E_v^r \end{bmatrix} = \frac{e^{jkr}}{r} \begin{bmatrix} S_{hh} & S_{hv} \\ S_{vh} & S_{vv} \end{bmatrix} \begin{bmatrix} E_h^t \\ E_v^t \end{bmatrix}, \quad (1)$$

where k denotes wavenumber and r is the distance between radar and target. The subscript of the EM field components E_i^j denotes horizontal (h) or vertical (v) polarisation, which is the most common set of orthogonal polarisations, while the superscript indicates transmitted (t) or received (r) wave. The scattering coefficients S_{ij} are subscripted with the associated receive and transmit polarisation, in that order. Together, they form the scattering matrix, denoted $\mathbf{S} = [S_{ij}]$.

The scattering matrix can be reduced to one of the vectors

$$\mathbf{s} = \begin{bmatrix} S_{hh} \\ (S_{hv} + S_{vh})/\sqrt{2} \\ S_{vv} \end{bmatrix} \text{ or } \mathbf{k} = \frac{1}{\sqrt{2}} \begin{bmatrix} S_{hh} + S_{vv} \\ S_{hh} - S_{vv} \\ S_{hv} + S_{vh} \end{bmatrix}. \quad (2)$$

The lexicographic scattering vector, denoted \mathbf{s} , is the vectorised version of \mathbf{S} after the cross-polarisation terms S_{hv}

and S_{vh} have been averaged, assuming reciprocity of the target. The scaling with a factor $\sqrt{2}$ is done to preserve total power of the signal. The Pauli basis scattering vector, denoted \mathbf{k} , is a linear transformation of \mathbf{s} , which provides physical interpretations of its elements in terms of basic scattering mechanisms [5].

2.1. Gaussian Model

It is commonly assumed that the scattering vector elements are jointly circular complex Gaussian, even though this model only encompasses variability due to speckle, and not texture, which is discussed in the Sec. 2.2. The matrix \mathbf{S} and the vectors \mathbf{s} and \mathbf{k} are single-look complex format representations of PolSAR data. The following derivations shall use \mathbf{s} as the scattering vector, but would be equivalent for \mathbf{k} .

Multilook PolSAR data is commonly represented by

$$\mathbf{C} = \frac{1}{L} \sum_{i=1}^L \mathbf{s}_i \mathbf{s}_i^H, \quad (3)$$

known as the sample covariance matrix. It is formed as the mean Hermitian outer product of the single-look scattering vectors $\{\mathbf{s}_i\}_{i=1}^L$, where L is the nominal number of looks. The superscript H means complex conjugate transpose. Assume that \mathbf{s} is zero mean and circular complex multivariate Gaussian, denoted $\mathbf{s} \sim \mathcal{N}_d^{\mathbb{C}}(\mathbf{0}, \mathbf{\Sigma})$, where $\mathbf{0}$ is a column vector of zeros, d is the dimension of \mathbf{s} , and $\mathbf{\Sigma} = \mathbb{E}\{\mathbf{s}\mathbf{s}^H\}$ is the covariance matrix of \mathbf{s} . The probability density function (pdf) of \mathbf{s} is thus

$$p_{\mathbf{s}}(\mathbf{s}; \mathbf{\Sigma}) = \frac{1}{\pi^d |\mathbf{\Sigma}|} \exp(-\mathbf{s}^H \mathbf{\Sigma}^{-1} \mathbf{s}), \quad (4)$$

where $|\cdot|$ is the determinant operator. It follows that if $L \geq d$ and the \mathbf{s}_i in (3) are independent, then \mathbf{C} follows the nonsingular complex Wishart distribution [6]:

$$p_{\mathbf{w}}(\mathbf{C}; L, \mathbf{\Sigma}) = \frac{L^{Ld} |\mathbf{C}|^{L-d}}{|\mathbf{\Sigma}|^L \Gamma_d(L)} \exp(-L \operatorname{tr}(\mathbf{\Sigma}^{-1} \mathbf{C})), \quad (5)$$

where $\operatorname{tr}(\cdot)$ is the trace operator. The normalisation constant $\Gamma_d(L)$ is the multivariate Gamma function, defined as

$$\Gamma_d(L) = \pi^{d(d-1)/2} \prod_{i=0}^{d-1} \Gamma(L-i), \quad (6)$$

where $\Gamma(L)$ is the standard Euler gamma function. In reality, the \mathbf{s}_i are correlated, and this is compensated for by replacing L with an equivalent number of looks, $L_e \geq L$, in order to obtain consistency between moments of the theoretical model and sample moments of the data. This approximation provides a good model for the distribution of \mathbf{C} , denoted $\mathbf{C} \sim \mathcal{W}_d^{\mathbb{C}}(L_e, \mathbf{\Sigma})$.

2.2. Product Model

In addition to speckle, the randomness of a SAR measurement can also be attributed to texture. The notion of texture represents the natural spatial variation of the radar cross section, which is generally not perfectly homogeneous for pixels that are thematically mapped as one class. Several statistical models exist that incorporate texture, either by assuming statistics that imply a non-Gaussian scattering vector, or explicitly modelling texture as a separate random variable (rv). The latter case leads to a doubly stochastic model with a compounded distribution.

The well known product model, reviewed in [7, 8, 9], is shown to be both mathematically tractable and successful for modelling purposes. In the multilook polarimetric version [1], the polarimetric covariance matrix \mathbf{C} is decomposed as a product of two independent stochastic processes with individual distributions:

$$\mathbf{C} = z \mathbf{W}. \quad (7)$$

One process, $\mathbf{W} \sim \mathcal{W}_d^{\mathbb{C}}(L_e, \mathbf{\Sigma})$, models speckle. The other process generates texture, represented by the scalar rv $z \in \mathbb{R}^+$, under the assumption that texture is independent of polarisation. The pdf of \mathbf{C} depends on the distribution of z , which is normalised to unit mean.

Gamma Distributed Texture

The first covariance matrix distribution derived from the product model used the gamma distribution to model z [2]. A gamma distributed rv $z > 0$ has density

$$p_z(z; \alpha, \mu) = \left(\frac{\alpha}{\mu}\right)^{\alpha} \frac{z^{\alpha-1}}{\Gamma(\alpha)} \exp\left(-\frac{\alpha}{\mu} z\right), \quad (8)$$

with shape parameter $\alpha > 0$ and mean value $\mu = \mathbb{E}\{z\} > 0$. This is denoted $z \sim \mathcal{G}(\alpha, \mu)$. The unitary mean texture rv is thus $z \sim \mathcal{G}(\alpha, 1)$. This leads to the matrix-variate \mathcal{K} distribution for \mathbf{C} [1, 2]:

$$p_{\mathbf{C}}(\mathbf{C}; L_e, \mathbf{\Sigma}, \alpha) = \frac{2 |\mathbf{C}|^{L_e-d} (L_e \alpha)^{\frac{\alpha+L_e d}{2}}}{|\mathbf{\Sigma}|^{L_e} \Gamma_d(L_e) \Gamma(\alpha)} (\operatorname{tr}(\mathbf{\Sigma}^{-1} \mathbf{C}))^{\frac{\alpha-L_e d}{2}} \times K_{\alpha-L_e d}(2\sqrt{L_e \alpha \operatorname{tr}(\mathbf{\Sigma}^{-1} \mathbf{C})}). \quad (9)$$

$K_n(\cdot)$ is the modified Bessel function of the second kind with order n . See [1, 10] for a detailed derivation.

Inverse Gamma Distributed Texture

The family of generalised inverse Gaussian distributions was proposed in [1] as a model for z . The gamma distribution is one special case. The inverse gamma distribution is another, which has been promoted in particular as

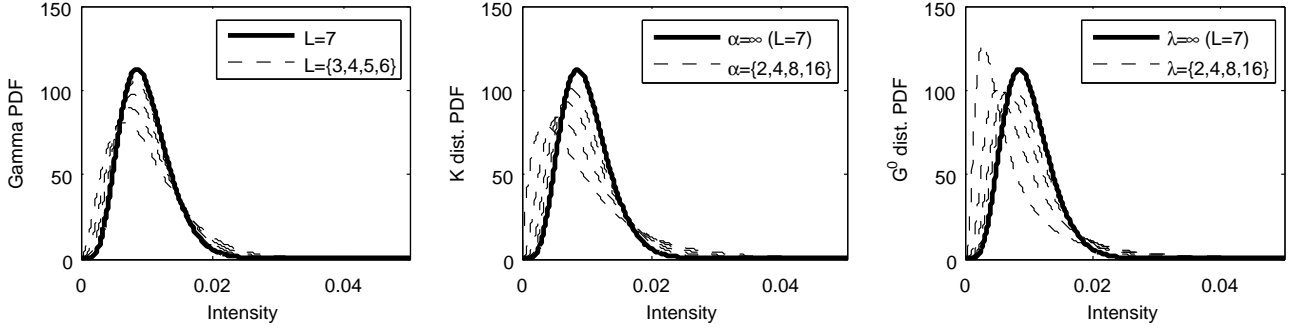


Figure 1. Examples of single intensity marginal densities for a polarimetric covariance matrix modelled by the relaxed Wishart distribution (left), matrix-variate \mathcal{K} distribution (middle), and matrix-variate \mathcal{G}^0 distribution (right).

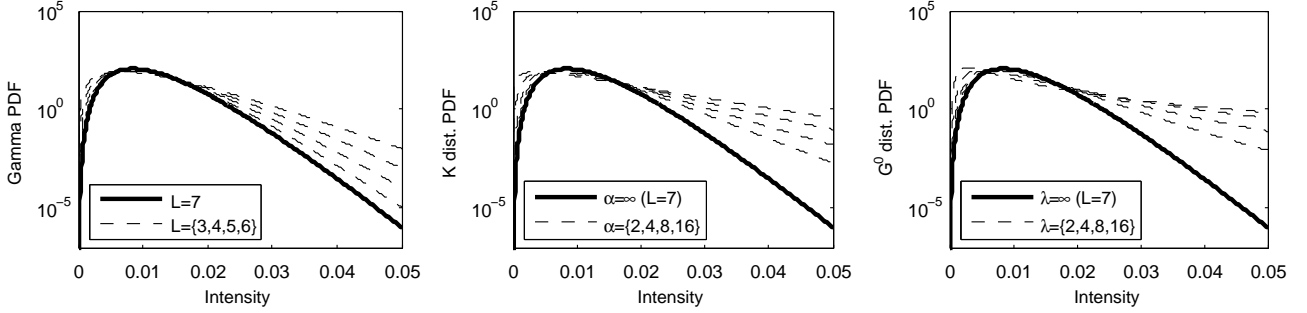


Figure 2. Same as Fig. 1 with logarithmic second axis to emphasize differences at the tails.

a good model for strongly heterogeneous clutter [1, 11]. Its pdf is given by

$$p_z(z; \lambda, \nu) = (\lambda\nu)^\lambda \frac{z^{-\lambda-1}}{\Gamma(\lambda)} \exp\left(-\frac{\lambda\nu}{z}\right). \quad (10)$$

This is denoted $z \sim \mathcal{G}^{-1}(\lambda, \nu)$, with shape parameter $\lambda > 0$ and $\nu > 0$. The normalised texture rv becomes $z \sim \mathcal{G}^{-1}(\lambda, (\lambda-1)/\lambda)$, which leads to the matrix-variate \mathcal{G}^0 distribution for \mathbf{C} [1]:

$$\begin{aligned} p_{\mathbf{C}}(\mathbf{C}; L_e, \boldsymbol{\Sigma}, \lambda) \\ = \frac{L_e^{L_e d} |\mathbf{C}|^{L_e-d} \Gamma(L_e d + \lambda) (\lambda-1)^\lambda}{|\boldsymbol{\Sigma}|^{L_e} \Gamma_d(L_e) \Gamma(\lambda)} \\ \times (L_e \text{tr}(\boldsymbol{\Sigma}^{-1} \mathbf{C}) + \lambda - 1)^{-\lambda-L_e d}. \end{aligned} \quad (11)$$

For interpretation purposes, we note that $z \rightarrow 1$ and the distributions in Eqs. (9) and (11) converge in distribution to the complex Wishart distribution in Eq. (5) as $\alpha \rightarrow \infty$ and $\lambda \rightarrow \infty$, respectively. Thus, high values of α and λ imply little texture, whereas low values refer to significant texture and non-Gaussianity.

2.3. Relaxed Wishart Model

The standard Wishart model in Eq. (5) is parametrised by a constant L_e , which is estimated for the data set as a whole [3, 4]. We introduce a relaxed Wishart (\mathcal{RW}) model, whose functional form is identical. The difference is that L_e is replaced with a variable shape param-

eter, $\mathcal{L} \leq L_e$. Depending on the application, \mathcal{L} is allowed to vary between classes (classification), segments (segmentation), or pixels (e.g., change detection). The new distribution is denoted by $\mathbf{C} \sim \mathcal{RW}_d^{\mathcal{C}}(\mathcal{L}, \boldsymbol{\Sigma})$.

The motivation for this approach is explained by Fig. 1. It is not possible to visualise the effect of the distribution parameters directly on the pdfs in Eqs. (5), (9), and (11). We therefore plot their marginal densities for a single polarisation intensity. The respective marginals are gamma distributed, \mathcal{K}_I distributed, and \mathcal{G}_I^0 distributed. For the latter two, the superscripted I denotes the multilook intensity version of the given distribution family.

In all the plots, the continuous curve represents the limiting case defined by the standard Wishart model, with a gamma distributed marginal pdf. In the left panel, the dashed curves show the evolution of the pdf under the \mathcal{RW} model as \mathcal{L} is lowered from the limit of $\mathcal{L} = L_e$. The same evolution is illustrated for the \mathcal{K} distribution (middle panel) and the \mathcal{G}^0 distribution (right panel) for decreasing values of the respective texture parameters, α and λ . We observe that the effect of varying \mathcal{L} resembles that induced by α and λ , even though a greater variation in shape is possible for the distributions based on the product model. Fig. 2 uses a logarithmic scale to highlight the heavy tails of the \mathcal{K} and \mathcal{G}^0 distribution, which is less prominent for the marginal pdf of the \mathcal{RW} distribution. We still conclude that \mathcal{L} can be interpreted as a texture parameter alongside α and λ . Thus, the \mathcal{RW} distribution implicitly models texture up to a moderate level.

3. GOODNESS-OF-FIT EVALUATION

This section discusses evaluation of the goodness-of-fit (GoF) for the matrix-variate density models of multilook PolSAR data. GoF testing in the literature has been limited to visual inspection of how well marginal densities of intensity fit histograms of the data. Classical statistical distribution tests, such as the Kolmogorov-Smirnov test or the Anderson-Darling test, are impractical in this case. As noted in [1], this is because they require binning of the domain of \mathbf{C} , which is the cone of positive definite matrices.

We here propose an alternative space where GoF evaluation can be performed. The idea is that GoF can be assessed by comparing theoretical moments of the models with sample moments computed from the data. We first define a new kind of matrix moments, that we call log-determinant cumulants. Closed form expressions for the candidate models are then derived. These are seen to have favourable properties that can be utilised to visualise the texture modelling capabilities of the models.

3.1. Log-determinant Cumulants

The following derivation is based on the application of second kind statistics, following the terminology introduced in [12]. Whereas the commonly known characteristic function is defined as the Fourier transform of a pdf, the second kind characteristic function is the Mellin transform of the pdf. This function can be used to generate moments and cumulants of the second kind, also termed log-moments and log-cumulants.

Let ξ be a real, positive rv with pdf $p_\xi(\xi)$. Start by defining the r th-order log-moment of ξ as

$$m_r(\xi) = \mathbb{E}\{(\ln \xi)^r\} = \left. \frac{d^r \phi_\xi(s)}{ds^r} \right|_{s=1}, \quad (12)$$

where $\phi_\xi(s)$ is the Mellin transform of $p_\xi(\xi)$ and $s \in \mathbb{C}$ [12]. Then define the r th-order log-cumulant of ξ as

$$\kappa_r(\xi) = \left. \frac{d^r \psi_\xi(s)}{ds^r} \right|_{s=1}, \quad (13)$$

where $\psi_\xi(s) = \ln \phi_\xi(s)$. Relations between some low-order log-moments and log-cumulants are given by

$$\kappa_1 = m_1, \quad (14)$$

$$\kappa_2 = m_2 - m_1^2, \quad (15)$$

$$\kappa_3 = m_3 - 3m_1m_2 + 2m_1^3. \quad (16)$$

It follows from a fundamental property of the Mellin transform [12] that for a product of independent random variables, $\xi = \rho \cdot \zeta$, with $\rho, \zeta > 0$:

$$\kappa_r(\xi) = \kappa_r(\rho) + \kappa_r(\zeta), \quad \forall r \in \mathbb{N}. \quad (17)$$

This equips us to derive the log-cumulant of $|\mathbf{C}|$, which will be referred to as the log-determinant cumulant (LDC) of \mathbf{C} .

Note that $|\mathbf{C}| = |z\mathbf{W}| = z^d |\mathbf{W}|$. Thus,

$$\kappa_r(|\mathbf{C}|) = d^r \kappa_r(z) + \kappa_r(|\mathbf{W}|). \quad (18)$$

The log-cumulants of z have been derived in [12]. For $z \sim \mathcal{G}(\alpha, \mu)$ it was shown that

$$\kappa_1(z) = \ln\left(\frac{\mu}{\alpha}\right) + \Psi^{(0)}(\alpha), \quad (19)$$

$$\kappa_r(z) = \Psi^{(r-1)}(\alpha), \quad r > 1, \quad (20)$$

where $\Psi^{(r)}(z)$ is Euler's polygamma function of order r . The log-cumulants of $z \sim \mathcal{G}^{-1}(\lambda, \nu)$ were found as

$$\kappa_1(z) = \ln \lambda \nu - \Psi^{(0)}(\lambda), \quad (21)$$

$$\kappa_r(z) = (-1)^r \Psi^{(r-1)}(\lambda). \quad (22)$$

The LDCs of the Wishart distributed \mathbf{W} can be deduced from results found in [3, 4] as

$$\kappa_1(|\mathbf{W}|) = \ln |\boldsymbol{\Sigma}| + \sum_{i=0}^{d-1} \Psi^{(0)}(L_e - i) - d \ln L_e \quad (23)$$

$$\kappa_r(|\mathbf{W}|) = \sum_{i=0}^{d-1} \Psi^{(r-1)}(L_e - i), \quad r > 1. \quad (24)$$

This completes the expression in Eq. (18) for our candidate models.

3.2. Log-determinant Cumulant Diagram

Note that the LDCs are matrix-variate generalisation of the log-cumulants derived in [12] for the single polarisation product model. As in the one-dimensional case, we can utilise the fact the LDCs do not depend on the scale parameter $\boldsymbol{\Sigma}$ for $r > 1$. More specifically, the $\kappa_{r>1}(|\mathbf{C}|)$ depend only on the texture parameters:

$$\kappa_{r>1}^{\mathcal{W}}(|\mathbf{C}|) = \sum_{i=0}^{d-1} \Psi^{(r-1)}(L_e - i) \quad (25)$$

$$\kappa_{r>1}^{\mathcal{K}}(|\mathbf{C}|) = d^r \Psi^{(r-1)}(\alpha) + \sum_{i=0}^{d-1} \Psi^{(r-1)}(L_e - i) \quad (26)$$

$$\kappa_{r>1}^{\mathcal{G}^0}(|\mathbf{C}|) = (-d)^r \Psi^{(r-1)}(\lambda) + \sum_{i=0}^{d-1} \Psi^{(r-1)}(L_e - i) \quad (27)$$

$$\kappa_{r>1}^{\mathcal{RW}}(|\mathbf{C}|) = \sum_{i=0}^{d-1} \Psi^{(r-1)}(\mathcal{L} - i), \quad (28)$$

where the superscript of κ indicates which model the expression describes.

By plotting two LDCs of different orders against each other, we obtain a curve in LDC space which depicts the

paired LDC values that can be attained under a given model. We refer to this as an LDC diagram. Sample LDCs calculated from data can be overlaid the model curves, and the diagram used to assess how well the data are described by different models, and which model provides the best fit. Diagrams of second and third order log-cumulants were plotted in [12, 13], and we will use the same orders for our LDC diagrams. Remark that the bias and variance of the sample LDCs are expected to increase rapidly with order.

4. EXPERIMENTS

4.1. Marginal Densities of Intensity

We have extracted three test samples from an L-band quadrature polarisation image acquired by the airborne NASA/JPL AIRSAR sensor over Flevoland, the Netherlands, in August 1989. The samples are taken from some of the more textured areas in the image. One is from a forest area and the other two from different crops. Marginal densities of the intensity in the HH, HV, and VV channels for the forest sample are shown in Fig. 3. These densities describe unfiltered data. Fig. 4 describes the same data after they have been filtered with a refined Lee filter [14] of window size 7×7 .

The first observation is that the data are well described by all the models in Fig. 3. By zooming in on the densities, it may be concluded by visual inspection that the \mathcal{K} model provides the best fit, followed by the \mathcal{RW} model. The ENL estimated for the data set, and used to parametrise the standard Wishart, \mathcal{K} , and \mathcal{G}^0 model, is $L_e = 3.3$. This constant is replaced with $\mathcal{L} = 2.53$ for the \mathcal{RW} model. The texture parameters α and λ are estimated with the method described for the \mathcal{K} distribution in [10].

Fig. 4 shows that the models have very different GoF for speckle filtered data. The Wishart model is the worst fit, and none of the distributions based on the product model produce an adequate result either. Only the \mathcal{RW} model seems to do a good job. The ENL was estimated to $L_e = 48$, and is reduced to $\mathcal{L} = 27$ for the \mathcal{RW} model. The marginal densities of the other two test samples yield very similar results, both before and after speckle filtering, and are therefore not shown.

4.2. Log-determinant Cumulant Diagrams

Fig. 5 shows a LDC diagram where $\kappa_3(|\mathbf{C}|)$ is plotted against $\kappa_2(|\mathbf{C}|)$, with analogy to the log-cumulant diagrams in [12, 13]. The Wishart model has no texture parameter, and its LDCs are therefore constant, equal to the contribution $\kappa_r(|\mathbf{W}|)$, $r = 1, 2$ of the Wishart distributed speckle matrix. These constants are indicated in the figure by the dotted lines, intersecting at the point which describes the Wishart model. The possible LDC pairs of the \mathcal{K} , \mathcal{G}^0 , and \mathcal{RW} models lie on a curve parametrised by

α , λ , and \mathcal{L} , respectively. The asymptotic behaviour of these curves, as the texture parameters decrease towards their lower limits, is indicated on the figure. At the upper limit, the curves all converge to the Wishart case. Sample LDCs of the three test samples are plotted as points in green (forest), cyan, and magenta (two different crops). We plot a collection of sample LDC estimates, obtained by bootstrap sampling of the full test samples, in order to illustrate the dispersion of the sample estimates. ENL estimates for each test sample are shown in the figure.

Diagrams of data before and after speckle filtering are presented in the left and right panel, respectively. For the unfiltered data, the LDC diagram clearly indicates that the \mathcal{RW} distribution is the best model for the forest sample. The other test samples are less textured, and all models are adequate. For the speckle filtered data, the LDC diagram suggests that the \mathcal{RW} model fits best for the forest sample and the first crop sample (cyan), while the \mathcal{K} model performs best for the second crop sample (magenta). For the crop samples, both the \mathcal{RW} model and the \mathcal{K} model fit reasonably well. The Wishart and the \mathcal{G}^0 model are inadequate in all cases. The LDC diagram indicates good fit of the \mathcal{K} model to the crop samples, which is not compatible with observations of the marginal densities (not shown). This prompts us to reconsider the estimator for α (and λ) in future work.

The success of the \mathcal{RW} distribution in modelling of speckle filtered data, and the relative failure of the others, can be explained by a discussion of the nature of adaptive speckle filters. Adaptive speckle filters apply variable smoothing by consideration of local homogeneity. Hence, the ENL is mapped from a constant value to a dispersed range of values. This is not modelled appropriately, neither by the Wishart distribution nor the other distributions based on the product model. The \mathcal{RW} distribution, on the other hand, apparently represents a better approach.

5. CONCLUSIONS

We have proposed a relaxed Wishart distribution where the equivalent number of looks of the standard Wishart model has been replaced by a variable shape parameter. We have further derived the log-determinant cumulants of the polarimetric covariance (or coherency) matrix under the product model, and demonstrated how they can be utilised in visual inspection of goodness-of-fit for matrix-variate distributions. Experimental results show that for a moderate level of texture, the newly proposed density can compete with densities derived from the product model with regards to modelling of unfiltered PolSAR data. For data that are processed with an adaptive speckle filter, the relaxed Wishart model is shown to perform better. Based on the very promising results, we suggest that the relaxed Wishart distribution should be tested more extensively on other data sets and with different speckle filters. It should also be applied to model-based classification, change detection, and other image analysis tasks.

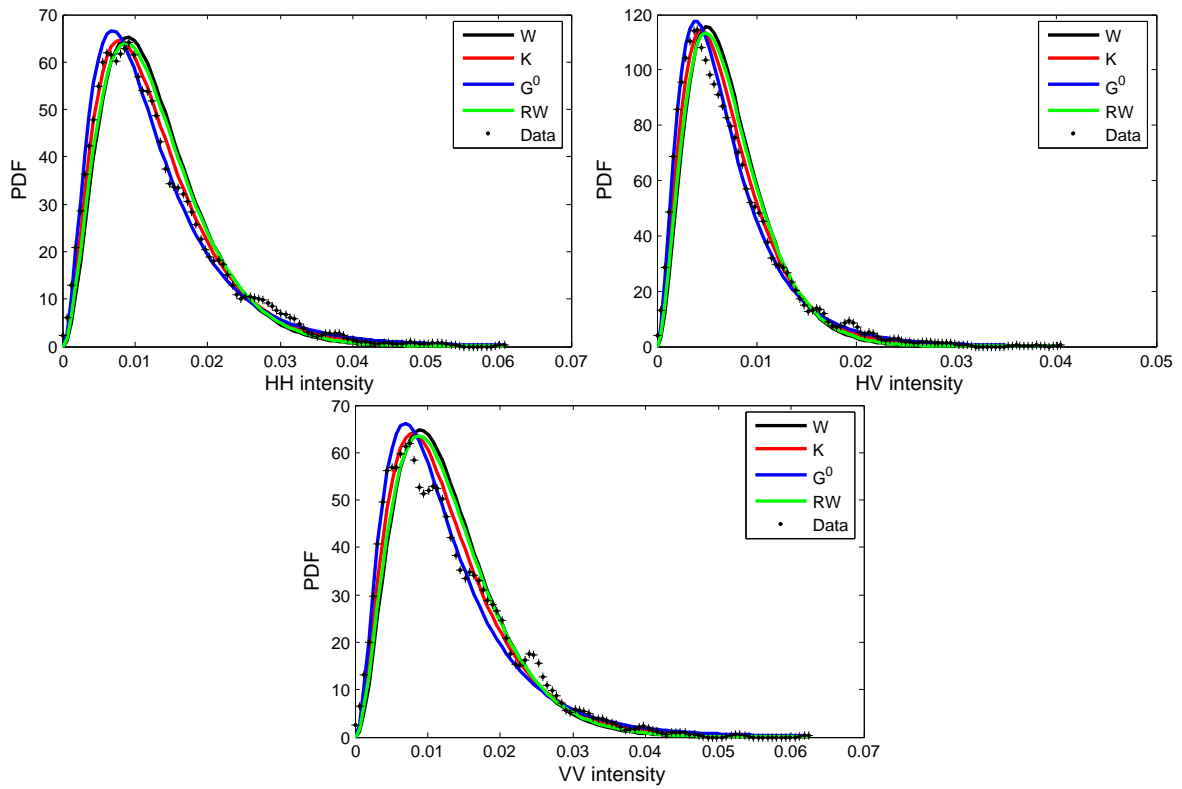


Figure 3. Comparison of marginal densities of the Wishart, \mathcal{RW} , \mathcal{K} , and \mathcal{G}^0 distribution with data histograms for a textured forest area in the AIRSAR L-band image of Flevoland. No speckle filter applied.

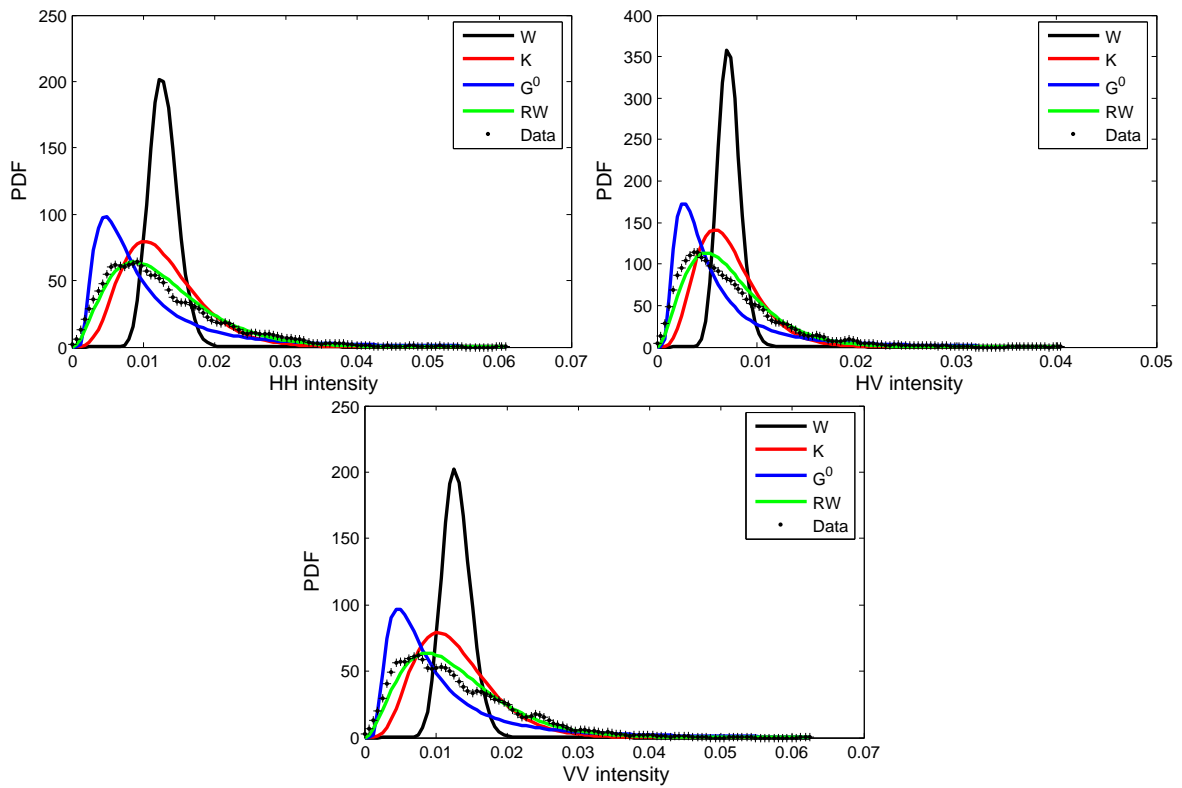


Figure 4. Comparison of marginal densities of the Wishart, \mathcal{RW} , \mathcal{K} , and \mathcal{G}^0 distribution with data histograms for a textured forest area in the AIRSAR L-band image of Flevoland. Modified Lee filter with window size 7×7 applied.

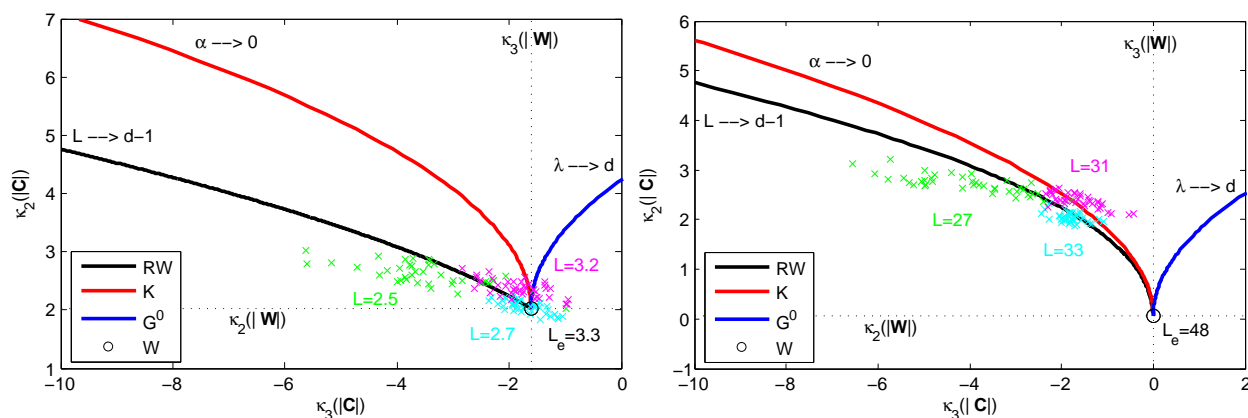


Figure 5. Diagram of second and third-order log-determinant cumulants (LDCs) for the Wishart, \mathcal{RW} , \mathcal{K} , and \mathcal{G}^0 distribution. Sample LDCs for three test samples from the AIRSAR L-band image of Flevoland are plotted in green (forest), cyan and magenta (two different crops). Results are shown before (left panel) and after (right panel) speckle filtering with a 7×7 refined Lee filter.

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Appendix B

Introduction to Second Kind Statistics:
Application of Log-Moments and
Log-Cumulants to Analysis of Radar
Images

Introduction to Second Kind Statistics: Application of Log-Moments and Log-Cumulants to Analysis of Radar Images

Jean-Marie Nicolas and Stian Normann Anfinsen (translator)



Abstract—Statistical methods classically used to analyse a probability density function (pdf) are founded on the Fourier transform, on which useful tools such the first and second characteristic function are based, yielding the definitions of moments and cumulants. Yet this transformation is badly adapted to the analysis of probability density functions defined on \mathbb{R}^+ , for which the analytic expressions of the characteristic functions may become hard, or even impossible to formulate. In this article we propose to substitute the Fourier transform with the Mellin transform. It is then possible, inspired by the precedent definitions, to introduce *second kind statistics*: second kind characteristic functions, second kind moments (or log-moments), and second kind cumulants (or log-cumulants). Applied to traditional distributions like the gamma distribution or the Nakagami distribution, this approach gives results that are easier to apply than the classical approach. Moreover, for more complicated distributions, like the \mathcal{K} distributions or the positive α -stable distributions, the second kind statistics give expressions that are truly simple and easy to exploit. The new approach leads to innovative methods for estimating the parameters of distributions defined on \mathbb{R}^+ . It is possible to compare the estimators obtained with estimators based on maximum likelihood theory and the method of moments. One can thus show that the new methods have variances that are considerably lower than those mentioned, and slightly higher than the Cramér-Rao bound.

Index Terms—Probability density functions defined on \mathbb{R}^+ , gamma distribution, Nakagami distribution, characteristic functions, parameter estimation, Mellin transform

1 INTRODUCTION

ESTIMATION of the parameters of a probability density functions (pdf) is a topic of major significance in pattern recognition. Starting from these estimates, segmentation and classification algorithms can be implemented, both in the field of signal processing and image processing. In signal processing, the intrinsic knowledge of the nature of the data (provided by an acoustic sensor, electromagnetic sensor, etc.) allows us to make realistic

assumptions about the suitable distribution models. In particular, many techniques are based on the additive noise model with a noise term that is assumed to be Gaussian. Traditionally, if one describes a random phenomenon by a pdf, one will also introduce the concept of the characteristic function, defined as the Fourier transform \mathcal{F} of the pdf. For example, if $p_x(u)$ is the pdf modelling a random variable X , the characteristic function $\Phi_x(v)$ is obtained by the relation [1]:

$$\Phi_x(v) = \mathcal{F}[p(u)](v) = \int_{-\infty}^{+\infty} e^{jvu} p_x(u) du. \quad (1)$$

The second characteristic function is defined as the logarithm of the characteristic function:

$$\Phi_x(v) = \log(\Psi_x(v)). \quad (2)$$

By taking account of properties of the Fourier transform, it is easy to show that moments of order n are obtained by derivation of the characteristic function:

$$\begin{aligned} m_n &= \int_{-\infty}^{+\infty} u^n p_x(u) du \\ &= (-j)^n \left. \frac{d^n \Psi_x(v)}{dv^n} \right|_{v=0} \end{aligned} \quad (3)$$

and cumulants of order n by derivation of the second characteristic function:

$$\kappa_{x(r)} = (-j)^n \left. \frac{d^r \Psi_x(v)}{dv^r} \right|_{v=0}.$$

Moreover, if a phenomenon is analysed, described by a pdf q_y , which is perturbed by an additive noise, described by its pdf r_z , one knows that the output signal is described by the pdf p_x given as

$$p_x = q_y * r_z, \quad (4)$$

with the operator $*$ denoting convolution. It is known that the characteristic functions and the cumulants can be written:

$$\Phi_x(s) = \Phi_y(s) \Phi_z(s) \quad (5)$$

$$\Psi_x(s) = \Psi_y(s) + \Psi_z(s) \quad (6)$$

$$\kappa_{x(r)} = \kappa_{y(r)} + \kappa_{z(r)} \quad \forall r \quad (7)$$

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However, in image processing the problems are different. It should be noted first of all that the pixel values are positive or zero (we will not discuss here the analysis of images defined by complex values), and that the noise is often multiplicative. Also, the preceding model must undergo some adaptations to be applicable as it is. One approach often proposed is to perform a logarithmic transformation, which is possible since the pdf is defined on \mathbb{R}^+ . Several remarks can then be made:

- The analytical calculation of the characteristic functions defined on \mathbb{R}^+ is sometimes hard, even impossible for certain distributions, as we will see in section 3.2.
- No complete methodology is proposed for logarithmically transformed data. Calculation of moments on logarithmic scale (that one may call log-moments) is carried out analytically starting from Eq. (3). It requires a change of variable (thus a rewriting of the pdf for this new variable) and is carried out in a specific way for each pdf. This approach requires a good knowledge of integral transforms and of the properties of special functions.
- In traditional statistics, the Gaussian distribution is the reference, which corresponds to the log-normal distribution on a logarithmic scale. However, in many examples, this law does not describe the studied phenomenon well. In particular, the speckle (clutter) observed in images obtained by coherent illumination (e.g., laser, radar, ultrasound) follows, for intensity images, the gamma distribution, which we will study in more detail in this article and which tends asymptotically towards a degenerated Gaussian distribution.

As we will show, a new methodology based on another integral transformation, the Mellin transform [2], makes it possible to perform a more effective analysis of practically important distributions defined on \mathbb{R}^+ . This methodology, that we propose to call *second kind statistics*, uses the same framework as traditional statistics for the definition of the characteristic functions (simply by replacing the Fourier transform with the Mellin transform in Eq. (1)) and the same construction of the moments and cumulants (by derivation of the characteristic functions). This leads naturally to the definitions of *second kind moments* and *second kind cumulants*. We shall see why we propose to call these new entities *log-moments* and *log-cumulants*. Thanks to this approach, it is possible to analyse in a simpler way distributions with two or three parameters that have traditionally been used for imagery: the gamma distribution, Nakagami distribution, and \mathcal{K} distribution. Then, we will see how to tackle more complex problems like the distributions of the Pearson system, additive mixtures and distributions with heavy tails (i.e., distributions for which the mo-

ments cannot be defined starting from a certain order¹). Finally, we will analyse why the parameter estimators of these distributions based on the log-moments and log-cumulants have a lower variance than those obtained from the traditional moments and cumulants.

The remark can be made that a formalism with such similarity to the existing definitions cannot lead to intrinsically new results. It should be stressed that the essential contribution of this framework is to offer a signal and image processing methodology which proves, for certain applications, considerably easier to use than the traditional approaches. The major goal of this article is to illustrate its simplicity of implementation as well as its flexibility in use.

2 DEFINITION OF THE SECOND KIND CHARACTERISTIC FUNCTIONS

The objective of this section is to propose a formalism of *second kind statistics* based on the Mellin transform and redefine some elements of traditional statistics, namely the characteristic function yielding moments and cumulants, as outlined in the introduction.

2.1 First Characteristic Function of the Second Kind

Let X be a positive-valued random variable whose pdf, $p_x(u)$, is defined for $u \in \mathbb{R}^+$. The *first characteristic function of the second kind* is defined as the Mellin transform \mathcal{M} of $p_x(u)$:

$$\phi_x(s) = \mathcal{M}[p_x(u)](s) = \int_0^{+\infty} u^{s-1} p_x(u) du \quad (8)$$

provided that this integral converges, which is verified in general only for values of s located inside a strip delimited by two lines parallel to the secondary axis, i.e.

$$s = a + jb, \quad a \in]a_1; a_2[, \quad b \in \mathbb{R}$$

with a_2 commonly approaching $+\infty$, just as a_1 approaches $-\infty$. As the Mellin transform has an inverse [2], knowing $\phi_x(s)$, one can deduce $p_x(u)$ thanks to the relation:

$$p_x(u) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} u^{-s} \phi_x(s) ds$$

given that c is confined within the strip where the first characteristic function is defined (i.e, $c \in]a_1; a_2[$). Note that if $p_x(u)$ is a pdf, the second kind characteristic function satisfies the fundamental property:

$$\phi_x(s)|_{s=1} = 1.$$

1. Translator's remark: Note that author uses a strict definition of heavy-tailed distributions. An alternative and more common definition is that heavy-tailed distributions are not exponentially bounded. That is, they have heavier tails than the exponential distribution. Since the context of the discussion is modelling of multilook intensity radar data, it would be natural to replace the exponential distribution with the (generalised) gamma distribution in this criterion.

By analogy, the *second kind moments*, \tilde{m}_ν ($\nu \in \mathbb{N}$) are defined by the relation:

$$\tilde{m}_\nu = \left. \frac{d^\nu \phi_x(s)}{ds^\nu} \right|_{s=1}. \quad (9)$$

By virtue of a fundamental property of the Mellin transform [Col59]:

$$\mathcal{M}[f(u)(\log u)^\nu](s) = \frac{d^\nu \mathcal{M}[f(u)](s)}{ds^\nu}$$

which is evaluated at $s = 1$, the second kind moments can be written in two different ways:

$$\tilde{m}_\nu = \left. \frac{d^\nu \phi_x(s)}{ds^\nu} \right|_{s=1} \quad (10)$$

$$= \int_0^{+\infty} (\log u)^\nu p_x(u) du. \quad (11)$$

Eq. (11) suggests that we refer to the second kind moments as *log-moments*, which is adopted for the remainder of the article.

We now introduce the second kind mean or *log-mean* \tilde{m} . This auxiliary variable is defined by the following relation

$$\log \tilde{m} = \tilde{m}_1 \Leftrightarrow \tilde{m} = e^{\tilde{m}_1}.$$

Note that this entity, which is in fact the geometric mean, takes its values in \mathbb{R}^+ (a suitable scale for the variable u), whereas the log-moments take their values in \mathbb{R} (on logarithmic scale). It is thus possible to compare the mean \tilde{m} and the log-mean \tilde{m} , and the practical importance will be demonstrated for the gamma distribution.

Just as one traditionally defines the central moments, we introduce the definition of the central log-moments of order n , \tilde{M}_n :

$$\begin{aligned} \tilde{M}_n &= \int_0^{+\infty} (\log u - \tilde{m}_1)^n p_x(u) du \\ &= \int_0^{+\infty} \left(\log \frac{u}{\tilde{m}} \right)^n p_x(u) du. \end{aligned} \quad (12)$$

In particular, one readily finds the expression

$$\tilde{M}_2 = \tilde{m}_2 - \tilde{m}_1^2.$$

Thanks to this formalism, it is possible to obtain an analytical expression for the log-moments by simple derivation of the second kind characteristic function. We will look at the classical interpretation of the Mellin transform.

2.2 A First Interpretation of the Mellin Transform

By comparison of the moment definition in Eq. (3) and the definition of the first characteristic of the second kind in Eq. (8), one can write the generalised moments, m_ν :

$$m_\nu = \phi_x(s)|_{s=\nu+1} = \int_0^{+\infty} u^\nu p_x(u) du. \quad (13)$$

For $\nu \in \mathbb{N}$, these are the traditional moments. For $\nu \in \mathbb{R}^+$, we have the fractional moments, that have been used

by some authors (like the use of FLOM: Fractional Low Order Moments, in [3]). Provided that the Mellin transform is defined for values of $\nu \in \mathbb{R}^-$, it is justified to use *lower order moments* [4]. Lastly, in addition to moments defined for a value $\nu = a$ (i.e., traditional moments, fractional moments, or lower order moments), one can define moments of complex order with $\nu = a + jb$ for all b , this because the pdf $p_x(u)$ is positive by definition, a property which is trivial to verify.

2.3 Second Kind Cumulants or Log-Cumulants

Still by analogy with classical statistic for scalar real random variables defined on \mathbb{R} , the *second characteristic function of the second kind* is defined as the natural logarithm of the first characteristic function of the second kind:

$$\psi_x(s) = \log(\phi_x(s)). \quad (14)$$

The derivative of the second characteristic function of the second kind, evaluated at $s = 1$, defines *second kind cumulants* of order n :

$$\tilde{\kappa}_{x(n)} = \left. \frac{d^n \psi_x(s)}{ds^n} \right|_{s=1}. \quad (15)$$

Since formally, second kind cumulants are constructed according to the same rules as traditional cumulants, the relations between log-moments and log-cumulants are identical to the relations existing between moments and cumulants. For instance, the three first log-cumulants can be written as:

$$\begin{aligned} \tilde{\kappa}_1 &= \tilde{m}_1 \\ \tilde{\kappa}_2 &= \tilde{m}_2 - \tilde{m}_1^2 \\ \tilde{\kappa}_3 &= \tilde{m}_3 - 3\tilde{m}_1\tilde{m}_2 + 2\tilde{m}_1^3 \end{aligned}$$

As in the case of log-moments, we adopt the name *log-cumulants* for the second kind cumulants.

2.4 Some Properties of Log-Moments and Log-Cumulants

The utilisation of the Mellin transform requires knowledge about some of its specific properties. In particular, let us point out the definition of the Mellin convolution (which is an associative and commutative operation):

$$\begin{aligned} h = f \hat{\star} g &\Leftrightarrow h = \int_0^\infty f(y)g\left(\frac{u}{y}\right) \frac{dy}{y} \\ &\Leftrightarrow h = \int_0^\infty g(y)f\left(\frac{u}{y}\right) \frac{dy}{y}, \end{aligned} \quad (16)$$

Its fundamental property is similar to the convolution property of the Fourier transform:

$$\mathcal{M}[f \hat{\star} g](s) = \mathcal{M}[f](s) \mathcal{M}[g](s).$$

Note that if f and g are pdfs, then h is also a pdf (i.e., $h(u) \geq 0 \forall u \in \mathbb{R}^+$ and $\mathcal{M}[h]|_{s=1} = 1$).

The use of this operator finds an immediate application in the study of multiplicative noise. Let Y and Z

be two independent random variables whose respective pdfs, q_y and r_z , are defined on \mathbb{R}^+ . Consider a random variable X constructed by a multiplication of these two variables. It is thus a model of multiplicative noise. It is then shown that the pdf of X , p_x , is obtained as the Mellin convolution of q_y and r_z [5], [6]:

$$p_x = q_y \hat{\star} r_z.$$

The properties deduced in the following are formally identical to those obtained in the case of a traditional convolution (Eqs. (5)-(7)). If ϕ_x is the second kind characteristic function of X , ϕ_y is the second kind characteristic function of Y and ϕ_z is the second kind characteristic function of Z , the following relations are obtained:

$$\begin{aligned} \phi_x(s) &= \phi_y(s) \phi_z(s) \\ \psi_x(s) &= \psi_y(s) + \psi_z(s) \\ \tilde{\kappa}_{x(n)} &= \tilde{\kappa}_{y(n)} + \tilde{\kappa}_{z(n)} \quad \forall n \in \mathbb{N} \end{aligned} \quad (17)$$

It is noted in particular that in the case of multiplicative noise, the log-cumulants are additive. This property is not surprising since the common method used to handle multiplicative noise, transformation into logarithmic scale, allows us to treat noise of multiplicative nature like additive noise.

Finally note the following property:

$$u(f \hat{\star} g) = (u f) \hat{\star} (u g).$$

One can also, in a step similar to that of traditional convolution, define the inverse convolution (a non-commutative and non-associative operator). If the ratio

$$\frac{\mathcal{M}[f](s)}{\mathcal{M}[g](s)}$$

is defined in the vicinity of $s = 1$ such that the inverse Mellin transform can be evaluated, the following relation is posed:

$$h = f \hat{\star}^{-1} g \Leftrightarrow \mathcal{M}[h](s) = \frac{\mathcal{M}[f](s)}{\mathcal{M}[g](s)}.$$

With the above notation we have, given that the pdfs p_x , q_y and r_z exist:

$$p_x = q_y \hat{\star}^{-1} r_z,$$

from which we deduce:

$$\begin{aligned} \phi_x(s) &= \frac{\phi_y(s)}{\phi_z(s)} \\ \psi_x(s) &= \psi_y(s) - \psi_z(s) \\ \tilde{\kappa}_{x(n)} &= \tilde{\kappa}_{y(n)} - \tilde{\kappa}_{z(n)} \quad \forall n \in \mathbb{N} \end{aligned} \quad (18)$$

Finally, it can be useful to utilise the *Mellin correlation* (also a non-associative and non-commutative operation), which is defined by the relation:

$$h = f \hat{\otimes} g \Leftrightarrow \mathcal{M}[h](s) = \mathcal{M}[f](s) \mathcal{M}[g](2-s). \quad (19)$$

A pdf must satisfy $\mathcal{M}[h]|_{s=1} = 1$, to which h complies. Starting from this relation and using the same notation, we can, provided that p_x satisfies

$$p_x = q_y \hat{\otimes} r_z,$$

deduce the following:

$$\begin{aligned} \phi_x(s) &= \frac{\phi_y(s)}{\phi_z(2-s)} \\ \psi_x(s) &= \psi_y(s) - \psi_z(2-s) \\ \tilde{\kappa}_{x(n)} &= \tilde{\kappa}_{y(n)} + (-1)^n \tilde{\kappa}_{z(n)} \quad \forall n \in \mathbb{N} \end{aligned} \quad (20)$$

The following expression can then be shown:

$$h = f \hat{\otimes} g \Leftrightarrow h = \int_0^\infty f(uy) g(y) y dy. \quad (21)$$

We also note the property:

$$u(f \hat{\otimes} g) = (u f) \hat{\otimes} \left(\frac{g}{u}\right) \quad (22)$$

In fact, the interpretation of the Mellin correlation is founded on the analysis of the inverse distribution, i.e., the distribution $p_I(u)$ of the random variable $Y = 1/X$, where the random variable X follows the distribution $p(u)$. The relation between these distributions are known to be:

$$p_I(u) = \frac{1}{u^2} p\left(\frac{1}{u}\right).$$

By taking account of a fundamental property of the Mellin transform:

$$\mathcal{M}\left[\frac{1}{u} f\left(\frac{1}{u}\right)\right](s) = \mathcal{M}[f(u)](1-s),$$

it is easily deduced that

$$\mathcal{M}[p_I](s) = \mathcal{M}[p](2-s). \quad (23)$$

It is then seen that the Mellin correlation of a pdf q_y of the random variable Y and a pdf r_z of the random variable Z ,

$$q_y \hat{\otimes} r_z,$$

is simply a way to establish the pdf of the random variable Y/Z .

Lastly, as for the traditional characteristic function, it is interesting to note that the second kind characteristic function can be expanded in terms of log-cumulants:

$$\begin{aligned} \psi_x(s) &= \tilde{\kappa}_{x(1)}(s-1) + \frac{1}{2!} \tilde{\kappa}_{x(2)}(s-1)^2 \\ &+ \frac{1}{3!} \tilde{\kappa}_{x(3)}(s-1)^3 + \dots \end{aligned}$$

2.5 Theorem of Existence of Log-Moments and Log-Cumulants

We have just seen that the theoretical introduction of the log-moments and log-cumulants does not pose any formal problem. However, the existence of these entities has not been proven, and an interrogation into the requirements for their existence is needed. In this section we will present a theorem of strong conditions, that generally verify the existence of the log-moments and log-cumulants for the distribution usually applied in signal and image processing.

TABLE 1

Properties of the Mellin convolution, the inverse Mellin convolution, and the Mellin correlation of two distributions defined on \mathbb{R}^+ : p_A and p_B , with second kind characteristic functions ϕ_A and ϕ_B , and log-cumulants $\tilde{\kappa}_{A,n}$ and $\tilde{\kappa}_{B,n}$.

	Characteristic function	Cumulant
$p_A \hat{\star} p_B$	$\phi_A(s) \phi_B(s)$	$\tilde{\kappa}_{A,n} + \tilde{\kappa}_{B,n}$
$p_A \hat{\star}^{-1} p_B$	$\frac{\phi_A(s)}{\phi_B(s)}$	$\tilde{\kappa}_{A,n} - \tilde{\kappa}_{B,n}$
$p_A \hat{\otimes} p_B$	$\phi_A(s) \phi_B(2-s)$	$\tilde{\kappa}_{A,n} + (-1)^n \tilde{\kappa}_{B,n}$

Let $p(u)$ be a probability distribution defined on \mathbb{R}^+ , whose second kind characteristic function is $\phi(s)$. This pdf satisfies the relations:

- $p(u) \geq 0 \quad \forall u \geq 0$
- $\int_0^{+\infty} p(u) du = \phi(s)|_{s=1}$

Theorem 1: If a pdf has a second kind characteristic function defined on the set $\Omega =]s_A, s_B[$, where $s = 1 \in \Omega$, then all of its log-moments and log-cumulants exist.

Proof: The existence of the log-moments and log-cumulants depends by the convergence of the integral

$$\int_0^{+\infty} (\log u)^n p(u) du.$$

In order to study this improper integral, we will study its behaviour at 0 and at the limit to infinity.

- **close to infinity:** Let $\alpha \in \Omega$ such that $\alpha > 1$. Thus, $\exists \alpha > 1$ such that:

$$\phi(\alpha) = \int_0^{+\infty} u^{\alpha-1} p(u) du < \infty$$

which amounts to saying that the moments of $p(u)$ (integer order or fractional) can be calculated for all orders between 1 and α . Assume an integer $n \geq 1$. Two cases then arise:

- For $\forall x > 1$ we have $(\log x)^n < x^{\alpha-1}$. In this case, knowing that $p(u)$ is a pdf and satisfies $p(u) \geq 0$, one can write

$$\begin{aligned} & \lim_{b \rightarrow \infty} \int_1^b (\log u)^n p(u) du \\ & \leq \lim_{b \rightarrow \infty} \int_1^b u^{\alpha-1} p(u) du \leq \phi(\alpha) \end{aligned}$$

which demonstrates the convergence of the integral as $x \rightarrow \infty$.

- There exists a constant $c > 1$ such that $(\log c)^n = c^{\alpha-1}$. Then, for $\forall x > c$ we have $(\log x)^n \leq x^{\alpha-1}$. By an identical argument as for the previous case, we deduce that

$$\lim_{b \rightarrow \infty} \int_c^b (\log u)^n p(u) du \leq \phi(\alpha)$$

which demonstrates the convergence of the integral as $x \rightarrow \infty$.

- **close to 0:**

- First of all, consider the particular case where the pdf is bounded. Assume that

$$\exists A : \quad \forall u \in [0, 1], \quad p(u) \leq A,$$

and calculate the limit

$$\lim_{a \rightarrow 0} \int_a^1 (\log u)^n p(u) du.$$

Since the pdf is bounded, we have for $\forall a \in]0, 1[$ that

$$\begin{aligned} \left| \int_a^1 (\log u)^n p(u) du \right| & \leq \left| \int_a^1 (\log u)^n A du \right| \\ & \leq A \left| \int_a^1 (\log u)^n du \right|. \end{aligned}$$

The following property

$$\left| \lim_{a \rightarrow 0} \int_a^1 (\log u)^n du \right| = \Gamma(n+1)$$

proves the convergence at 0.

- In the general case, the study of the convergence starts from the variables change $x \rightarrow \frac{1}{x}$, thereafter utilising the convergence property that we have just shown for the case $x \rightarrow \infty$.

We deduce that if a probability distribution with bounded values possesses moments (fractional or integer ordered) of order strictly larger than 0 and strictly smaller than 0, then all its log-moments and log-cumulants exist. \square

Note that a far more elegant and concise proof, founded on the properties of analytical functions, can be worked out without major problems based on the assumption that $\phi(s)$ is holomorphic [7], and thus differentiable up to all orders at $s = 1$.

2.6 Comparison with Logarithmic Transformation

At this stage, one can wonder what the advantages of this new approach are, and whether a simple transformation into logarithmic scale would lead to the same result. We will show that in order to calculate a characteristic function after logarithmic transformation, one effectively has to calculate the Mellin transform of the original pdf.

We shall consider a random variable x with density defined over real positive numbers. Its pdf, $p_x(u)$, is thus defined for $u \in \mathbb{R}$, and the characteristic function is written:

$$\Phi_x(v) = \int_0^{+\infty} e^{jvu} p_x(u) du.$$

Then perform a logarithmic transformation. The new random variable y is described the pdf $q_y(w)$, defined for $w \in \mathbb{R}$, with $w = \log u$. This pdf results from p_x with the relation given by

$$q_y(w) = e^w p_x(e^w).$$

Now calculate the characteristic function of the random variable y :

$$\begin{aligned}
 \Phi_y(v) &= \int_{-\infty}^{+\infty} e^{jvw} q_y(w) dw \\
 &= \int_{-\infty}^{+\infty} e^{jvw} e^w p_x(e^w) dw \\
 &= \int_0^{+\infty} e^{jv \log u} p_x(u) du \text{ with } u = e^w \\
 &= \int_0^{+\infty} u^{jv} p_x(u) du.
 \end{aligned} \tag{24}$$

The relation in (24) is recognised as the Mellin transform of $p_x(u)$ at $s = 1 + jv$:

$$\Phi_y(v) = \phi_x(s)|_{s=1+jv}. \tag{25}$$

This relation shows that if one knows the Mellin transform of a pdf (i.e., its second kind characteristic function), then one knows the ordinary characteristic function in logarithmic scale.

On logarithmic scale, moments and cumulants are deduced by differentiation (simple or logarithmic) of expression (25), which is equivalent to what was defined in Eq. (10). This is another way to justify the terms *log-moments* and *log-cumulants*. We note, however, that the second kind statistics represent a generic method to find log-moments and log-cumulants directly without requiring a variable change (logarithmic transformation) and also without having to calculate the new distribution for the transformed variable.

Moreover, we will see that in the cases generally encountered in signal and image processing, and where the entities are defined on \mathbb{R}^+ , it is easier to calculate the Mellin transform than the Fourier transform. Thus, our approach simplifies the working of the problem. In addition, when the Mellin transform is known, one automatically obtains:

- the moments, by inserting positive integers for the Mellin transform variable s , and
- the log-moments, by differentiating the Mellin transform with respect to s and evaluating at $s = 1$.

This should be appreciated by any practitioner, since, by applying a single transformation to the distribution, both moments and log-moments are obtained.

2.7 Comparison between Integral Transforms

The use in this context of an ignored transform: the Mellin transform, may surprise, since there exist more common invertible transforms, such as the Laplace transform, that could potentially play an important role in the study of distributions defined on \mathbb{R}^+ . At this stage, it is important to observe what the relations between the Fourier transform (\mathcal{F}), the Laplace transform (\mathcal{L}), and the Mellin transform (\mathcal{M}) are. When it exists, the Laplace transform of a pdf $p(u)$ is written:

$$\mathcal{L}[p(u)](\sigma) = \int_0^{\infty} e^{-\sigma u} p(u) du$$

while the first characteristic function of this pdf is written

$$\mathcal{F}[p(u)](v) = \int_{-\infty}^{+\infty} e^{jvu} p(u) du.$$

The following relation is immediately deduced:

$$\mathcal{F}[p(u)](v)|_{v=-\frac{b}{2\pi}} = \mathcal{L}[p(u)](\sigma)|_{\sigma=jb}$$

Because the Laplace transform variable is a complex entity, one may consider that the Laplace transform could allow for an analytical continuation of the characteristic function [8]. However, the intrinsic properties of the Laplace transform are the same as those of the Fourier transform. A logarithmic transformation (in which case, one will have to use the bilateral Laplace transform) will in reality turn the Laplace transform into a Mellin transform:

$$\mathcal{L}[\tilde{p}(u)]|_{\sigma=a+jb} = \mathcal{M}[p(u)]|_{s=-a-jb} = \phi_x(s)|_{s=-a-jb}$$

where \tilde{p} is the pdf on logarithmic scale.

There are such strong relations between these transforms that, most likely, nothing fundamentally new will be found by the use of the Mellin transform. Therefore, it seems that the choice should be dictated by practical considerations. We have seen that the Mellin transform makes it possible to obtain traditional moments and log-moments at the same time, without having to derive the pdf on logarithmic scale. Moreover, the Mellin transform of the experimental distributions commonly used in signal and image processing can be found in tables. This justifies a further look into the use of this rather ignored transform. This is the pragmatic view which motivates the derivations of this article.

3 FUNDAMENTAL EXAMPLES

We will illustrate the new approach by applying it to distributions used to model synthetic aperture radar (SAR) images. These are the gamma and the generalised gamma distribution (intensity images with fully developed speckle), the Rayleigh and the Nakagami distribution (amplitude images with fully developed speckle), and finally the \mathcal{K} distribution (an intensity distribution modelling fully developed speckle modulated multiplicatively by gamma distributed texture). Even if some of the results obtained are trivial, it seems important to be able to carry out comparisons with these simple and well-known cases, in particular in order to handle the problem of estimating the distribution parameters, an aspect which will be looked at in Section 5.

3.1 Gamma and Generalised Gamma Distribution

The two parameter gamma distribution, denoted $\mathcal{G}[\mu, L]$, is a type III solution of the Pearson system [9]. It is defined on \mathbb{R}^+ as

$$\mathcal{G}[\mu, L](u) = \frac{1}{\Gamma(L)} \frac{L}{\mu} \left(\frac{Lu}{\mu} \right)^{L-1} e^{-\frac{Lu}{\mu}} \text{ with } L > 0 \tag{26}$$

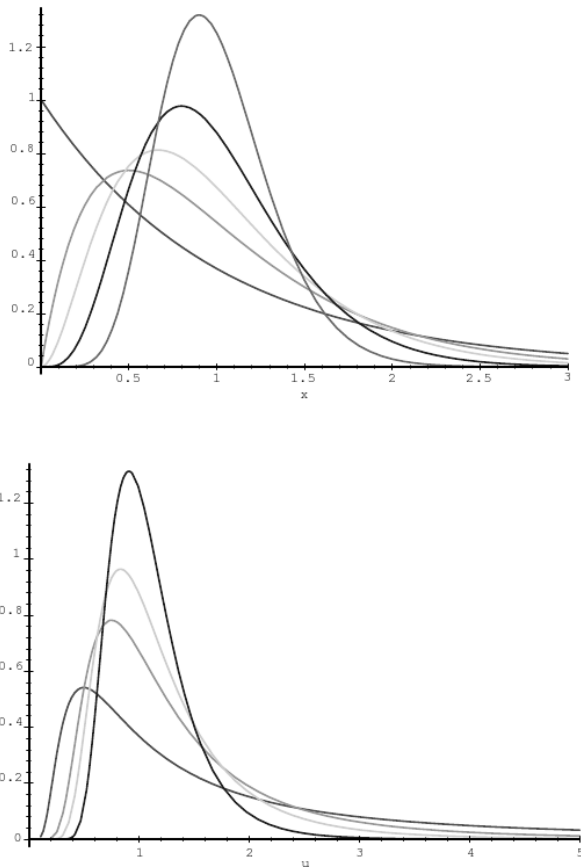


Fig. 1. Top: Gamma distribution $\mathcal{G}[\mu, L]$ with $\mu = 1$ and $L = 1, 2, 3, 5$ and 10 . Bottom: Inverse gamma distribution $\mathcal{IG}[\mu, L]$ with $\mu = 1$ and $L = 1, 2, 3, 5$ and 10 .

We see that μ is a scale parameter and that L is a shape parameter (Figure 1).

The particular case of $L = 1$ corresponds to the true gamma distribution², which is well-known from the radar literature as a model of fully developed speckle in single-look images. The case of $L = \frac{1}{2}$ gives the χ^2 distribution.

The Fourier transform tables show that the characteristic function is written as:

$$\Phi(\nu) = \left(\frac{L}{\mu}\right)^L \frac{e^{jL \arctan(\frac{\nu\mu}{L})}}{\left(\nu^2 + \frac{L^2}{\mu^2}\right)^{\frac{L}{2}}} \quad (27)$$

whose complicated expression makes it difficult to use in practice.

On the other hand, by use of known Mellin transforms that can be found in tables [2], [10], the second kind characteristic function can be expressed in terms of the gamma function as:

$$\phi_x(s) = \mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)} \quad (28)$$

2. Translator's remark: Note that the author uses the terms 'generalised gamma distribution' and '(true) gamma distribution' for the distributions more commonly referred to as the 'gamma distribution' and the 'exponential distribution', respectively.

The classical moments $m_n, \forall n \in \mathbb{N}$ are much easier to derive from this function than from (27):

$$m_n = \mu_n \frac{\Gamma(L+n)}{L^n \Gamma(L)} \quad (29)$$

from which we have the well-known moments:

$$m_1 = \mu \quad m_2 = \frac{L+1}{L} \mu^2.$$

This equation system is analytically invertible, and from the first two moments we derive the following relations for the parameters μ and L :

$$\mu = m_1 \quad (30)$$

$$L = \frac{m_1^2}{m_2 - m_1^2} = \frac{1}{\frac{m_2}{m_1^2} - 1} \quad (31)$$

Note that this distribution is asymmetric, and its mode value is given by:

$$m_{\text{mode}} = \frac{L-1}{L} \mu \leq \mu.$$

We also remark that the second kind characteristic function can be separated into a first term, μ^{s-1} , and a second term that only depends on L , the shape parameter. As L goes to infinity, $\frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)}$ goes towards 1, and $\mathcal{G}[\mu, L]$ converges in distribution to the homothetic distribution $\mathcal{H}[\mu]$:

$$\begin{aligned} \phi_x(s) &\rightarrow \mu^{s-1} \\ \Leftrightarrow \mathcal{G}[\mu, L](u) &\rightarrow \mathcal{H}[\mu](u) = \frac{1}{\mu} \delta(\mu u - 1) \end{aligned}$$

We note that the homothetic distribution can be seen as a degenerate Gaussian distribution (i.e. with zero variance). It seems to confirm what many experts of radar imaging has pointed out, that the gamma distribution tends towards a Gaussian distribution as L goes to infinity, but by the alternative denotation we avoid abuse of language that can lead to confusion.

Another major point specific to the distributions defined on \mathbb{R}^+ rests on the fact that the Mellin transform of $\mathcal{G}[\mu, L]$ is defined for $s > 1 - L$. It is seen that, for $L > 1$, it is possible to have negative values of $s - 1$ and thus *lower order moments*. Qualitatively, the lower order moments - i.e. positive powers of $\frac{1}{u}$ - mainly reflect the weight of the distribution between 0 and μ , while the traditional moments - i.e. positive powers of u - rather analyse the distribution between μ and ∞ . Thanks to the lower order moments, it is possible to analyse selectively the left or the right tail of a probability distribution. The importance of this observation for asymmetrical distributions such as the gamma distribution is evident.

The first two log-cumulants of $\mathcal{G}[\mu, L]$ are expressed by the following relations, where $\Phi(\cdot)$ is the digamma function and $\Psi(r, \cdot)$ is the r -th order polygamma function, i.e. the r -th order derivative of the digamma function:

$$\tilde{\kappa}_{x(1)} = \log(\mu) + \Psi(L) - \log(L) \quad (32)$$

$$\tilde{\kappa}_{x(2)} = \Psi(1, L) \geq 0 \quad (33)$$

$$\tilde{\kappa}_{x(3)} = \Psi(2, L) \leq 0 \quad (34)$$

and it is trivial to show that

$$\tilde{\kappa}_{x(r)} = \Psi(r-1, L) \quad \forall r > 1,$$

which expresses the property that the log-cumulants depend only on L from second order and upwards.

We note that the property

$$\lim_{L \rightarrow \infty} (\Psi(L) - \log(L)) = 0$$

associated with the fact that the polygamma functions go towards 0 at infinity, can be used to show that the gamma distribution converges towards the homothetic distribution as L goes to infinity.

Remark that the third order log-cumulant is negative. This illustrates that, for the gamma distribution, the left tail is heavier than the right tail of the distribution, which decreases very quickly as the argument approaches infinity.

The log-mean is written:

$$\tilde{m} = \mu \frac{e^{\Psi(L)}}{L} \quad (35)$$

It is interesting to note the two following points:

- $\frac{e^{\Psi(L)}}{L} \leq 1$: The log-mean is less than the mean value. Note that this property is valid for all L .
- $\frac{e^{\Psi(L)}}{L} \geq \frac{L-1}{L}$: The log-mean is larger than the mode value.

A more complete analysis would show that the log-mean is also lower than the median value, defined by

$$\int_0^{m_{med}} \mathcal{G}(u) du = 0.5.$$

It can also be justified to use the log-mean instead of the traditional mean. This gives interesting results in certain applications of image processing [11].

Finally, by a logarithmic transformation, the gamma distribution $\mathcal{G}[\mu, L](u)$ becomes the Fisher-Tippett distribution $\mathcal{G}_{\mathcal{FT}}[\tilde{\mu}, L](w)$ with $\tilde{\mu} = \log \mu$ and $w = \log u$:

$$\mathcal{G}_{\mathcal{FT}}[\tilde{\mu}, L](w) = \frac{L^L}{\Gamma(L)} e^{L(w-\tilde{\mu})} e^{-Le^{(w-\tilde{\mu})}}$$

Its characteristic function is obtained by taking the Fourier transform. Unfortunately, the required relation is not found in tables. This is commonly circumvented by showing that the evaluation amounts to calculating a Mellin transform. In effect, one applies (25) unknowingly.

To conclude, it is seen that in the case of the generalised gamma distribution, the second kind statistics approach allows us:

- to obtain a simpler expression for the second kind characteristic function than for the classical characteristic function.
- to estimate the distribution parameters more efficiently by inversion of Eqs. (32) and (33).
 - The shape parameter L is easily derived from the second order log-cumulant, even if no analytical formulation can be found, since the

polygamma functions are monotonous and easy to invert numerically (Tabulation can also be used to save computation time). The variance of this estimator is evaluated in Section 5, and we will see that it is notably lower than the variance obtained with the method of moments estimator, as defined in (31).

– After L is known, μ can be derived from the expression of the first order log-cumulant.

- to propose a “typical value” for use in image processing, lying between the mode and the mean, which realistically represents a sample if it can be regarded as homogeneous.

3.2 Rayleigh and Nakagami Distribution

We will now handle a problem specific to SAR imagery, namely the transformation of intensity data to amplitude data. Even if models have simple expressions for intensity data (the gamma distribution is known to all scientific communities), the images are quite often available as amplitude data, which will reveal new problems regarding parameter estimation. In this article, we will thus address the transformation from intensity data that follow the gamma distribution, to amplitude data with their resulting distribution.

The Nakagami distribution³ is the name which in the radar literature has been associated with amplitude data that follow a gamma distribution when transformed into the intensity domain. It is thus a two parameter distribution: $\mathcal{RN}[\mu, L]$, given by:

$$\mathcal{RN}[\mu, L](u) = \frac{2}{\mu} \frac{\sqrt{L}}{\Gamma(L)} \left(\frac{\sqrt{L}u}{\mu} \right)^{2L-1} e^{-\left(\frac{\sqrt{L}u}{\mu}\right)^2}. \quad (36)$$

For $L = 1$, one retrieves the Rayleigh distribution:

$$\mathcal{RN}[\mu, L=1](u) = \frac{2}{\mu} \left(\frac{u}{\mu} \right) e^{-\left(\frac{u}{\mu}\right)^2}.$$

The fundamental relation between the Nakagami distribution (for amplitude) and the generalised gamma distribution (for intensity, i.e. squared amplitude) is obtained by a simple variable change, which can be written as:

$$\mathcal{RN}[\mu, L](u) = 2u\mathcal{G}[\mu_G, L](u^2). \quad (37)$$

By means of the following two Mellin transform properties [2]:

$$\mathcal{M}[u^a f(u)](s) = \mathcal{M}[f(u)](s+a)$$

$$\mathcal{M}[f(u^a)](s) = \frac{1}{a} \mathcal{M}[f(u)]\left(\frac{s}{a}\right)$$

3. It is important to return to Nakagami the paternity of this distribution described by two parameters: mean and shape parameter, which has often been wrongly associated with the generalised gamma distribution. The formalism was proposed in 1942 by Nakagami in an exhaustive study of the “m-distribution”. It was not published in English until 1960 [12]. Of course, this distribution is for instance found in [13] as the result of transformations starting from the gamma distributions. However, it seems that Nakagami performed the first complete study.

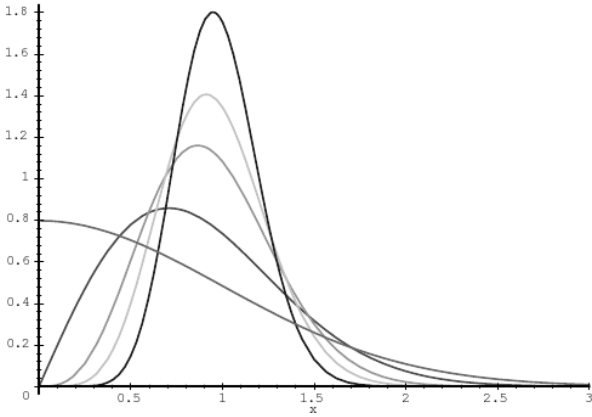


Fig. 2. Rayleigh-Nakagami distribution $\mathcal{RN}[\mu, L]$ with $\mu = 1$ and $L = 0.5, 1, 2, 3, 5$.

and knowing $\phi_{\mathcal{G},x}(s)$, the second kind characteristic function of the gamma distribution, the second kind characteristic function of the Nakagami distribution can be derived directly as:

$$\phi_{\mathcal{RN},x}(s) = \phi_{\mathcal{G},x}\left(\frac{s+1}{2}\right)$$

which, by inserting $\mu_{\mathcal{G}} = \mu^2$, allows us to write:

$$\phi_{\mathcal{RN},x}(s) = \mu^{s-1} \frac{\Gamma\left(\frac{s-1}{2} + L\right)}{L^{\frac{s-1}{2}} \Gamma(L)}.$$

This reasoning applies also elsewhere, regardless of the power to which u is raised in the change of variable. It is easily shown that for $v = u^\alpha$, we have

$$p_u(u) = \alpha u^{\alpha-1} p_v(u^\alpha)$$

and the second kind characteristic function of the random variable u is derived directly from the properties of the Mellin transform as:

$$\phi_u(s) = \phi_v\left(\frac{s + \alpha - 1}{\alpha}\right). \quad (38)$$

Note that this result would be useful for the analysis of the Weibull distribution [14], another well-known radar distribution, which we will not address in this article.

The classical moments of the Nakagami distribution follow directly from $\phi_{\mathcal{RN},x}(s)$:

$$m_1 = \frac{\Gamma\left(L + \frac{1}{2}\right)}{\sqrt{L}\Gamma(L)}\mu \quad m_2 = \mu^2.$$

Take note of a peculiarity of this distribution: There is a very simple relation between the parameter μ and the second order moment, not the first order moment. On the other hand, the implicit expression of L obtained through the first order moment is very hard to handle. We cannot obtain a simple inversion formula (as in the gamma distribution case, where (31) gave L directly in terms of m_1 and m_2) to solve for L .

The mode of this pdf is

$$m_{\text{mode}} = \sqrt{\frac{2L-1}{2L}}\mu$$

The log-cumulants are derived directly from those of the gamma distribution as:

$$\begin{aligned} \tilde{\kappa}_{\mathcal{RN},x(r)} &= \left. \frac{d^r \psi_{\mathcal{RN}}(s)}{ds^r} \right|_{s=1} \\ &= \left. \frac{d^r \log \phi_{\mathcal{RN}}(s)}{ds^r} \right|_{s=1} \\ &= \left. \frac{d^r \log \phi_{\mathcal{G}}\left(\frac{s+1}{2}\right)}{ds^r} \right|_{s=1} \\ &= \left(\frac{1}{2}\right)^r \left. \frac{d^r \log \phi_{\mathcal{G}}(s')}{ds'^r} \right|_{s'=1} \\ &= \left(\frac{1}{2}\right)^r \tilde{\kappa}_{\mathcal{G},x(r)} \end{aligned}$$

From this we deduce:

$$\begin{aligned} \tilde{\kappa}_{x(1)} &= \log(\mu) + \frac{1}{2}\Psi(L) - \frac{1}{2}\log(L) \\ \tilde{\kappa}_{x(2)} &= \frac{1}{4}\Psi(1, L) \end{aligned}$$

and for all $r > 1$:

$$\tilde{\kappa}_{x(r)} = \left(\frac{1}{2}\right)^r \Psi(r-1, L)$$

More generally, it is shown for $v = u^\alpha$ that

$$\tilde{\kappa}_{p_u,x(r)} = \left(\frac{1}{\alpha}\right)^r \tilde{\kappa}_{p_v,x(r)}$$

In this case, L can be calculated directly if the second order log-cumulant is known. The problem we meet is of the same kind as for the gamma distribution, namely inversion of polygamma functions.

The log-mean is written

$$\tilde{m} = \mu \frac{e^{\frac{\Psi(L)}{2}}}{\sqrt{L}} \quad (39)$$

To conclude, the motivation of our approach is seen from the fact that the analytical expressions of the log-moments and the log-cumulants have a complexity comparable with the case of the gamma distribution, which is not the case in traditional statistics, where a simple relation between the first two moments and the shape parameter L cannot be obtained.

3.3 Inverse Gamma Distribution

The inverse gamma distribution is another two parameter distribution which is also a solution of the Pearson system (the type V solution). It is expressed as

$$\mathcal{IG}[\nu, L](u) = \frac{1}{\Gamma(L)} \frac{1}{L\nu} \left(\frac{L\nu}{u}\right)^{L+1} e^{-\frac{L\nu}{u}}$$

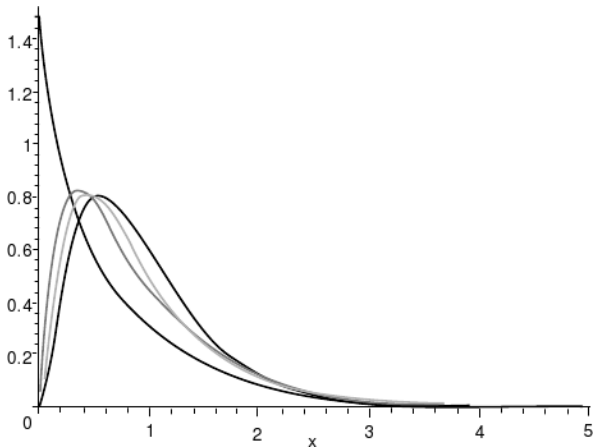


Fig. 3. \mathcal{K} distribution: Fully developed speckle ($L = 3$) modulated multiplicatively by a Rayleigh distribution with $\mu = 1$ and 1, 3, 5 and 10.

where $L \geq 0$ and $\nu > 0$. Its second kind characteristic function is written:

$$\phi_x(s) = \nu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)}.$$

It is seen that the n -th order moments of the inverse gamma distribution are not defined for $n \geq L$. The inverse gamma distribution is thus an example of heavy tailed distributions. Its log-cumulants, that exist for all orders, are written:

$$\begin{aligned} \tilde{\kappa}_{x(1)} &= \log(\nu) - \Phi(L) + \log(L) \\ \tilde{\kappa}_{x(2)} &= \Psi(1, L) \\ \tilde{\kappa}_{x(r)} &= (-1)^r \Psi(r-1, L) \quad \forall r > 1 \end{aligned}$$

For even r , these are the same as those of the gamma distribution. For odd r , they are opposite (See the more general relation in (20)). As for the gamma distribution, from second order and upwards, the log-cumulants depend only on the shape parameter L . Note that the third order log-cumulant is positive, which is a sufficient condition for being heavy-tailed.

This distribution could also have been introduced as the inverse of the gamma distribution (cf. Section 2.4), which would make it possible to deduce the log-cumulants directly. However, it was important to recall that the inverse gamma distribution is a particular solution of the Pearson system and associated with its own share of work in the literature.

3.4 \mathcal{K} Distribution

With the \mathcal{K} distribution, we will show that second order statistics provide an estimation method for the parameters of a complex distribution by simple application of the results already achieved for the gamma distribution. The \mathcal{K} distribution $\mathcal{K}[\mu, L, M]$ has three parameters and

is defined as

$$\begin{aligned} \mathcal{K}[\mu, L, M](u) &= \frac{1}{\Gamma(L)\Gamma(M)} \frac{2LM}{\mu} \left(\frac{LMu}{\mu}\right)^{\frac{M+L}{2}-1} \\ &\times K_{M-L} \left[2 \left(\frac{LMu}{\mu}\right)^{1/2} \right] \end{aligned} \quad (40)$$

where $K_n(\cdot)$ is the modified Bessel function of the second kind with order n . On this form, calculations of moments and log-moments require good knowledge of Bessel function properties as well as tables of transforms of Bessel functions.

In fact, the \mathcal{K} distribution is the distribution which is followed by a random variable defined as the product of two independent variables that are both gamma distributed. Note that this definition made it possible for Lomnicki [15] to retrieve Eq. (40) using, already at this time, the Mellin transform.

More precisely, it is possible to define the $\mathcal{K}[\mu, L, M]$ distribution as a Mellin convolution of two gamma distributions [6]:

$$\mathcal{K}[\mu, L, M] = \mathcal{G}[1, L] \star \mathcal{G}[\mu, M]$$

This definition greatly simplifies the calculations of the second kind characteristic function and thus the moments and log-cumulants. In effect, from the properties of the Mellin convolution (Section 2.4) and knowing the characteristics of the gamma distribution, one can write the second kind characteristic function of the \mathcal{K} distribution like a product of the second kind characteristic functions of the gamma distributions $\mathcal{G}[1, L]$ and $\mathcal{G}[\mu, M]$:

$$\phi_x(s) = \mu^{s-1} \frac{\Gamma(L+s-1)\Gamma(M+s-1)}{L^{s-1}\Gamma(L)M^{s-1}\Gamma(M)} \quad (41)$$

which allows us to immediately deduce the classical moments m_1 and m_2 without using the definition of the \mathcal{K} distribution, and thus without needing to know the properties of the Bessel function:

$$m_1 = \mu \quad m_2 = \mu^2 \frac{L+1}{L} \frac{M+1}{M}.$$

In the same manner, we obtain directly the first two log-cumulants as the sum of the log-cumulants of the gamma distributions $\mathcal{G}[1, L]$ and $\mathcal{G}[\mu, M]$:

$$\tilde{\kappa}_{x(1)} = \log \mu + (\Psi(L) - \log(L)) + (\Psi(M) - \log(M)) \quad (42)$$

$$\tilde{\kappa}_{x(2)} = \Psi(1, L) + \Psi(1, M) \quad (43)$$

$$\tilde{\kappa}_{x(3)} = \Psi(2, L) + \Psi(2, M) \quad (44)$$

and we can show that for all $r > 1$:

$$\tilde{\kappa}_{x(r)} = \Psi(r-1, L) + \Psi(r-1, M).$$

Finally, the normalised second order moment is written:

$$\tilde{M}_2 = \Psi(1, L) + \Psi(1, M) \quad (45)$$

and the log-mean:

$$\tilde{m} = \mu \frac{e^{\Psi(L)}}{L} \frac{e^{\Psi(M)}}{M}. \quad (46)$$

Also here, it is easy to derive the shape parameters L and M by virtue of the second and third order log-cumulants (Eqs. (43) and (44)). A simple numerical method is proposed and tested in [14]. The scale parameter is derived from the first order log-cumulant (Eq. (42)).

This method is much easier than the traditional method of moments, which results in a third degree equation. Note also that maximum likelihood estimation cannot be applied for this distribution [16].

4 APPLICATIONS

Second kind statistics prove easy to put into practice in the framework of fundamental probability distribution defined on \mathbb{R}^+ . Except for the introduction of the gamma function and its logarithmic derivatives (the polygamma functions), the obtained expressions contain no difficult terms. On the contrary, they are simple and easy to comprehend.

We will now look at several innovative applications of this model:

- A new approach to analysis of the three parameter distributions used to model SAR imagery
- The case of additive mixtures of gamma distributions, for which the traditional approaches lead to expressions that are very hard to handle
- The case of positive α -stable distributions, used for instance by Pierce to characterise sea clutter [17], for which it is difficult to estimate the parameters. (The analytical expression of the pdf for such heavy-tailed distributions is generally not known.)
- Finally, another example of the α -stable distribution proposed by Kuruoğlu and Zerubia [18], which can be seen as a generalisation of the Rayleigh distribution.

We will start by pointing out a method classically used to characterise these distributions: the use of the parameters of asymmetry β_1 (skewness) and peakedness β_2 (kurtosis).

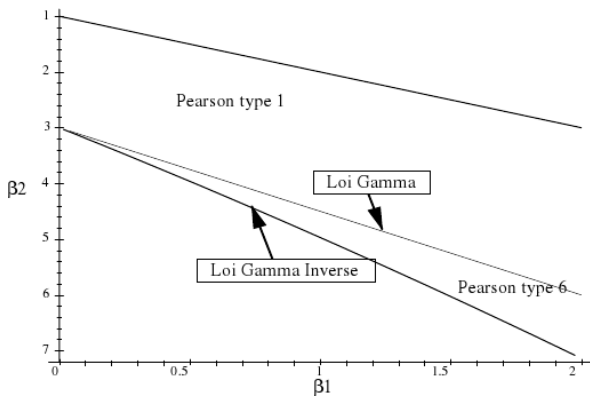


Fig. 4. The Pearson system displayed in a (β_1, β_2) diagram.

4.1 (β_1, β_2) Diagram

Traditionally, the skewness and kurtosis are used to characterise distributions belonging to the Pearson system. These two coefficients are written in terms of the second, third and fourth order moment:

$$\beta_1 = \frac{M_3^2}{M_2^3}$$

$$\beta_2 = \frac{M_4}{M_2^2}$$

The curves obtained for the Pearson system are shown in Figure 4 in their classical representation. The characteristic point of $(\beta_1 = 0, \beta_2 = 3)$ corresponds to the Gaussian case (It is invariant with respect to variance).

Because of the choice of squaring the third order central moment in β_1 , this coefficient is not able to distinguish between distributions that have skewness of the same magnitude but with opposite sign. Thus, it is not possible to separate between “standard” distributions and heavy-tailed distributions. Hence, this diagram seems to be badly adapted to the distributions defined on \mathbb{R}^+ .

4.2 Characterisation of Texture Distributions in the $(\tilde{\kappa}_3, \tilde{\kappa}_2)$ Diagram

The (β_1, β_2) diagram is founded on the calculation of traditional centred moments and aims at comparing distributions against the reference Gaussian distribution, for which the skewness is zero and the kurtosis is directly related to the variance (σ). It is then natural to propose a similar approach, founded on the functions of second kind statistics. We propose in this section a new method: the (κ_3, κ_2) diagram, that is, the representation of third order log-cumulants versus second order log-cumulants (that are always positive or zero for pdfs defined on \mathbb{R}^+).

In this diagram, the origin corresponds to the homothetic distribution. Because of the asymptotic behaviour of the gamma distribution and the inverse gamma distribution at $L \rightarrow \infty$, these distributions are represented by curves joining at the origin. As noted, the gamma

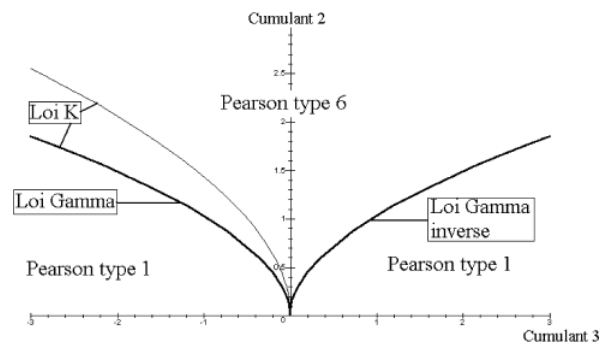


Fig. 5. The Pearson system and the \mathcal{K} distribution displayed in a (κ_3, κ_2) diagram.

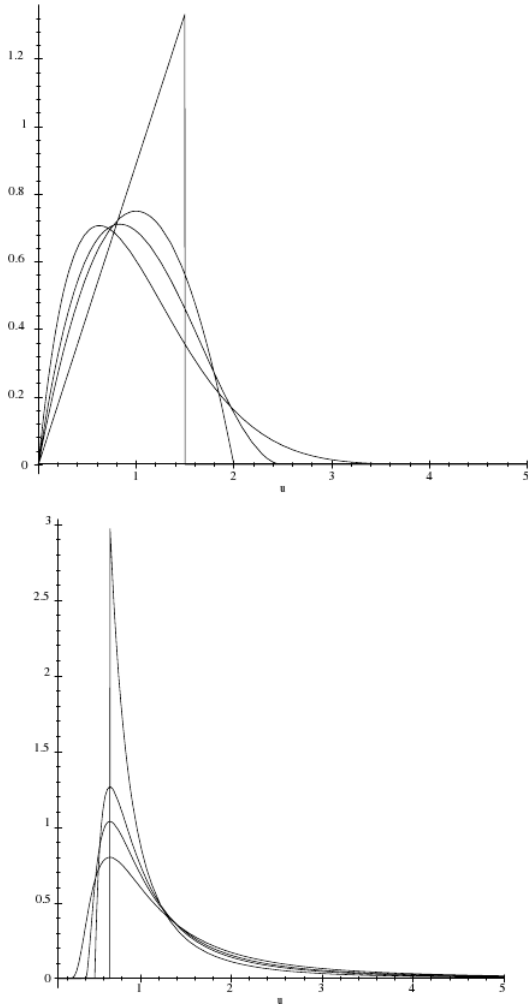


Fig. 6. The Pearson type I distribution (top) and the inverse Pearson type I distribution (bottom).

distribution has negative values for the third-order log-cumulants, while the heavy-tailed inverse gamma distribution has positive values. Note that it is easy to show that the log-normal distribution, for which all log-cumulants of order $n > 2$ is zero, occupies the second axis.

Figure 5 places the gamma distribution and the inverse gamma distribution in the diagram, together with the \mathcal{K} distribution (which occupies a surface above the gamma distribution, limited upwards by a curve defined by the distribution $\mathcal{K}[L, L]$) as well as the Pearson distributions of type I (standard and inverse) and type VI. We will see in the following section that the inverse Pearson distributions of type I find their natural place in this diagram, but offer some theoretical surprises.

4.3 An Original Approach to Characterisation of Three-Parameter Distributions Used for SAR Imagery

Knowing the two elementary two parameter distributions (the gamma distribution and the inverse gamma distribution), it falls natural to make use of these as basic

generating functions to obtain a kind of grammar by using elementary operations like the Mellin convolution and the inverse Mellin convolution (One could also have used the Mellin correlation instead of the Mellin convolution while inverting one of the distributions). Assume that we have two distributions p_A and p_B with respective second kind characteristic functions ϕ_A and ϕ_B and log-cumulants $\tilde{\kappa}_A$ and $\tilde{\kappa}_B$. Applying a Mellin convolution or an inverse Mellin convolution will correspond to forming the product or ratio of their second kind characteristic functions, and the sum or difference of the log-cumulants, respectively (See Table 1).

The characteristic functions of the distributions obtained by direct or inverse convolution of the two normalised distributions: the gamma distribution ($\mathcal{G}[1, L]$) or the inverse gamma distribution ($\mathcal{G}[1, M]$), are included in Table 2. From these expressions, and by consulting tables of the Mellin transform (and also using the properties of the transform), it is possible to retrieve the analytical expressions of these distributions without further calculations [14]. Furthermore, while considering only the second and third order second kind cumulants, Table 3 summarises the result obtained by direct or inverse convolution of the two normalised distributions: the gamma distribution ($\mathcal{G}[1, L]$) or the inverse gamma distribution ($\mathcal{G}[1, M]$). Recall that the second and third order second kind cumulants only depend on the shape parameter.

The distributions traditionally used in processing of SAR data are found in this table. These are

- The \mathcal{K} distribution
- The solutions of the Pearson system [6] corresponding to the distributions defined on \mathbb{R}^+ , that is, in addition to the gamma and inverse gamma distribution, also the type I solutions (also known as the beta distribution) and the type VI solutions (known as the Fisher distribution).

Moreover, uncommon distributions are generated by this algebraic method. It provides:

- A new distribution which is effectively the inverse Pearson distribution of type I, denoted $\mathcal{IPI}[\xi, L, M]$:

$$\mathcal{IPI}[\xi, L, M] = \frac{M}{L\xi} \frac{\Gamma(M)}{\Gamma(L)\Gamma(M-L)} \left(\frac{L\xi}{Mu}\right)^{L+1} \left(1 - \frac{L\xi}{Mu}\right)^{M-L-1} \quad (47)$$

with $u \geq \frac{L\xi}{M}$ and $M \geq L + 1$.

This expression is derived simply by means of the Mellin transform tables [2], since the distribution can be expressed by the relation:

$$\mathcal{IPI}[\xi, L, M] = \mathcal{IG}[\xi, L] \hat{\star}^{-1} \mathcal{IG}[1, M]$$

whose characteristic function is

$$\xi^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)} \frac{M^{1-s}\Gamma(M)}{\Gamma(M+1-s)}$$

TABLE 2

Second kind characteristic functions of the different distributions obtained by Mellin convolution (direct and inverse) of the gamma distribution $\left(\phi(s) = \mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)}\right)$ and inverse gamma distribution $\left(\phi(s) = \mu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)}\right)$. The distributions whose names are typeset with boldface correspond to new analytical expressions. The second and third order log-cumulants are included in Table 3.

$\hat{\star} \nearrow$	$\mathcal{G}[1, M]$	$\mathcal{IG}[1, M]$
$\mathcal{G}[\mu, L]$	\mathcal{K} distribution $\mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)} \frac{\Gamma(M+s-1)}{M^{s-1}\Gamma(M)}$	Pearson VI $\mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)} \frac{\Gamma(M+1-s)}{M^{1-s}\Gamma(M)}$
$\mathcal{IG}[\mu, L]$	Pearson VI $\mu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)} \frac{\Gamma(M+s-1)}{M^{s-1}\Gamma(M)}$	\mathcal{IK} distribution $\mu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)} \frac{\Gamma(M+1-s)}{M^{1-s}\Gamma(M)}$
$\hat{\star}^{-1} \nearrow$	$\mathcal{G}[1, M]$	$\mathcal{IG}[1, M]$
$\mathcal{G}[\mu, L]$	Pearson I $\mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)} \frac{M^{s-1}\Gamma(M)}{\Gamma(M+s-1)}$	Bessel $\mu^{s-1} \frac{\Gamma(L+s-1)}{L^{s-1}\Gamma(L)} \frac{M^{1-s}\Gamma(M)}{\Gamma(M+1-s)}$
$\mathcal{IG}[\mu, L]$	Bessel $\mu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)} \frac{M^{s-1}\Gamma(M)}{\Gamma(M+s-1)}$	Inverse Pearson I $\mu^{s-1} \frac{\Gamma(L+1-s)}{L^{1-s}\Gamma(L)} \frac{M^{1-s}\Gamma(M)}{\Gamma(M+1-s)}$

TABLE 3

Second and third order log-cumulants of the different distributions obtained by Mellin convolution (direct and inverse) of the gamma distribution and the inverse gamma distribution (cf. Table 2). To simplify the presentation, only the second and third order log-cumulants are included in the table. The distributions whose names are typeset in boldface correspond to new analytical expressions.

$\hat{\star} \nearrow$	$\mathcal{G}[1, M]$	$\mathcal{IG}[1, M]$	$\hat{\star}^{-1} \nearrow$	$\mathcal{G}[1, M]$	$\mathcal{IG}[1, M]$
$\mathcal{G}[1, L]$	\mathcal{K} distribution $\Psi(1, L) + \Psi(1, M)$ $\Psi(2, L) + \Psi(2, M)$	Pearson VI $\Psi(1, L) + \Psi(1, M)$ $\Psi(2, L) - \Psi(2, M)$	$\mathcal{G}[1, L]$	Pearson I $\Psi(1, L) - \Psi(1, M)$ $\Psi(2, L) - \Psi(2, M)$	Bessel $\Psi(1, L) - \Psi(1, M)$ $\Psi(2, L) + \Psi(2, M)$
$\mathcal{IG}[\mu, L]$	Pearson VI $\Psi(1, L) + \Psi(1, M)$ $-\Psi(2, L) + \Psi(2, M)$	\mathcal{IK} distribution $\Psi(1, L) + \Psi(1, M)$ $-\Psi(2, L) - \Psi(2, M)$	$\mathcal{IG}[\mu, L]$	Bessel $\Psi(1, L) - \Psi(1, M)$ $-\Psi(2, L) - \Psi(2, M)$	Inverse Pearson I $\Psi(1, L) - \Psi(1, M)$ $-\Psi(2, L) + \Psi(2, M)$

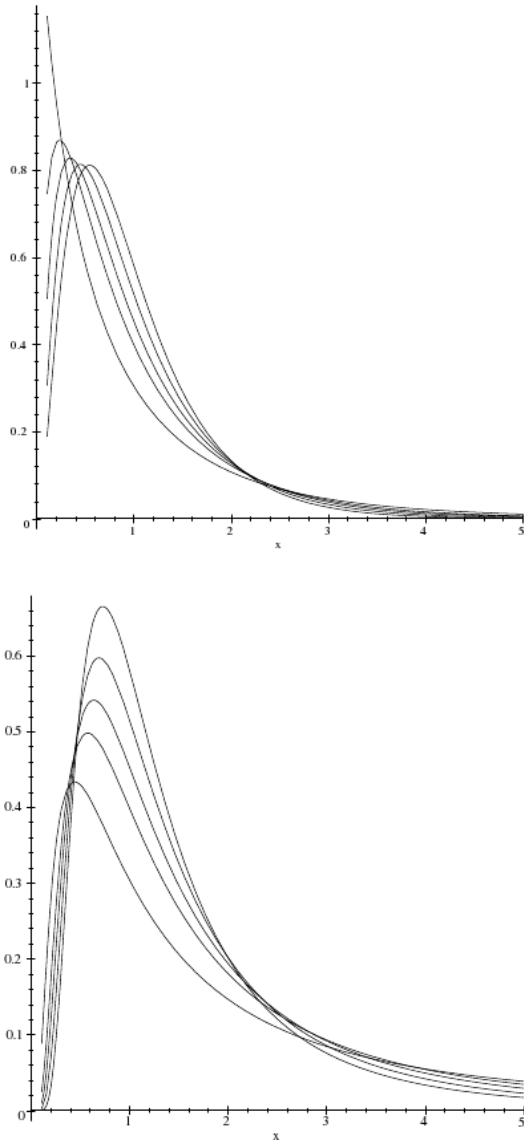


Fig. 7. The \mathcal{K} distribution (top) and the inverse \mathcal{K} distribution (bottom).

where the last expression is found in the Mellin transform tables. Figure 6 allows a comparison between the \mathcal{IP} distribution and the \mathcal{IPI} distribution for the same set of parameters. Recall that the Pearson type I distribution is expressed as:

$$\mathcal{IP}[\xi, L, M] = \frac{L}{M\xi} \frac{\Gamma(M)}{\Gamma(L)\Gamma(M-L)} \left(\frac{Lu}{M\xi}\right)^{L-1} \left(1 - \frac{Lu}{M\xi}\right)^{M-L-1}$$

with $u \leq \frac{M}{L\xi}$ and $M \geq L + 1$.

Curiously, this distribution is never mentioned in the classical works [9], [19], whereas they characterise the gamma distribution and the inverse gamma distribution separately. Moreover, the inverse Pearson type I distribution has the peculiar property of being localised, in the (β_1, β_2) diagram, between the gamma distribution

and the inverse gamma distribution, that is, exactly where the Pearson type VI solution is found. Indeed, the case $M \rightarrow \infty$ corresponds to the inverse gamma distribution and the zone corresponding to the inverse Pearson type I distribution cannot have ambiguities, whereas the (κ_3, κ_2) diagram separate well between the solutions of the Pearson system.

- The inverse \mathcal{K} distribution, which is also uncommon, is expressed as:

$$\mathcal{IK}[\mu, L, M](u) = \frac{1}{L\Gamma(L)M\Gamma(M)} \frac{2}{\mu} \left(\frac{LM\mu}{u}\right)^{\frac{M+L}{2}+1} \times K_{M-L} \left[2 \left(\frac{LM\mu}{u}\right)^{\frac{1}{2}} \right] \quad (48)$$

Figure 7 allows a comparison between the \mathcal{K} distribution and the inverse \mathcal{K} distribution for the same set of parameters. As for the \mathcal{K} distribution, the modelling through the Mellin convolution makes it easy to show that the inverse \mathcal{K} distribution tends to an inverse gamma distribution as one of the shape parameters (L or M) goes to infinity. One thus has a three parameter distribution which is heavy-tailed.

- The combinations $\mathcal{G} \hat{\star}^{-1} \mathcal{IG}$ and $\mathcal{IG} \hat{\star}^{-1} \mathcal{G}$ have known analytical solutions that include Bessel functions. However, they are not probability distributions, as the condition $p_x(u) \geq 0$ is not satisfied.

We see that insightful interpretations can be made based on the second and third order log-cumulants, $\tilde{\kappa}_2$ and $\tilde{\kappa}_3$. By the simple analysis of these entities, we can effectively get an idea about the flexibility of a certain distribution compared to the generalised gamma distribution, its inverse, and the other distributions that cover the log-cumulant space: the \mathcal{K} distribution and its inverse, and the solutions of the Pearson system. The analysis of the second and third order log-cumulants can also account for more complex models. We shall see that the same diagram can be used to analyse an additive mixture of gamma distributions, and propose an original and simple practical method to determine the model parameters.

4.4 Additive Mixture of Gamma Distributions

Additive mixtures of gamma distributions are important practical modelling tools (in particular for SAR imagery). Contrarily to the Gaussian case, a unimodal pdf is generally obtained, except when the (two) initial distributions are very different (see Figure 8). However, we will show that there exists a simple solution to determine the parameters of the mixture by analysing this problem aided by second and third order log-cumulants.

Consider the following mixture of gamma distributions:

$$\lambda \mathcal{G}[\mu, L] + \lambda' \mathcal{G}[\mu', L]$$

with $\lambda \geq 0$, $\lambda' \geq 0$ and $\lambda + \lambda' = 1$. In this model, L has the same value for the two gamma distributions.

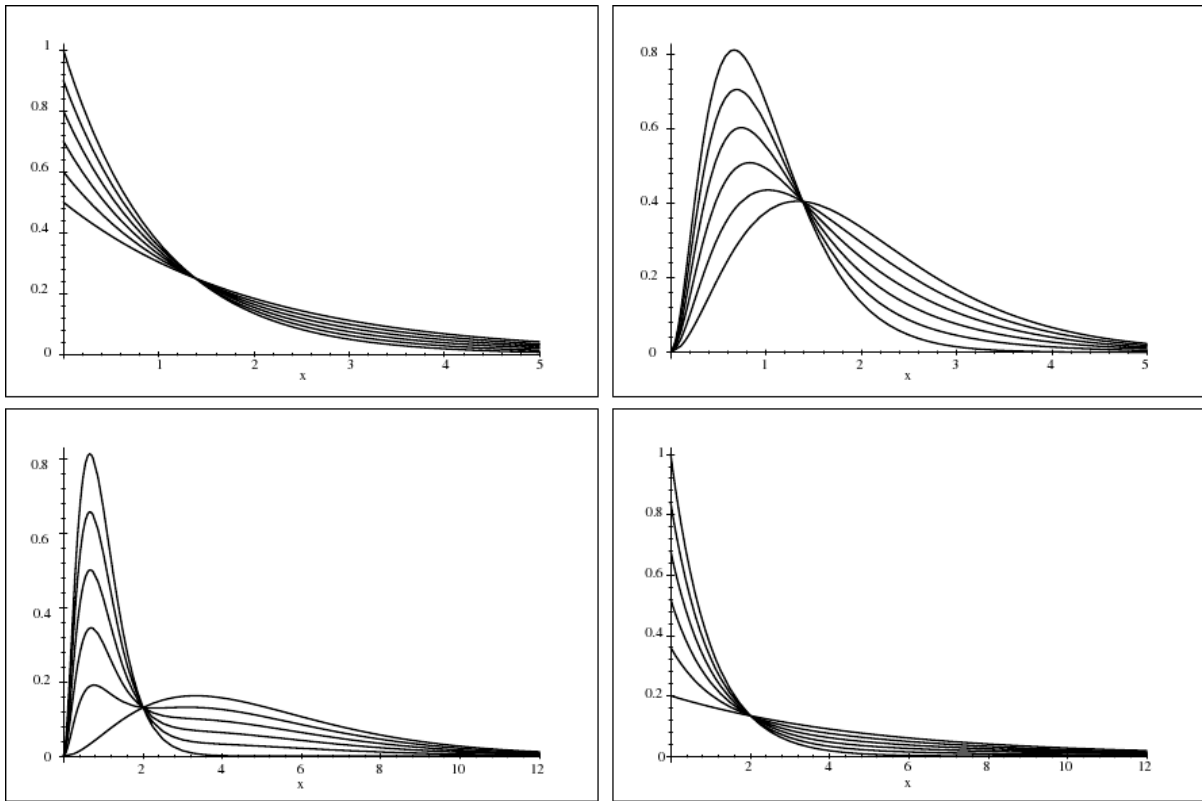


Fig. 8. Examples of additive mixtures of gamma distributions. The left column shows distributions with $L = 1$, and the right column with $L = 3$. In the first row, $\rho = 2$ ($\mu = 1$ and $\mu' = 2$). In the second row, $\rho = 5$ ($\mu = 1$ and $\mu' = 5$). λ takes the values 0, 0.2, 0.4, 0.6, 0.8 and 1.

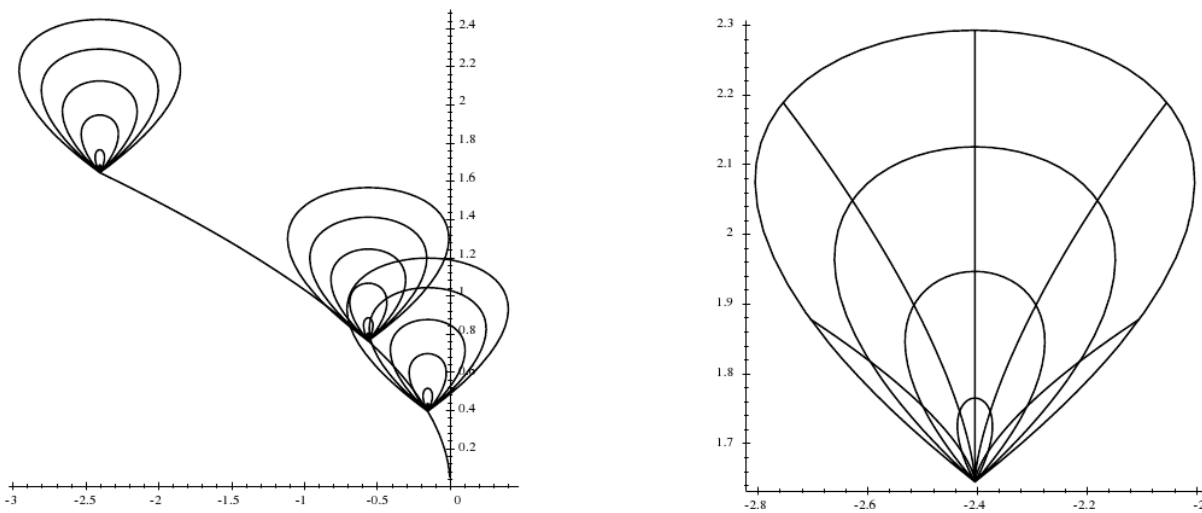


Fig. 9. (κ_2, κ_3) diagram for a mixture of gamma distributions described by the parameters λ (mixing proportion) and ρ (ratio of component means). To the left, for λ varied between 0 and 1, three diagrams are traced out corresponding to several values of ρ for three values of L ($L = 1, 3$ and 5). In the same plot, the gamma distribution is represented by a line which spans $L \in [1, \infty]$. To the right, for a fixed value of $L = 1$, one separately varies ρ between 0 and 5 (five curves, with λ taking the values 0.1, 0.3, 0.5, 0.7 and 0.9) and λ between 0 and 1 (four closed curves, with ρ taking values 2, 3, 4 and 5), placing both diagrams in the same figure.

The model can be rewritten by defining the variable ρ such that $\mu' = \rho\mu$, which makes it possible to write the mixture in the following form:

$$\lambda\mathcal{G}[\mu, L] + (1 - \lambda)\mathcal{G}[\rho\mu, L] \quad (49)$$

The mixture is then defined by a gamma distribution $\mathcal{G}[\mu, L]$ (corresponding to only one of the mixture components) and two parameters describing the mixture: λ and ρ . The second kind characteristic function is written:

$$\phi(s) = (\lambda + (1 - \lambda)\rho^{s-1})\mu^{s-1} \frac{\Gamma(L + s - 1)}{L^{s-1}\Gamma(L)}$$

Based on this expression, calculation of the log-cumulants can be carried out directly, giving the following expressions:

$$\begin{aligned} \tilde{\kappa}_{x(1)} &= \Psi(L) - \log L + \log \mu + (1 - \lambda) \log \rho \\ \tilde{\kappa}_{x(2)} &= \Psi(1, L) + \log(\rho)^2 \lambda(1 - \lambda) \\ \tilde{\kappa}_{x(3)} &= \Psi(2, L) + \log(\rho)^3 \lambda(1 - \lambda)(2\lambda - 1) \end{aligned}$$

We remark that starting from the second order, the log-cumulants do not depend on μ , and they have the same values as the standard gamma distribution for the limiting values $\lambda = 0$ and $\lambda = 1$.

We assume that the entity L is known (L can be perceived as a function of the instrument, thus it will be known by the processor). This leads to:

$$\begin{aligned} \overline{\tilde{\kappa}_{x(2)}} &= \tilde{\kappa}_{x(2)} - \Psi(1, L) \\ \overline{\tilde{\kappa}_{x(3)}} &= \tilde{\kappa}_{x(3)} - \Psi(2, L) \end{aligned}$$

Parameters λ and ρ are then given by the solutions of quadratic equation, which gives:

$$\begin{aligned} \lambda &= \frac{1}{2} \left(1 \pm \frac{\overline{\tilde{\kappa}_{x(3)}}}{\sqrt{4\overline{\tilde{\kappa}_{x(2)}}^3 + \overline{\tilde{\kappa}_{x(3)}}^2}} \right) \\ \rho &= e^{\frac{\sqrt{4\overline{\tilde{\kappa}_{x(2)}}^3 + \overline{\tilde{\kappa}_{x(3)}}^2}}{\overline{\tilde{\kappa}_{x(2)}}}} \end{aligned}$$

The evolution of the different parameters in the $(\tilde{\kappa}_3, \tilde{\kappa}_2)$ diagram is shown in Figure 9. It is interesting to notice that the shape of these curves does not depend on L .

If the obtained results are compared with those found in the literature, it is noticed that this approach relies on one assumption only: knowledge of the parameter L , while analyses of mixtures by classical methods require the additional knowledge of μ [20].

4.5 Positive α -Stable Distributions

We will now apply the methodology proposed in this article to the case of a 'heavy-tailed' distribution, for which neither the analytical form of the pdf nor moments from a certain order and upwards are known. This prohibits the method of moments.

A positive α -stable distribution has a pdf characterised by two parameters: α and γ . It cannot in general be

defined, other than by its characteristic function $\Phi(\nu)$, which is written (according to Pierce [17]) as:

$$\Phi(\nu) = e^{-\gamma|\nu|^\alpha(1+j\operatorname{sgn}(\nu)\tan(\frac{\alpha\pi}{2}))}$$

with

$$\operatorname{sgn}(\nu) = \begin{cases} 1, & \nu > 0 \\ 0, & \nu = 0, \\ -1, & \nu < 0 \end{cases} \quad 0 < \alpha < 1, \quad \gamma > 0.$$

Except for certain particular values of α , the analytical expression of the distribution is not known.

One can nevertheless express the moments of this distribution (including fractional ones) as:

$$m_\nu = \frac{\gamma^{\frac{\nu}{\alpha}} \sin(\pi\nu)\Gamma(\nu + 1) \left(1 + \left(\tan\left(\frac{\pi\alpha}{2}\right)\right)^2\right)^{\frac{\nu}{2\alpha}}}{\alpha \sin\left(\frac{\pi\nu}{\alpha}\right) \Gamma\left(1 + \frac{\nu}{\alpha}\right)} \quad (50)$$

These moments are only defined for $\nu < \alpha < 1$, which means that even the first moment is not defined. This is evidently a heavy-tailed distribution.

It is nevertheless possible, by an analytical continuation, to derive the second order characteristic function, which is written:

$$\phi(s) = \frac{\gamma^{\frac{s-1}{\alpha}} \sin(\pi(s-1))\Gamma(s) \left(1 + \left(\tan\left(\frac{\pi\alpha}{2}\right)\right)^2\right)^{\frac{s-1}{2\alpha}}}{\alpha \sin\left(\frac{\pi(s-1)}{\alpha}\right) \Gamma\left(1 + \frac{s-1}{\alpha}\right)}$$

It is seen that this function is well defined in a complex neighbourhood around the value $s = 1$. The existence theorem in Section 2.5 thus confirms that the log-moments and log-cumulants of all orders exist, whereas the moments of orders $\nu \geq \alpha$ are not defined.

Even though the analytical form is rather complicated, it is still possible to obtain simple expressions of the log-cumulants. Note that these expressions are only analytical continuations because the gamma functions in the derivatives of $\phi(s)$ have discontinuities at $s = 1$, a value at which they must be evaluated when calculating the log-cumulants. It is then necessary to study the limit at $s = 1$ in order to obtain the analytical expressions. The following results were established with assistance of mathematical computation software Maple, as a result of lengthy analytical developments:

$$\begin{aligned} \tilde{\kappa}_1 &= \frac{(1 - \alpha)\Psi(1)}{\alpha} + \frac{-\log\left(\cos\left(\frac{\pi\alpha}{2}\right)\right)}{\alpha} + \frac{\log\gamma}{\alpha} \\ \tilde{\kappa}_2 &= \frac{(1 - \alpha^2)}{\alpha^2} \Psi(1, 1) \\ \tilde{\kappa}_3 &= \frac{\alpha^3 - 1}{\alpha^3} \Psi(2, 1) \end{aligned} \quad (51)$$

These expressions, that eventually appear as rather simple, illustrate well the strength of our new approach. The two parameters of the distribution are easily derived since:

- The parameter α can be estimated from the second order log-cumulant as:

$$\alpha = \sqrt{\frac{\Psi(1, 1)}{\Psi(1, 1) + \tilde{\kappa}_2}}$$

- When α is known, γ can be obtained from the first log-cumulant as:

$$\gamma = e^{\alpha \tilde{\kappa}_1 - (1-\alpha)\Psi(1) + \log(\cos(\frac{\pi\alpha}{2}))}$$

- By combining the previous expressions, one can also write γ in terms of $\tilde{\kappa}_1$ and $\tilde{\kappa}_2$:

$$\gamma = \exp\left(\sqrt{\frac{\Psi(1,1)}{\Psi(1,1) - \tilde{\kappa}_2}} \tilde{\kappa}_1 - \left[1 - \sqrt{\frac{\Psi(1,1)}{\Psi(1,1) - \tilde{\kappa}_2}}\right] \Psi(1) + \log\left[\cos\left(\frac{\pi}{2} \sqrt{\frac{\Psi(1,1)}{\Psi(1,1) - \tilde{\kappa}_2}}\right)\right]\right)$$

We note that the Mellin transform sheds, on the theoretical side, a new and recent light on the heavy-tailed distributions [21].

4.6 A Generalisation of the Rayleigh Distribution

Another example drawn from the α -stable distributions inspired Kuruoğlu and Zerubia to propose a generalisation of the Rayleigh distribution [18]. The pdf has two parameters (α and γ) and its analytical expression is given by the following integral equation:

$$p(u) = u \int_0^\infty v e^{-\gamma v^\alpha} J_0(uv) dv \quad (52)$$

where J_0 is the Bessel function of the second kind. This distribution falls into the heavy-tailed category, since its moments are not defined from a certain order and upwards, with this order given as: $\min(\alpha, 2)$.

To calculate its second kind characteristic function, two approaches are possible:

- The first search for the Mellin transform of this expression [14] led to the following result:

$$\phi(s) = \frac{2^s \Gamma(\frac{s+1}{2}) \gamma^{\frac{s-1}{\alpha}} \Gamma(\frac{1-s}{\alpha})}{\Gamma(\frac{1-s}{2}) \alpha} \quad (53)$$

- A second approach consists of rewriting (52) on the form of a Mellin correlation:

$$p(u) = u \left(J_0(u) \hat{\otimes} \left(e^{-\gamma u^\alpha} \right) \right)$$

By using the property in (22), (53) is immediately retrieved.

It can be noted that at $s = 1$, the second kind characteristic function goes in the limit to the value 1, since

$$\lim_{s \rightarrow 1} \frac{\Gamma(\frac{1-s}{\alpha})}{\Gamma(\frac{1-s}{2})} = \frac{\alpha}{2}$$

It is thus a valid pdf.

In the vicinity of $s = 1$, this function is defined for $s < 1 + \min(\alpha, 2)$ and for $s > -1$. It is thus well defined in a vicinity of $s = 1$, hence it is legitimate to calculate its log-moments and log-cumulants. Note that the case

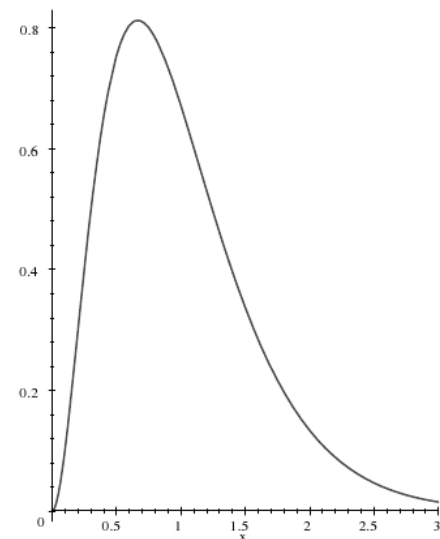
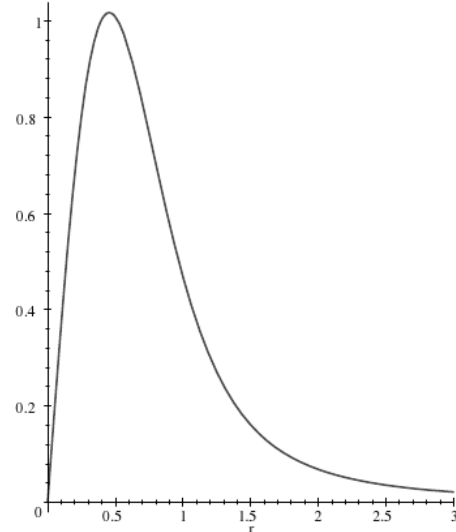
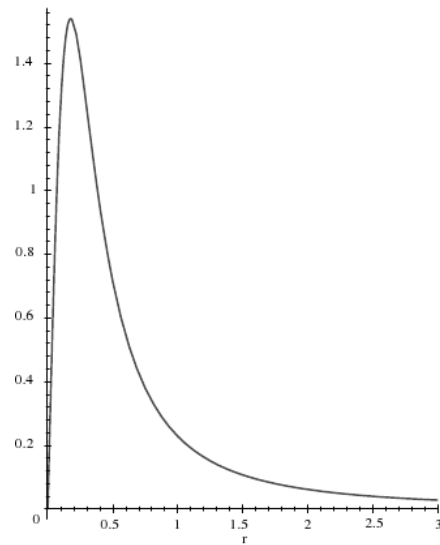


Fig. 10. Generalised Rayleigh distribution [18]: $\alpha = 1$, $\alpha = 1.5$ and $\alpha = 2$.

$\alpha = 2$ gives the Rayleigh distribution and the case $\alpha = 1$ gives the distribution

$$p(u) = \frac{\gamma u}{(u^2 + \gamma^2)^{\frac{3}{2}}}.$$

Figure 10 shows this distribution for $\alpha = 1$, $\alpha = 1.5$ and $\alpha = 2$ (the Rayleigh distribution).

Although the analytical form is rather complicated, it is possible to formulate the first and second order log-cumulants of the probability distribution. Again, the expressions obtained are analytical continuations because of discontinuities in the gamma functions. Also in this case, Maple was used to establish following expressions:

$$\begin{aligned}\tilde{\kappa}_1 &= -\Psi(1) \frac{1-\alpha}{\alpha} + \log\left(2\gamma^{\frac{1}{\alpha}}\right) \\ \tilde{\kappa}_2 &= \frac{\Psi(1,1)}{\alpha^2}.\end{aligned}$$

The equation system obtained is easy to handle. The distribution parameters are easily retrieved from the first two log-cumulants:

$$\begin{aligned}\alpha &= \sqrt{\frac{\Psi(1,1)}{\tilde{\kappa}_2}} \\ \gamma &= e^{\alpha \tilde{\kappa}_1 - \Psi(1)(1-\alpha)}.\end{aligned}$$

It is verified that for $\alpha = 2$, the log-cumulants of the Rayleigh distribution are retrieved (with $\mu = 2\sqrt{\gamma}$):

$$\begin{aligned}\tilde{\kappa}_1 &= \frac{1}{2}\Psi(1) + \log(2\sqrt{\gamma}) \\ \tilde{\kappa}_2 &= \frac{1}{4}\Psi(1,1).\end{aligned}$$

5 PARAMETER ESTIMATION

The proposal of a new methodology to evaluate the parameters of a probability distribution requires us to compare the results with those obtained by traditional methods in a realistic setting where N samples are available. In order to decide which method is the preferred one, it is important to establish the variance of the estimators. The goal of this section is to carry out an exhaustive comparison for a schoolbook example: the gamma distribution, for which the method of moments, the method of log-moments, and the method of maximum likelihood are applicable.

5.1 Traditional Methods: Method of Moments Estimation and Maximum Likelihood Estimation

5.1.1 Experimental Framework

Assume that a probability distribution is described by p parameters: α_j , $j \in [1, p]$. Estimation of the parameters describing this distribution is commonly performed by the two following approaches:

- The *method of moments* (MoM) consists of calculating the sample moments of order 1 to p in order to obtain a system of p equations in p unknown pdf

parameters. If N samples are available, x_i , $i \in [1, N]$, the r th order sample moment is expressed simply as

$$m_r = \frac{1}{N} \sum_{i=1}^N x_i^r.$$

In order to determine p parameters, it is necessary that all moments up to order p exist, which can pose a problem for instance for distributions with heavy tails. It is also possible to use fractional moments (like FLOM [3]), or lower (and even negative) order moments, whose possible existence is justified in section 2.2 [4], to obtain an equation system which can be solved. Note, however, that the expressions sometimes prove impossible to invert analytically, and the system may also be difficult to invert numerically.

- The *maximum likelihood* approach consists of regarding the N samples x_i as N independent realisations of the distribution which one seeks to estimate, so that they maximise the expression

$$\prod_{i=1}^N p_x(x_i)$$

or, equivalently,

$$\sum_{i=1}^N \log(p_x(x_i)).$$

With these expressions representing a maximum, calculation of partial derivatives for each parameter then makes it possible to obtain another system of p equations in p unknowns:

$$\frac{\partial \left(\sum_{i=1}^N \log(p_x(x_i)) \right)}{\partial \alpha_j} = 0. \quad (54)$$

The solution relies on the existence of the partial derivatives, which can pose a problem, as for the \mathcal{K} distribution [16].

5.1.2 Estimator Variance

With several applicable methods available, we must compare them to select the one which is likely to give the user the most reliable results. A natural approach is to seek the method which provides minimum variance for the estimator of a given parameter, knowing that one has a finite number of N samples.

It is known that for the distributions of the exponential family, maximum likelihood estimators attain the Cramer-Rao boundary. Provided that the p partial derivatives in Eq. (54) can be calculated analytically, and that the system of equations can be solved, one obtains p estimators whose variances are minimal. However, many existing distributions (such as the \mathcal{K} distribution) do not have analytical expressions for all partial derivatives,

which then renders the method of maximum likelihood inadequate.

In this case, the use of the method of moments is justified, even if the estimator variance thus obtained is higher. The variance of estimators obtained by the method of moments can be derived through an approach suggested by Kendall [19]. Let m_1 and m_2 be the estimates of the first two moments, and $g(m_1, m_2)$ a function depending only on these two entities. We seek to calculate the variance of the function $g(m_1, m_2)$ by linearising it and writing it as a first-order expansion around the values of the theoretical moments, $m_{0,1}$ and $m_{0,2}$:

$$g(m_1, m_2) = g(m_{0,1}, m_{0,2}) + (m_1 - m_{0,1}) \frac{\partial g}{\partial m_1}(m_{0,1}, m_{0,2}) + (m_2 - m_{0,2}) \frac{\partial g}{\partial m_2}(m_{0,1}, m_{0,2}).$$

After having verified that the $\partial g / \partial m_i$ are not both zero in the point $(m_{0,1}, m_{0,2})$, the variance of g is established as the quadratic error between $g(m_1, m_2)$ and $g(m_{0,1}, m_{0,2})$ due to the following formula [19, Eq. (10.12)]:

$$\begin{aligned} & \text{Var}\{g(m_1, m_2)\} \\ &= \text{E} \left\{ [g(m_1, m_2) - g(m_{0,1}, m_{0,2})]^2 \right\} \\ &= \text{E} \left\{ \left[(m_1 - m_{0,1}) \frac{\partial g}{\partial m_1}(m_{0,1}, m_{0,2}) + (m_2 - m_{0,2}) \frac{\partial g}{\partial m_2}(m_{0,1}, m_{0,2}) \right]^2 \right\} \\ &= \frac{\partial g}{\partial m_1}(m_{0,1}, m_{0,2})^2 \text{Var}\{m_1\} \\ & \quad + \frac{\partial g}{\partial m_2}(m_{0,1}, m_{0,2})^2 \text{Var}\{m_2\} \\ & \quad + 2 \frac{\partial g}{\partial m_1}(m_{0,1}, m_{0,2}) \frac{\partial g}{\partial m_2}(m_{0,1}, m_{0,2}) \\ & \quad \quad \times \text{Cov}\{m_1, m_2\}. \end{aligned} \quad (55)$$

The method can obviously be generalised to functions utilising moments m_i of order i . The definition of the covariance matrix allows us to write:

$$\begin{aligned} \text{Var}\{m_i\} &= \frac{1}{N} (m_{2i} - m_i^2) \\ \text{Cov}\{m_i, m_j\} &= \frac{1}{N} (m_{i+j} - m_i m_j) \end{aligned}$$

5.2 Method of Log-Moments

We propose in this article a new method for analysis of pdfs defined on \mathbb{R}^+ based on log-moments and log-cumulants. We will see in this section how to implement it and how to calculate the variance of the estimators obtained.

5.2.1 Description

The method of log-moments (MoLM) consists of calculating estimates of log-moments and log-cumulants in order to obtain a system of p equations in p unknowns

(the parameters of the pdf). Assume that we have N samples $x_i, i \in [1, N]$ from the distribution to be estimated. The estimate of the p th order log-moment is expressed simply as

$$\tilde{m}_p = \frac{1}{N} \sum_{i=1}^N \log x_i^p.$$

To determine p parameters, it is necessary to check in advance that the log-moments up till order p exist. This is in general true, as stated by the theorem of existence, which has been verified for the distributions generally used in signal and image processing.

5.2.2 Estimator Variance

Since we use a logarithmic scale, the criterion of the quadratic error (applied in Eq. (55)), is replaced by another criterion which we will call "normalised quadratic error", E_{nq} , which is in fact the quadratic error calculated on a logarithmic scale:

$$E_{nq} = \text{E} \left\{ \left(\log \left(\frac{x}{y} \right) \right)^2 \right\}.$$

In the same spirit, we introduce the second kind variance and covariance, $\tilde{\text{Var}}$ and $\tilde{\text{Cov}}$, on the form

$$\begin{aligned} \tilde{\text{Var}}\{\tilde{m}_i\} &= \text{E} \left\{ [(\log x)^i - \tilde{m}_i]^2 \right\} = \frac{1}{N} (\tilde{m}_{2i} - \tilde{m}_i^2) \\ \tilde{\text{Cov}}\{\tilde{m}_i, \tilde{m}_j\} &= \text{E} \left\{ [(\log x)^i - \tilde{m}_i] [(\log x)^j - \tilde{m}_j] \right\} \\ &= \frac{1}{N} (\tilde{m}_{i+j} - \tilde{m}_i \tilde{m}_j) \end{aligned}$$

where \tilde{m}_i is the i th order log-moment.

With this new approach, and taking the preceding step as starting point, let the function g be expressed in terms of the first two estimated log-moments as $g(\tilde{m}_1, \tilde{m}_2)$. Then g can be expanded around the first two theoretical log-moments, $\tilde{m}_{0,1}$ and $\tilde{m}_{0,2}$, as

$$\begin{aligned} g(\tilde{m}_1, \tilde{m}_2) &= g(\tilde{m}_{0,1}, \tilde{m}_{0,2}) + (\tilde{m}_1 - \tilde{m}_{0,1}) \frac{\partial g}{\partial \tilde{m}_1}(\tilde{m}_{0,1}, \tilde{m}_{0,2}) \\ & \quad + (\tilde{m}_2 - \tilde{m}_{0,2}) \frac{\partial g}{\partial \tilde{m}_2}(\tilde{m}_{0,1}, \tilde{m}_{0,2}). \end{aligned}$$

After verifying that the $\partial g / \partial \tilde{m}_i$ are not both zero in $(\tilde{m}_{0,1}, \tilde{m}_{0,2})$, the variance of g is established by the same formula applied in the previous section.

$$\begin{aligned} \text{Var}\{g(\tilde{m}_1, \tilde{m}_2)\} &= \text{E} \left\{ [g(\tilde{m}_1, \tilde{m}_2) - g(\tilde{m}_{0,1}, \tilde{m}_{0,2})]^2 \right\} \\ &= \text{E} \left\{ \left[(\tilde{m}_1 - \tilde{m}_{0,1}) \frac{\partial g}{\partial \tilde{m}_1}(\tilde{m}_{0,1}, \tilde{m}_{0,2}) + (\tilde{m}_2 - \tilde{m}_{0,2}) \frac{\partial g}{\partial \tilde{m}_2}(\tilde{m}_{0,1}, \tilde{m}_{0,2}) \right]^2 \right\} \\ &= \frac{\partial g}{\partial \tilde{m}_1}(\tilde{m}_{0,1}, \tilde{m}_{0,2})^2 \tilde{\text{Var}}\{\tilde{m}_1\} + \frac{\partial g}{\partial \tilde{m}_2}(\tilde{m}_{0,1}, \tilde{m}_{0,2})^2 \tilde{\text{Var}}\{\tilde{m}_2\} \\ & \quad + 2 \frac{\partial g}{\partial \tilde{m}_1}(\tilde{m}_{0,1}, \tilde{m}_{0,2}) \frac{\partial g}{\partial \tilde{m}_2}(\tilde{m}_{0,1}, \tilde{m}_{0,2}) \times \tilde{\text{Cov}}\{\tilde{m}_1, \tilde{m}_2\}. \end{aligned}$$

As in the previous, this method can obviously be generalised to functions of the moments \tilde{m}_i of unspecified order i .

5.3 The Gamma Distribution Case

We will use the gamma distribution as an example to compare the available methods. This distribution is not heavy tailed, thus the method of moments can be used, as well as the method of log-moments. The partial derivatives with respect to the parameters are known, which makes it possible to apply maximum likelihood estimation.

5.3.1 Variance of the Gamma Distribution Parameter Estimators with the Method of Moments

The method of moments (MoM) utilises the first two moments to deduce estimates of L and μ (Eqs. (30) and (31)). The method of Kendall, presented in Section 5.1.2, gives the following variance for the estimators of μ and L :

$$\text{Var}_{MoM}\{\hat{\mu}\} = \frac{\mu^2}{NL} \quad (56)$$

$$\text{Var}_{MoM}\{\hat{L}\} = \frac{2L(L+1)}{N} \quad (57)$$

5.3.2 Variance of the Gamma Distribution Parameter Estimators with the Method of Log-Moments

The parameter L is derived from (33) as

$$\Psi(1, L) = \tilde{\kappa}_{x(2)}$$

which can be rewritten as a function of $(\tilde{m}_1, \tilde{m}_2)$:

$$\Psi(1, L) = \tilde{m}_2 - \tilde{m}_1^2.$$

One then carries out the limited expansion proposed in the previous, which requires the use of implicit differentiation. Although the expression brings into play the first to fourth order log-moments, the result can be simplified and we obtain:

$$\text{Var}_{MoLM}\{\hat{L}\} = \frac{1}{N} \frac{\Psi(3, L) + 2\Psi(1, L)^2}{\Psi(2, L)^2}. \quad (58)$$

Figure 11 (left panel) presents the ratio of the standard deviation for the variance of MoLM estimate of \hat{L} to the standard deviation for the variance of the MoM estimate of \hat{L} . The whole motivation for using the new method is evident for low values of L , where the improvement approaches 30%. When the variance of the different estimators is fixed, this results in the same amount of shrinking of the analysis window, and therefore a better spatial localisation of the estimate.

For the parameter μ , the calculation is much more elaborate, and we finally arrive at the following expression, whose interpretation is not simple, but which can easily be implemented numerically:

$$\begin{aligned} \text{Var}_{MoLM}\{\hat{L}\} &= -\frac{1}{N} \frac{\mu^2}{L^2 \Psi(2, L)^2} \\ &\times [2\Psi(1, L)L\Psi(3, L) - \Psi(1, L)^2 L^2 \Psi(3, L) \\ &\quad + 4\Psi(1, L)^3 L - 2\Psi(1, L)^4 L^2 - 2L\Psi(2, L)^2 \\ &\quad + \Psi(1, L)L^2 \Psi(2, L)^2 - 2\Psi(1, L)^2 - \Psi(3, L)] \end{aligned}$$

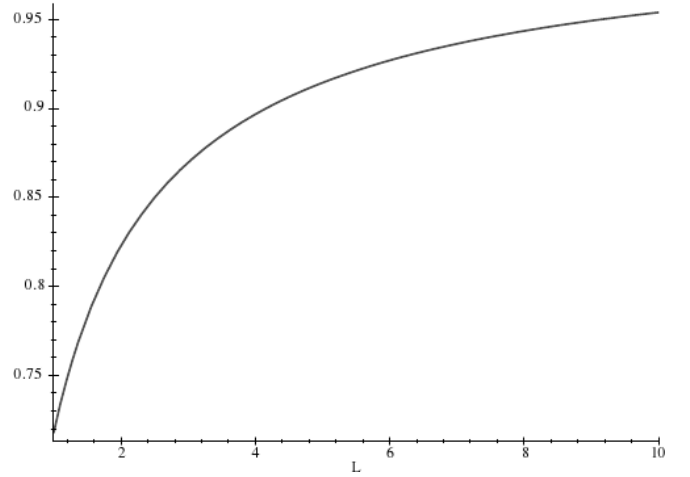


Fig. 12. Gamma distribution: Comparison of the variance of the estimator for μ by the method of log-moments with the method of moments. The curve represents the ratio of the standard deviations for values of L between 1 and 10.

Also Figure 12 presents the ratio of the standard deviation of $\hat{\mu}$ calculated by the MoLM to the standard deviation of $\hat{\mu}$ calculated by the MoM.

It can be noted that the MoM provides better results for low values of L . Recall moreover that this is also the maximum likelihood estimator and thus attains minimum variance (i.e. the Cramer-Rao bound).

5.3.3 Variance of the Gamma Distribution Parameter Estimators with the Method of Lower Order Moments

The existence of the second kind characteristic function for values of s lower than 1 justifies the use of the method of lower order moments (MoLOM), i.e. negative ones. In the case of the gamma distribution, it is known that the lower order moments exist for $\nu > -L$. For a given value of ν it is verified that $\nu > -L$ and using the three moments μ_ν , $\mu_{\nu+1}$ and $\mu_{\nu+2}$ it is easy to show that $\hat{\mu}$ and \hat{L} can be derived from the relation:

$$\begin{aligned} \hat{\mu} &= \frac{\hat{m}_{\nu+1}}{\hat{m}_\nu} (1 + \nu) - \nu \frac{\hat{m}_{\nu+2}}{\hat{m}_{\nu+1}} \\ \hat{L} &= \frac{1}{\hat{m}_\nu \left(\frac{\hat{m}_{\nu+2}}{\hat{m}_{\nu+1}} \right) - 1} - \nu \end{aligned}$$

For $\nu = 0$, this reduces to the MoM (Eqs. (30) and (31)).

The variances of the estimators for μ and L can be established by the method of Kendall, used in Section 5.3.1 (for the MoM). For L , the following expression is obtained:

$$\begin{aligned} \text{Var}_{MoLOM}\{\hat{L}\} &= \frac{1}{N} \frac{\Gamma(L)\Gamma(2\nu+L)}{\Gamma(\nu+L)^2} \\ &\times (2L(L+1) + \nu[4L(\nu+2) + (\nu+4)(\nu+1)^2]) \end{aligned} \quad (59)$$

The problem with this relation is that it has a minimum for ν , which cannot be expressed explicitly as a function of L . The optimal values of ν must be calculated numerically. Table 4 gives some values of ν as a function

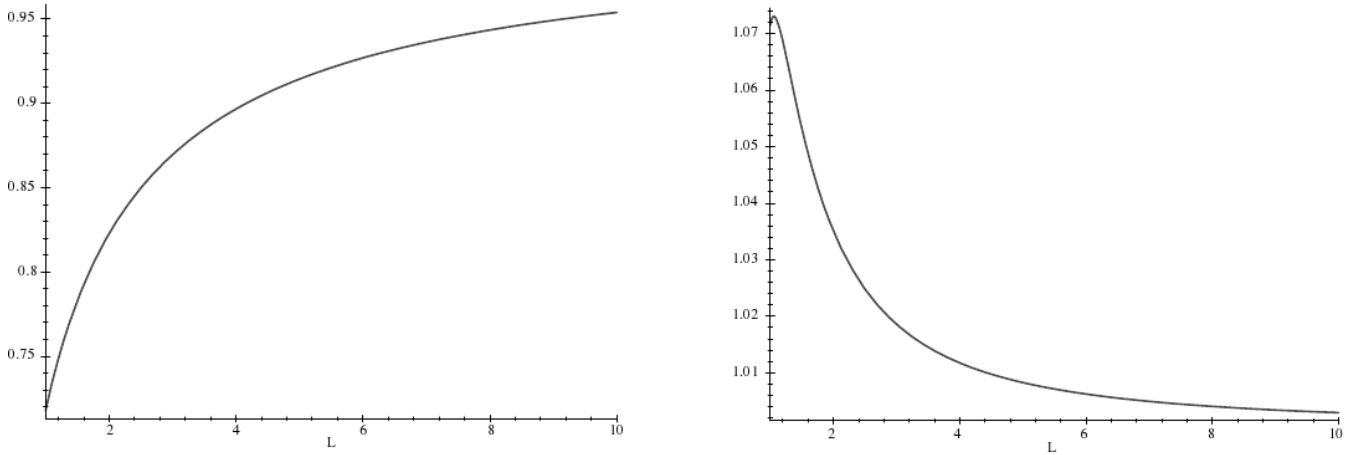


Fig. 11. To the left, a comparison of the variance of the estimator for L by the method of log-moments and the method of moments. To the right, a comparison of the variance of the estimators of L by the method of log-moments and the method of lower order moments (only for the value of $\nu = -0.35$). The curves represent the standard deviation for values of L between 1 and 10.

TABLE 4

Gamma distribution estimated with the method of lower order moments. Optimal values of the parameter ν that minimises the variance of \hat{L} as a function of L

L	ν_{opt}
1	-0.35
2	-0.44
3	-0.56
4	-0.59

of L . When information about L is absent, the choice of $\nu = -0.35$ seems to be a good compromise.

Figure 11 (right panel) presents the ratio of the standard deviation of \hat{L} calculated by the MoLM to the standard deviation of \hat{L} calculated by the MoLOM with $\nu = -0.35$. It is interesting to note that the MoLOM is slightly better than the MoLM. Nevertheless, if one wants to fully utilise this method, then one must know L to be able to choose the optimal value of ν . As the difference is altogether rather small, we promote the MoLM because it does not require us to determine a parameter in order to make optimal use of the method.

However, it is easily shown that minimum variance is obtained for $\nu = 0$, which is less than astonishing since this value corresponds to the maximum likelihood estimator.

5.3.4 Summary

We propose to summarise these results by posting in Table 5 the optimal window dimension for these three methods when we seek to reach an error of 10% for the estimate of the shape parameter L (i.e. the standard deviation is 10% of the value to be estimated). For

TABLE 5

Number of samples (and examples of the analysis window) needed to estimate the parameters L and μ of a gamma distribution with 10% error. The methods used are, for the shape parameter L , the method of moments (MoM), the method of lower order moments (MoLOM) with $\nu = -0.35$, and the method of log-moments (MoLM).

The Cramer-Rao bound (CRB) is calculated by the means of the Fisher information matrix. For μ , only the MoM is used.

Gamma distribution					
L	\hat{L}				$\hat{\mu}$
	MoM	MoLOM	MoLM	CRB	MoM
1	400	179	206	155	100
	20×20	13×13	14×14	12×12	10×10
2	300	189	203	172	50
	17×17	14×14	14×14	13×13	7×7
3	267	194	202	180	33
	16×16	14×14	14×14	13×13	6×6
5	240	197	201	187	20
	15×15	14×14	14×14	14×14	4×4
10	220	199	200	194	10
	14×14	14×14	14×14	14×14	3×3

the parameter μ , only the method of moments (which coincides with the maximum likelihood estimator) is used.

We first remark that for an identical relative error (10%), the estimate of L requires much more samples

TABLE 6

Number of samples (and examples of the analysis windows) needed to estimate the parameters σ and μ of a Gaussian distribution with 10% error.

Gaussian distribution		
σ	$\hat{\sigma}$	$\hat{\mu}$
1	200 14×14	100 10×10
0.707	200 14×14	50 7×7
0.577	200 14×14	33 6×6
0.447	200 14×14	20 4×4
0.316	200 14×14	10 3×3

than the estimate of μ . Secondly, the estimate of μ requires much less samples when L is large, i.e. the distribution is localised. Lastly, it is interesting to note a characteristic feature of the MoLM: It requires about the same number of samples for all values of L , whereas the MoM requires a much higher number of samples when L is small. From this, two remarks can be made:

- It can be shown that the variance of \hat{L} for the MoLM (Eq. (58)) is almost quadratic in L :

$$\frac{\Psi(3, L) + 2\Psi(1, L)^2}{\Psi(2, L)^2} \simeq 2L^2$$

Thus, if a constant relative error is sought, the number of samples is independent of L .

- It is interesting to analyse the same problem for the Gaussian distribution $\mathcal{N}[\mu, \sigma^2]$. It is easy to show that the variances of the estimators of μ and σ do not depend on σ for the MoM. They are written:

$$\begin{aligned} \text{Var}_{\mathcal{N}}(\mu) &= \frac{\sigma^2}{N} \\ \text{Var}_{\mathcal{N}}(\sigma) &= \frac{\sigma^2}{2N} \end{aligned}$$

By choosing Gaussian distributions with $\mu = 1$, the values of σ become comparable and an identical criterion, and the required window sizes can be calculated. These are included in Table 6.

It is seen that the MoM needs a constant number of samples to estimate the shape parameter σ , an analogy to the property of the MoLM for the gamma distribution.

To achieve this analysis, we calculate Fisher's infor-

mation matrix for the gamma distribution⁴:

$$\begin{bmatrix} \frac{L}{\mu^2} & 0 \\ 0 & \Psi(1, L) - \frac{1}{L} \end{bmatrix},$$

which allows the calculation of the Cramer-Rao bound, given in Table 5.

5.4 The Mixture of Gamma Distribution Case

The analytical calculation of the variance of the estimators in the mixture of gamma distributions case described in Section 4.4 does not pose any problem, except for the apparent complexity of the expressions obtained, whose length prohibits us from including them in a publication. Another possibility would be to assess them by numerical evaluation.

Table 7 presents the standard deviations of λ and ρ if the analysis is carried out in neighbourhood of 100 samples (a 10×10 window), for various values of ρ and λ .

Table 8 presents the optimal dimension of a square window which guarantees a maximum of 10% estimation error (where the error is defined as the ratio of the standard deviation to the estimated value). Note that for $L = 1$, a large window size is required, which is not surprising when recalling Figure 8 showing that a mixture of gamma distributions is generally unimodal.

6 CONCLUSIONS

Second kind statistics seem to be an innovative and powerful tool for the study of distributions defined on \mathbb{R}^+ . The analytical formulation of the log-moments and the log-cumulants is indeed particularly simple and easy to exploit. At least, this is true for the examples presented in this article, whereof some, such as the mixture distributions, are not commonplace. Moreover, the variance of the estimators thus defined approaches the minimal values reached by the maximum likelihood method, while avoiding some of the analytical pitfalls. This approach shows great potential in certain applications in SAR image processing (such as the characterisation of an optimal homomorphic filter [11]). One can reasonably question why this approach, in all its simplicity, has not been proposed before. Several reasons can be called upon:

- The first is based on the observation that a Mellin transform of a pdf is only a Fourier transform of the same pdf taken on a logarithmic scale. Even if this step is perfectly justified on the theoretical level, all the possible advantages of moving into the Mellin domain remain hidden, such as the use of existing tables of known Mellin transforms, or the direct use of the log-moments and log-cumulants that produce better estimates of the distribution parameters.

4. The diagonal form of this matrix justifies *a posteriori* the analytical expression of the gamma distribution that we chose, which differs slightly from the one found in reference book like [9], [19].

TABLE 7

Standard deviation (SD) for the estimates of λ and ρ in the case of a mixture of gamma distributions. The size of the analysis window is 10×10 . Since the SD is inversely proportional to the square root of the number of samples in the analysis window, the table can serve to determine the optimal window after a maximum error has been set.

$\rho = 2$			$\rho = 5$			$\rho = 10$		
λ	SD_λ	SD_ρ	λ	SD_λ	SD_ρ	λ	SD_λ	SD_ρ
.1	.553	3.874	.1	.071	2.638	.1	.043	3.867
.2	.617	1.967	.2	.083	1.606	.2	.052	2.481
.3	.676	1.235	.3	.092	1.183	.3	.059	1.918
.4	.730	.784	.4	.099	.925	.4	.063	1.591
.5	.780	.455	.5	.105	.749	.5	.065	1.380

TABLE 8

Optimal window size for a relative error of 10% in λ and ρ .

$\rho = 2$			$\rho = 5$			$\rho = 10$		
λ	SD_λ	SD_ρ	λ	SD_λ	SD_ρ	λ	SD_λ	SD_ρ
.1	553×553	194×194	.1	71×71	53×53	.1	43×43	39×39
.2	309×309	98×98	.2	41×41	32×32	.2	26×26	25×25
.3	225×225	62×62	.3	31×31	24×24	.3	20×20	19×19
.4	182×182	39×39	.4	25×25	19×19	.4	16×16	16×16
.5	156×156	23×23	.5	21×21	15×15	.5	13×13	14×14

- The analysis of the product model, which is reserved for particular processes like coherent imaging, has not received the same strong attention as the additive signal model. The philosophy adopted for the study of the product model has too often consisted of transformation into logarithmic scale, in order to use the known tools for the additive model. This reductional step quickly pose problems, as it requires large control of the analytical expressions thus obtained. It is probably the reason why non-experts have written off other distributions than the gamma distribution and the inverse gamma distribution, such as the \mathcal{K} distribution, for instance.
- Finally, the Mellin transform has been completely ignored. Its applications has been confined to certain specialised applications, which has unfortunately prevented diffusion of the method beyond the field of study (e.g., radar and sonar signals, number theory, ultrasound propagation in heterogeneous media, the Fourier-Mellin transform in image processing). Even if certain pieces of work, old [15] as well as recent [22], [21] ones, have shown its applicability in the field of probability, its use has been very restricted. Therefore, few people know the fundamental properties, or even the exact definition.

The unfortunate consequence of the confidentiality is that few research groups have worked on the subject. Therefore, powerful and sufficiently general numerical implementations of the analytical transform are still missing. These would make it possible to consider numerical deconvolutions of the probability distributions described by a Mellin convolution, and thus to recover significant parameters of a SAR scene [23].

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